

*Baryon-Baryon correlation from heavy-ion collisions and its implication to baryon-baryon interaction*

**Akira Ohnishi (YITP, Kyoto U.)**

*FemTUM19, Sep. 12-13, 2019, Munich, Germany*

*K. Morita, T. Furumoto, AO, PRC91('15)024916 ( $\Lambda\Lambda$ )*

*AO, K. Morita, K. Miyahara, T. Hyodo, NPA954 ('16)294 ( $\Lambda\Lambda, K-p$ )*

*K. Morita, AO, F. Etminan, T. Hatsuda, PRC94('16), 031901(R) ( $\Omega N$ )*

*S. Cho et al.(ExHIC Collab.), Prog.Part.Nucl.Phys.95('17)279 ( $\Lambda\Lambda, K-p$ )*

*T. Hatsuda, K. Morita, AO, K. Sasaki, NPA967('17), 856 ( $\Xi-p$ )*

*K. Morita, S. Gongyo, T. Hatsuda, T. Hyodo, Y. Kamiya, AO,  
arXiv:1908.05414 ( $N\Omega, \Omega\Omega$ )*

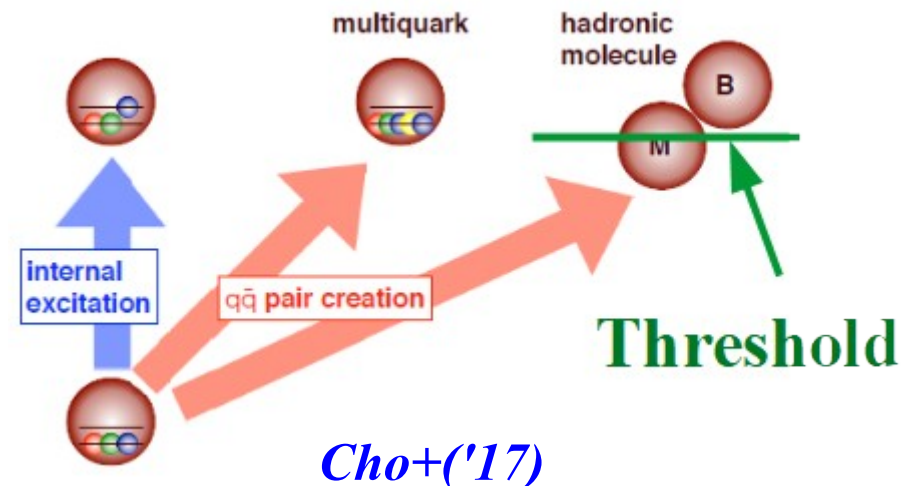


# Hadron Physics at High-Energy Colliders ?

- High-Energy Nuclear Collisions ( $\sqrt{s_{NN}}=40 \text{ GeV} - 14 \text{ TeV}$ ) as a Hadron Factory
  - Higgs, BSM, SUSY, QGP, ...
  - $dN/dy \sim 1000$  (RHIC, Au+Au)  $\rightarrow 10^3-10^5$  hadrons in one event
  - Various hadrons, nuclei ( $A \leq 4$ ) and anti-nuclei are formed.
  - Yield  $\sim$  Stat. Model calc.  
(Formation processes are too complicated to be out of statistical.)

- Trend in Hadron physics

- Quark-gluon structure of hadrons (Multi-quark or Hadronic molecule)
- Hadrons with heavy-quarks
- Hadron-Hadron interaction  $\Lambda N, \Sigma N, \Lambda \Lambda, \Xi N, \bar{K} N, \dots$
- Hadrons in nuclear matter / EOS of nuclear matter



# Outline

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- **Introduction: Hadron physics at High-Energy Colliders ?**
- **Correlation function and interaction**
  - **Two way to use the correlation function**
  - **Scattering length dependence of the correlation function**
- **Baryon-Baryon Correlation Function**
  - **Where is dibaryon ?**
  - **$p\Omega$  potential from lattice QCD**
  - **$p\Omega$  from STAR and ALICE**
  - **$\Lambda\Lambda$ ,  $p\Xi$  and  $\Omega\Omega$  correlations**
- **Summary**

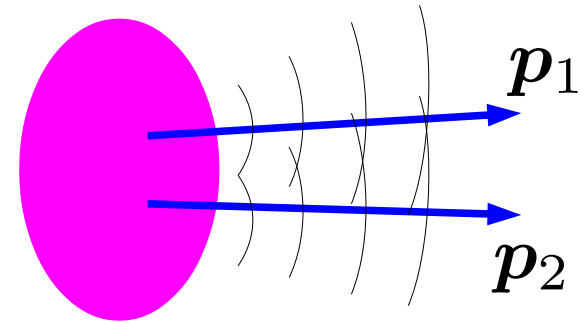
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# *Correlation Function and Interaction*

# Correlation Function

- Emitting source function

$$N_i(\mathbf{p}) = \int d^4x S_i(x, \mathbf{p})$$



- Two-particle momentum dist.

- Assumption: Two particles are produced independently, and the correlation is generated by the final state int.

*Koonin('77), Pratt+('86), Lednicky+('82)*

$$N_{12}(\mathbf{p}_1, \mathbf{p}_2) \simeq \int d^4x d^4y S_1(x, \mathbf{p}_1) S_2(y, \mathbf{p}_2) \underbrace{|\Psi_{\mathbf{p}_1, \mathbf{p}_2}(x, y)|^2}_{\text{two-body w.f.}}$$

$$\simeq \int d^4x d^4y S_1(x, \mathbf{p}_1) S_2(y, \mathbf{p}_2) \underbrace{|\varphi_{\mathbf{q}}(\mathbf{r})|^2}_{\text{relative w.f.}}$$

- Correlation function

$$C(\mathbf{p}_1, \mathbf{p}_2) = \frac{N_{12}(\mathbf{p}_1, \mathbf{p}_2)}{N_1(\mathbf{p}_1)N_2(\mathbf{p}_2)} \simeq \int d\mathbf{r} S_{12}(\mathbf{r}) |\varphi_{\mathbf{q}}(\mathbf{r})|^2$$

# Correlation Function

- **Example: Free identical bosons (spin 0, non-relativistic), Gaussian source (static, simultaneous, spherical)**

$$S(\mathbf{x}, \mathbf{p}) \propto \exp \left[ -\frac{\mathbf{x}^2}{2R^2} - \frac{\mathbf{p}^2}{2MT} \right]$$

$$S(\mathbf{x}, \mathbf{p}_1) S(\mathbf{y}, \mathbf{p}_2) \propto \exp \left[ -\frac{R_{\text{cm}}^2}{R^2} - \frac{\mathbf{r}^2}{4R^2} - \frac{\mathbf{P}^2}{4MT} - \frac{\mathbf{q}^2}{2\mu T} \right]$$

$$\begin{aligned} \Psi_{\mathbf{p}_1, \mathbf{p}_2}(\mathbf{x}, \mathbf{y}) &\propto \frac{1}{\sqrt{2}} \left[ e^{i\mathbf{p}_1 \cdot \mathbf{x} + i\mathbf{p}_2 \cdot \mathbf{y}} + e^{i\mathbf{p}_1 \cdot \mathbf{y} + i\mathbf{p}_2 \cdot \mathbf{x}} \right] \\ &= e^{i\mathbf{P} \cdot \mathbf{R}_{\text{cm}}} \times \sqrt{2} \cos \mathbf{q} \cdot \mathbf{r} \end{aligned}$$

- **Correlation function**

$$\begin{aligned} C(\mathbf{q}) &= (4\pi R^2)^{-3/2} \int d\mathbf{r} \exp \left[ -\frac{\mathbf{r}^2}{4R^2} \right] 2 \cos^2 \mathbf{q} \cdot \mathbf{r} \\ &= 1 + \exp(-4q^2 R^2) \end{aligned}$$

**Correlation Function  $\rightarrow$  Source Size**

# How can we measure the radius of a star ?

## Two photon intensity correlation

*Hanbury Brown & Twiss, Nature 10 (1956), 1047.*

- Simultaneous two photon observation probability is enhanced from independent emission cases  
 → angular diameter of Sirius=6.3 msec

**Recent data (Wikipedia)**  
**5.936±0.016 msec**

A TEST OF A NEW TYPE OF STELLAR INTERFEROMETER ON SIRIUS

By R. HANBURY BROWN

Jodrell Bank Experimental Station, University of Manchester

AND

DR. R. Q. TWISS

Services Electronics Research Laboratory, Baldock

NATURE November 10, 1956 VOL. 178

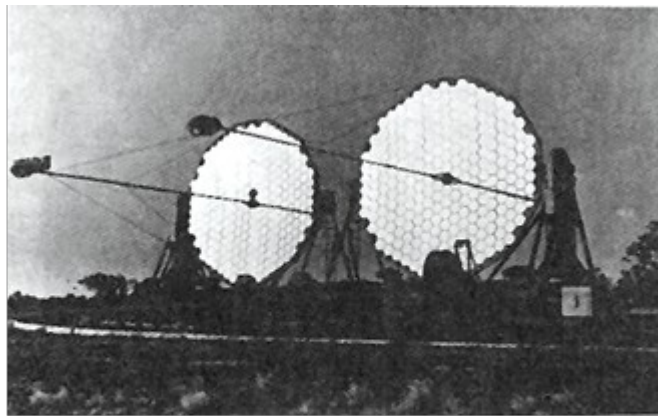


Figure 2. Picture of the two telescopes used in the HBT experiments. The figure was extracted from Ref.[1].

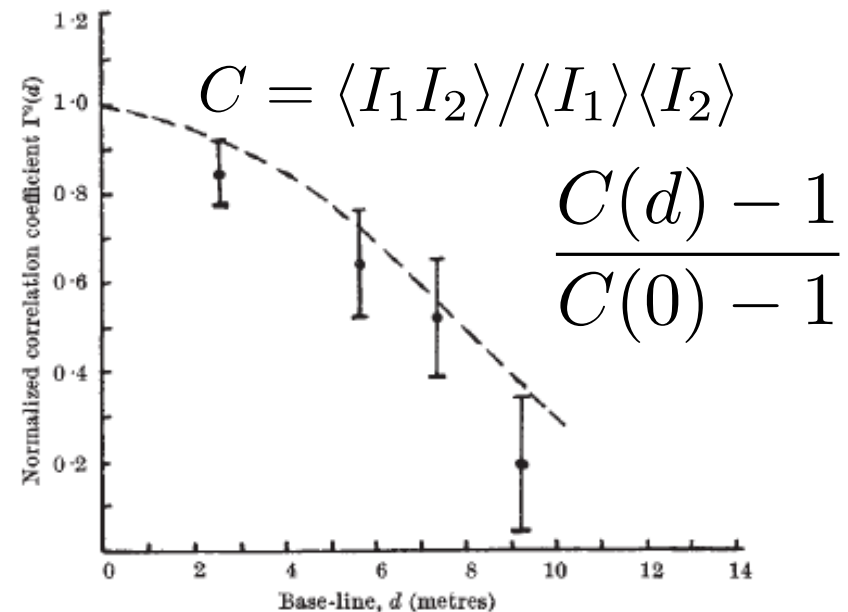


Fig. 2. Comparison between the values of the normalized correlation coefficient  $\Gamma^2(d)$  observed from Sirius and the theoretical values for a star of angular diameter  $0.0063''$ . The errors shown are the probable errors of the observations

*HBT telescope (from Goldhaber, ('91))*

*HBT ('56)*

# Two particle intensity correlation

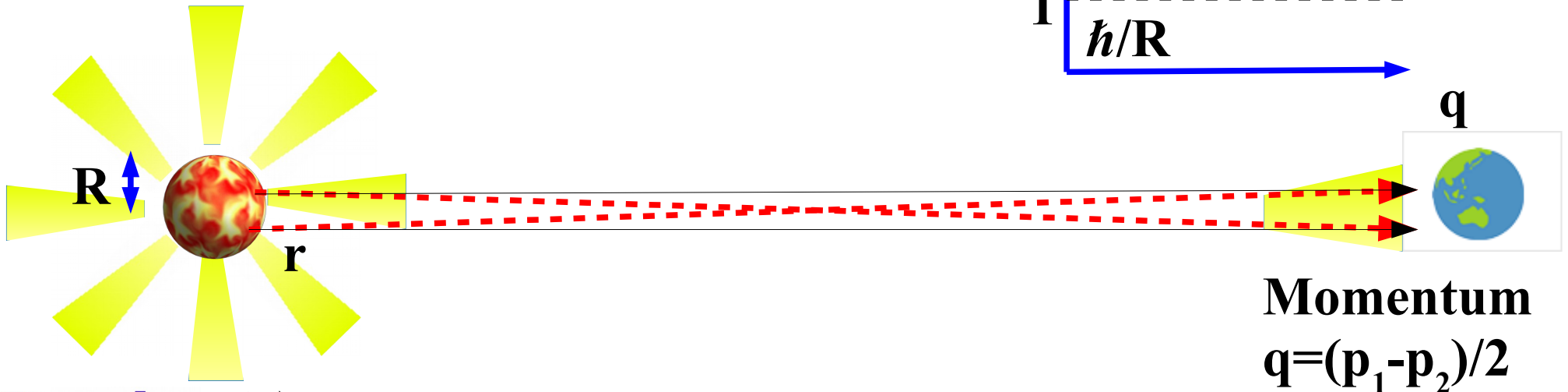
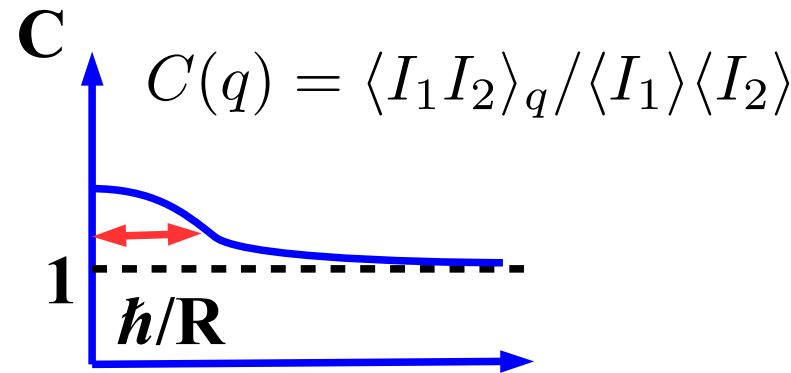
- Wave function symmetrization from quantum statistics

$$C(\mathbf{q}) = \int d^3r S(\mathbf{q}, \mathbf{r}) \left| \frac{1}{\sqrt{2}} (e^{i\mathbf{q}\cdot\mathbf{r}} + e^{-i\mathbf{q}\cdot\mathbf{r}}) \right|^2 \simeq \underline{1 + \exp(-4q^2 R^2)}$$

Source fn.  
( $\mathbf{r}$ =relative coordinate)  
(symmetrized w.f.)<sup>2</sup>

Static spherical source case

→ Small relative momenta are favored due to symmetrization of the relative wave function.



Momentum  
 $\mathbf{q} = (\mathbf{p}_1 - \mathbf{p}_2)/2$



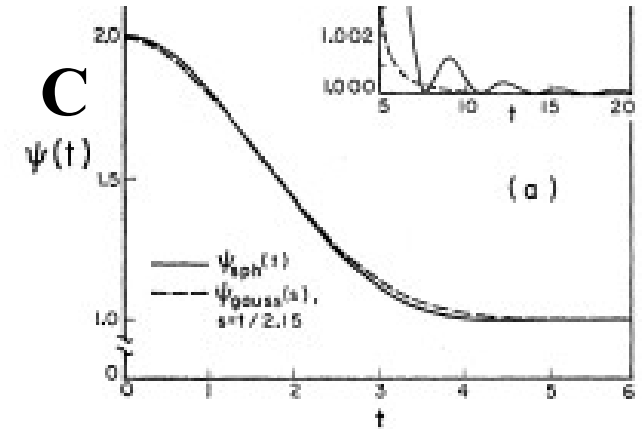
# How can we measure source size in nuclear reactions ?

## ■ Two pion interferometry

*G. Goldhaber, S. Goldhaber, W. Lee, A. Pais, Phys. Rev. 120 (1960), 300*

- Two pion emission probability is enhanced at small relative momenta

→ Pion source size  $\sim 0.75 \hbar / \mu c$



$q$  (relative momentum)

PHYSICAL REVIEW

VOLUME 120, NUMBER 1

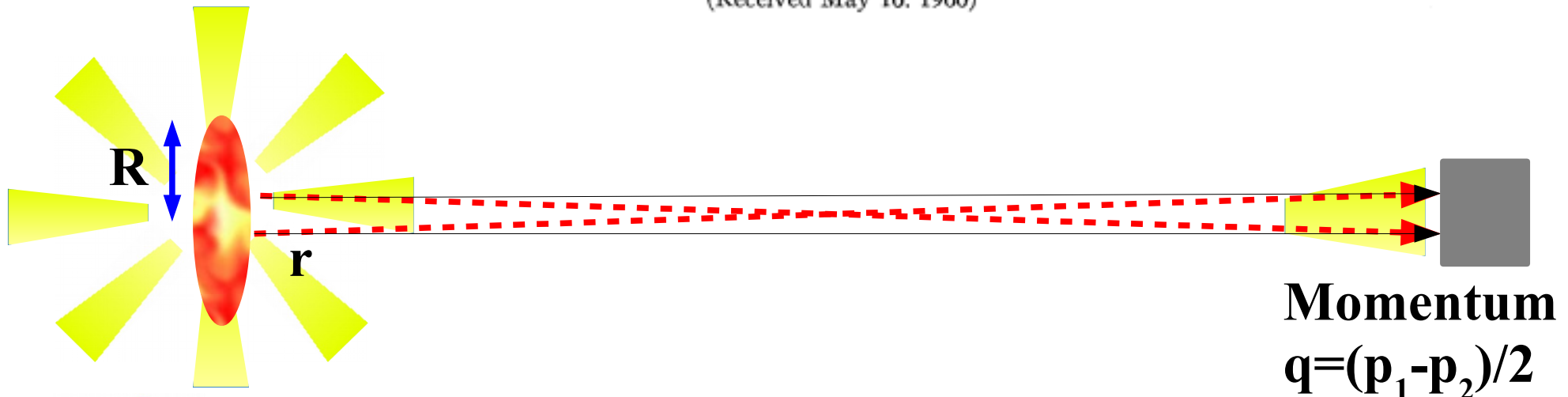
OCTOBER 1, 1960

## Influence of Bose-Einstein Statistics on the Antiproton-Proton Annihilation Process\*

GERSON GOLDHABER, SULAMITH GOLDHABER, WONYONG LEE, AND ABRAHAM PAIS†

*Lawrence Radiation Laboratory and Department of Physics, University of California, Berkeley, California*

(Received May 16, 1960)



# Another usage of correlation function

- **HBT, GGLP: Corr. Fn. + w.f. → Source Size**  
**Another way: Corr. Fn. + Source Size**  
→ wave function → hadron-hadron interaction
- **Effect of hadron-hadron interaction on the wave function**
  - **Assumption: Only s-wave (L=0) is modified.**
  - **Non-identical particle pair, Gauss source.**

$$\varphi_{\mathbf{q}}(\mathbf{r}) = e^{i\mathbf{q}\cdot\mathbf{r}} - j_0(qr) + \chi_q(r)$$

$$\rightarrow C(\mathbf{q}) = \int d\mathbf{r} S(r) |\varphi_{\mathbf{q}}(\mathbf{r})|^2$$

$$= 1 + \int d\mathbf{r} S(r) \{ |\chi_q(r)|^2 - |j_0(qr)|^2 \}$$

*K. Morita, T. Furumoto, AO, PRC91('15)024916*

*Corr. Fn. shows how much squared w. f. is enhanced  
→ Large CF is expected with attraction*

# Wave function around threshold (S-wave, attraction)

## Low energy w.f. and phase shift

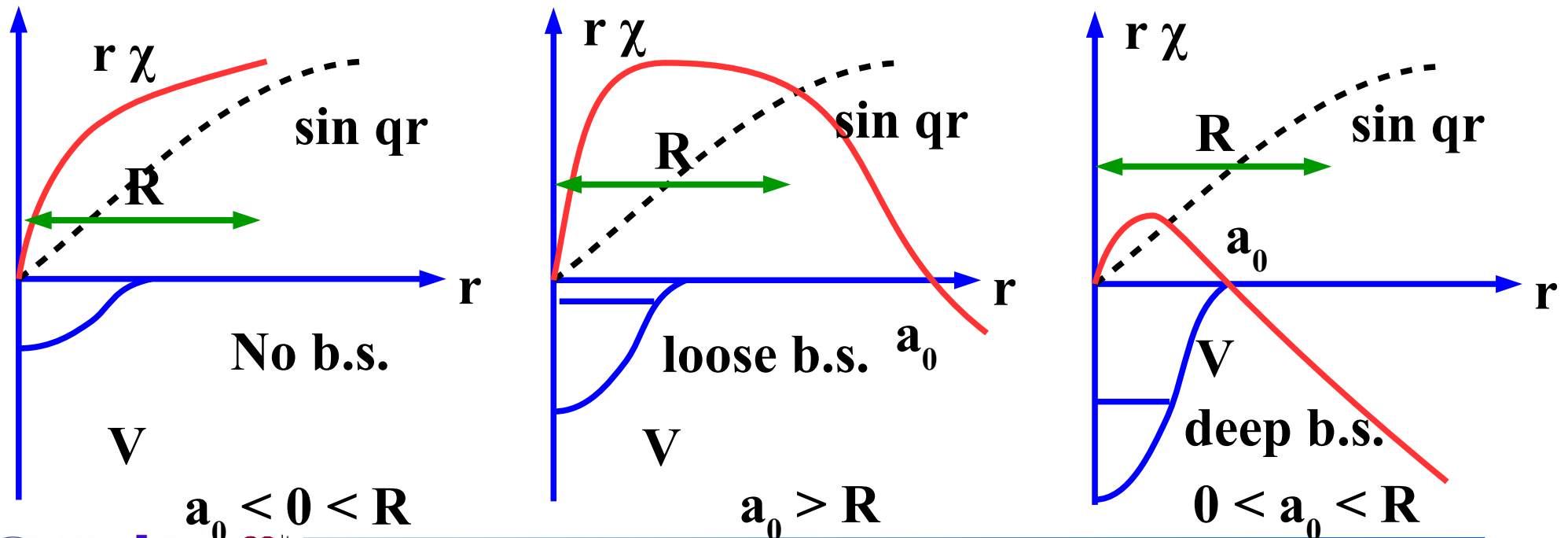
$$u(r) = qr\chi_q(r) \rightarrow \sin(qr + \delta(q)) \sim \sin(q(r - a_0))$$

$$q \cot \delta = -\frac{1}{a_0} + \frac{1}{2}r_{\text{eff}}q^2 + \mathcal{O}(q^4) \quad (\delta \sim -a_0q)$$

$a_0$  = scatt. length

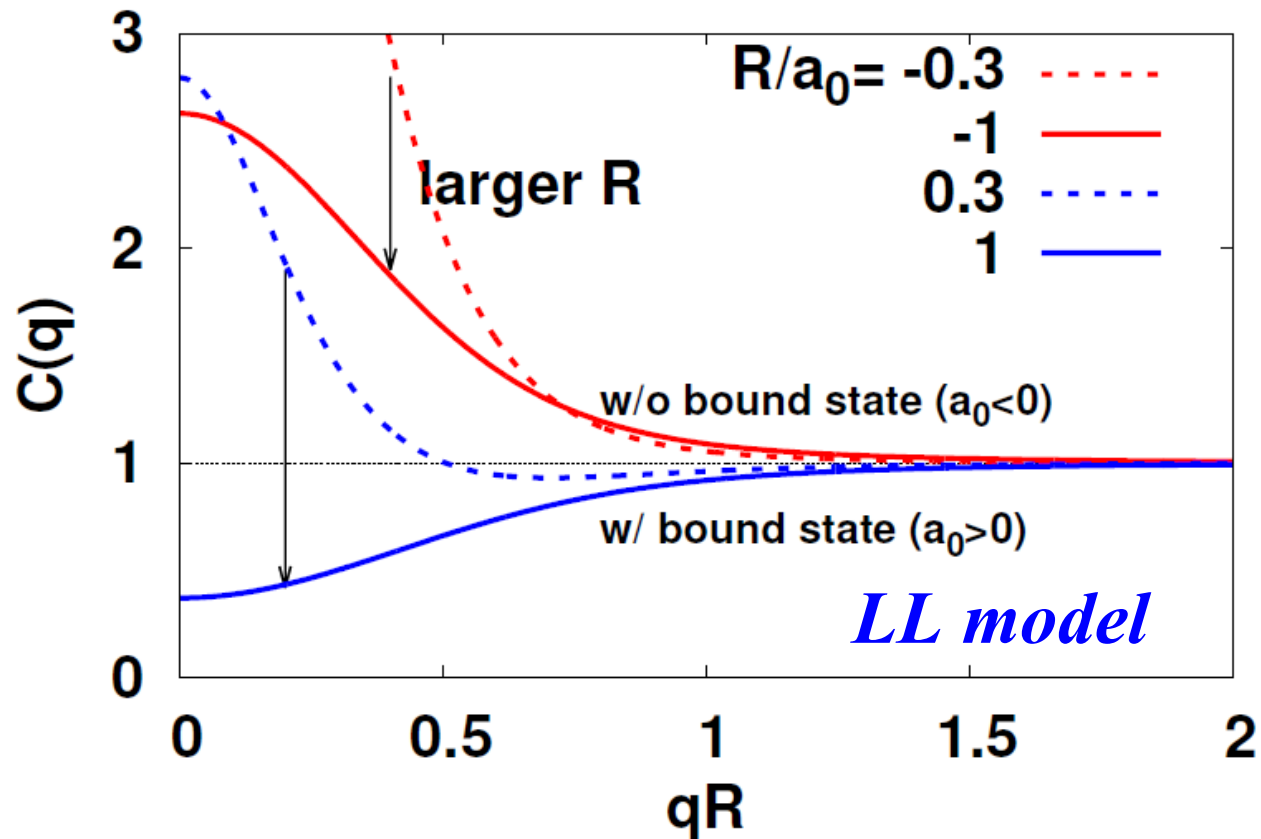
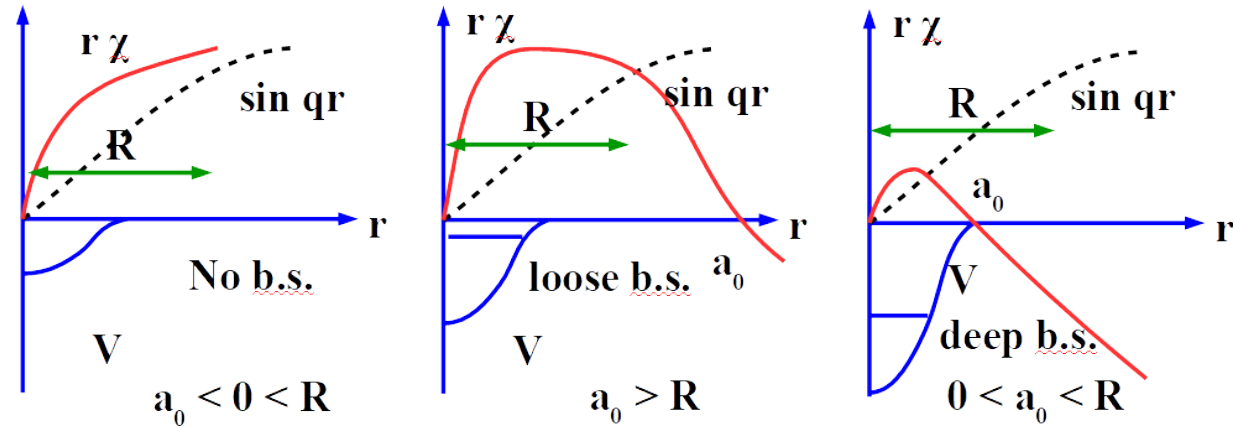
$r_{\text{eff}}$  = eff. range

- Wave function grows rapidly at small  $r$  with attraction.
- With a bound state ( $a_0 > 0$ ), a node appears around  $r = a_0$

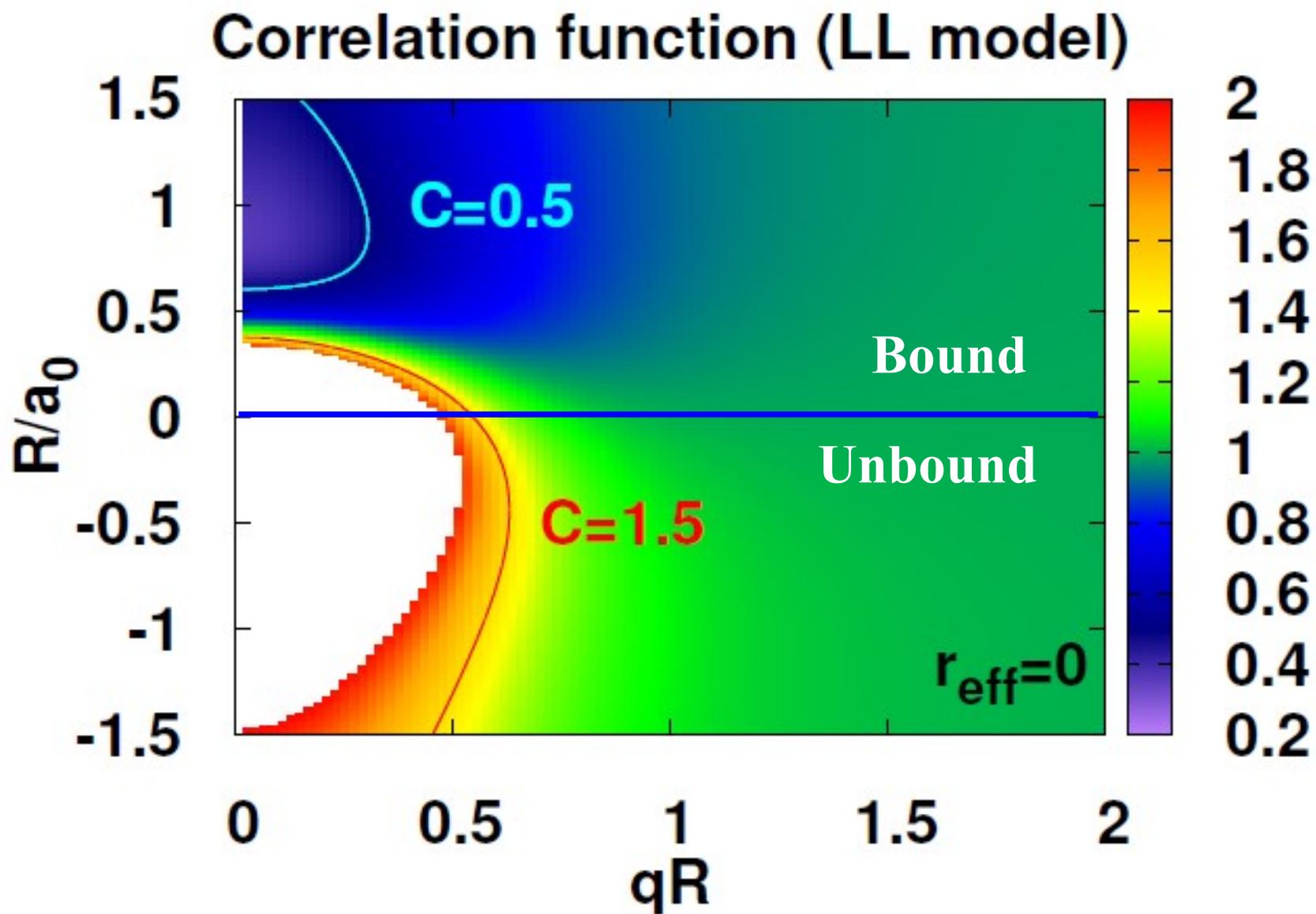


# From correlation function to hadron-hadron interaction

- Large  $|a_0|$  ( $|a_0| > R$ )  
→ Large  $C(q)$   
(unitary regime)
- w/o bound state  
( $a_0 < 0, |a_0| \sim R$ )  
→  $C(q) > 1$
- With bound state  
( $a_0 > 0, |a_0| \sim R$ )  
→ Region with  $C(q) < 1$  appears



# Correlation Function in LL model



*LL model: R. Lednicky, V. L. Lyuboshits ('82)*

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# *Baryon-Baryon Correlation Function (mainly on $p\Omega$ correlation)*

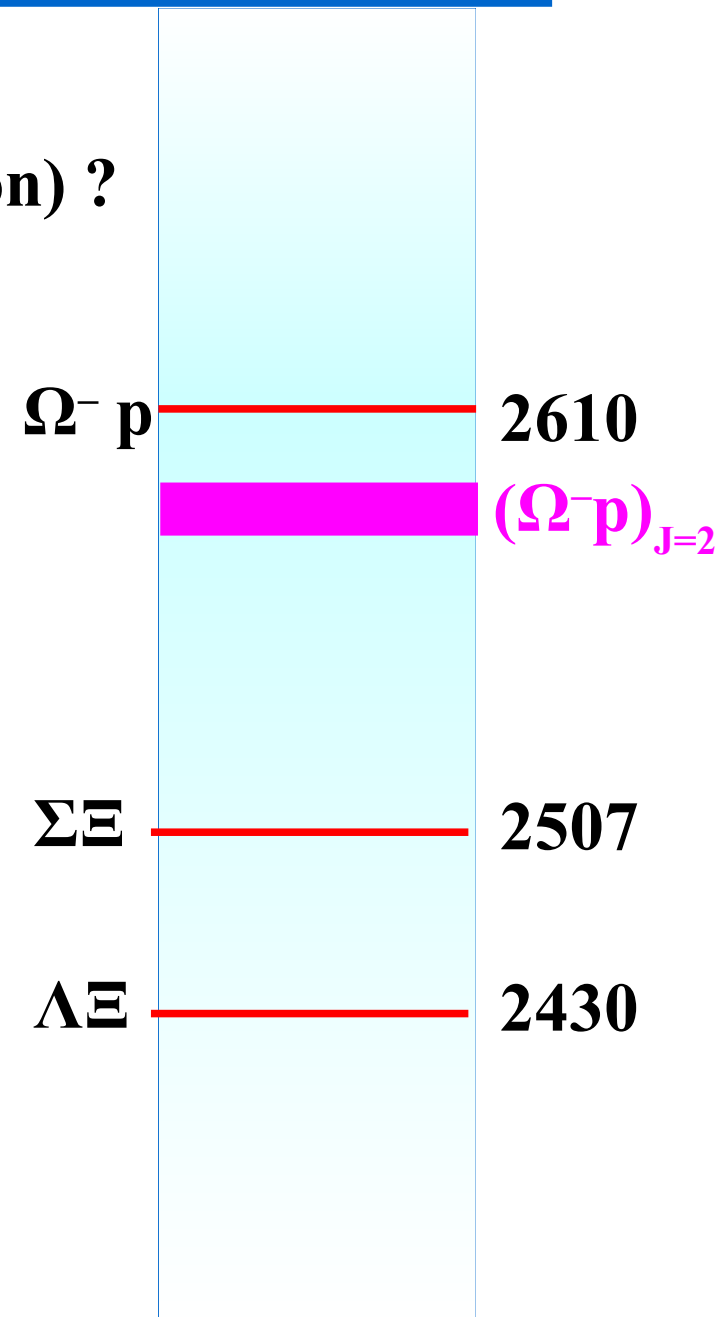
# Where is dibaryon ?

- Deuteron = First dibaryon (pn bound state)
  - H-particle: 6-quark state (uuddss)
    - Predicted (*Jaffe ('77)*), Ruled-out ( $\Lambda\Lambda$  nucl., *Takahashi+('01)*), Suggested as a resonance in exp. (*Yoon+ ('07)*) or as a bound state of  $\Xi N$  (*HAL QCD ('16)*)
  - Dibaryon would appear in channels, where *Oka ('88)*, *Gal ('16)*
    - The Pauli blocking of quarks does not operate,
    - and the Color-magnetic interaction is attractive
- Examples: H(= $\Lambda\Lambda$ - $N\Xi$ - $\Sigma\Sigma$ ),  $N\Omega$ ,  $N\Sigma^*$ ,  $d^*(=\Lambda\Lambda)$ .

*Let us examine the existence of dibaryon states by using the correlation function !*

# $\Omega N$ dibaryon

- $\Omega$  : sss,  $J\pi=3/2^+$ ,  $M=1672$  MeV
- Is there an  $\Omega N$  bound state ( $S=-3$  dibaryon) ?
  - Predicted as a dibaryon candidate  
*Goldman+ ('87), Oka ('88), Gal ('16)*
  - Lattice QCD predicts a bound state with narrow width for  $J=2$  ( $^5S_2$ )  
*Etminan+ (HAL QCD)('14), Iritani+ (HAL QCD) ('19)*
  - Correlation function is measurable !  
*Adam+ (STAR)('19), ALICE, in prep.*





# $\Omega N$ potential from lattice QCD

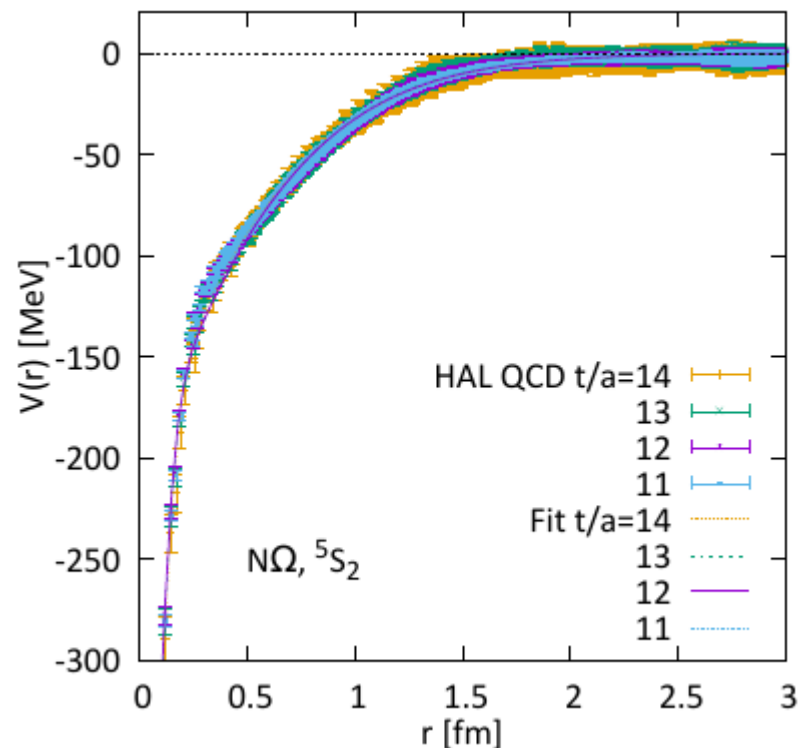
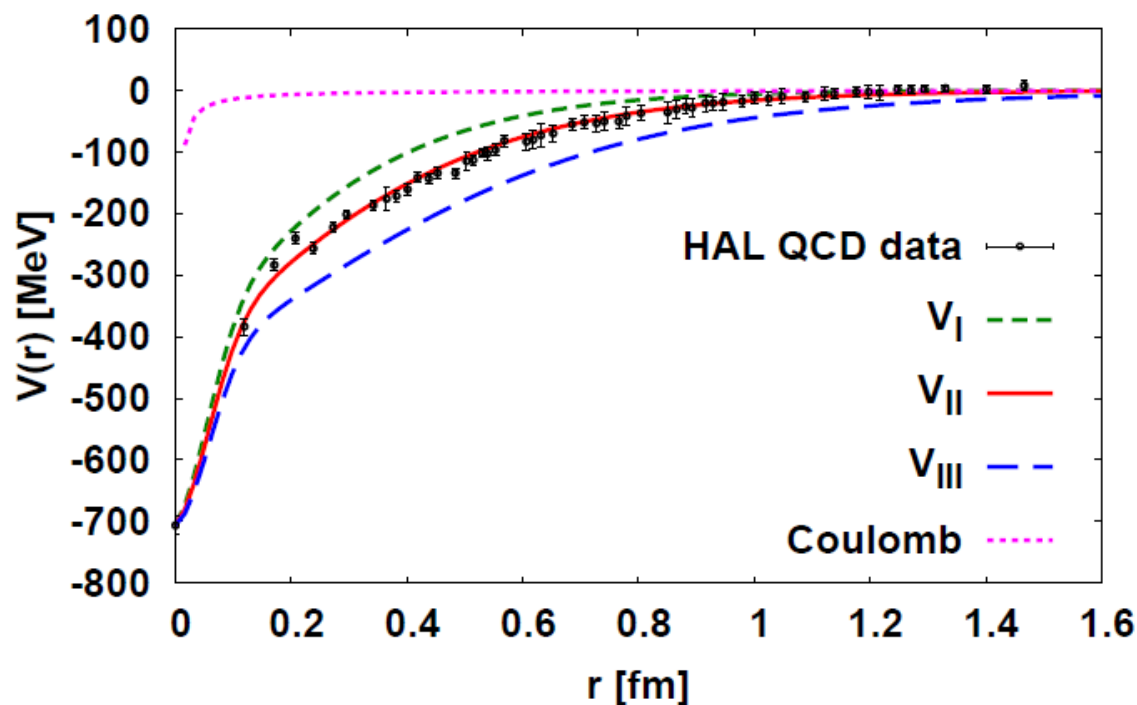
## ■ $\Omega N$ potential by HAL QCD Collab. (J=2)

- $m_\pi=875$  MeV, B.E.~ 19 MeV

*F. Etminan et al. (HAL QCD Collab.), NPA928('14)89.*

- $m_\pi=146$  MeV, B.E.~ 2.2 MeV

*T. Iritani et al. (HAL QCD Collab.), PLB 792('19)284.*



# Calculation Details

## ■ $N\Omega$ potential from HAL QCD Collab.

*Etminan+(HAL QCD) ('14), Iritani+ (HAL QCD)('19)*

- $J=1$  potential is uncertain  $\rightarrow$  Three models

Strong abs. at  $r < r_0$  ( $r_0 \sim 2$  fm) (*Morita+('16)*) (Standard)

Complete absorption  $\chi(J=1) = 0$  (Minimum)

Same w.f. as that with  $J=2$ ,  $\chi(J=1) = \chi(J=2)$  (Reference)

- Statistical Error can be evaluated by using Jackknife potentials.

## ■ Coulomb potential enhances CF even without strong int.

$\rightarrow$  Small-Large ratio of CF (*Morita+('16)*)

- Large source  $\rightarrow$  Coulomb force dominate

Small source  $\rightarrow$  Visible strong interaction effects

## ■ Source function: Blast wave, Gaussian source

# Emission Source Function

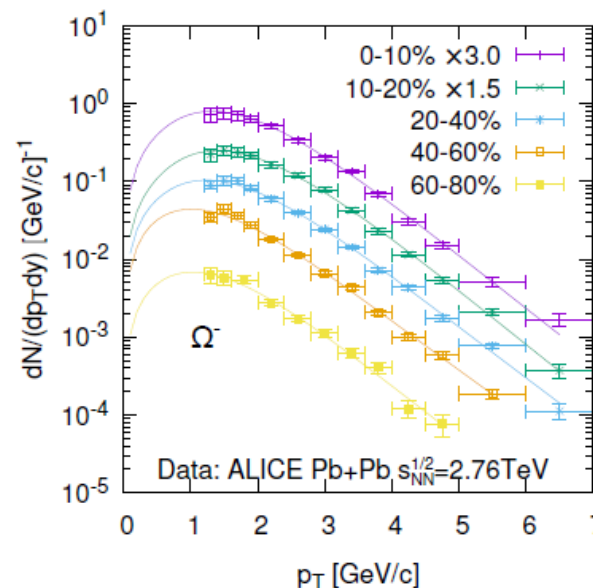
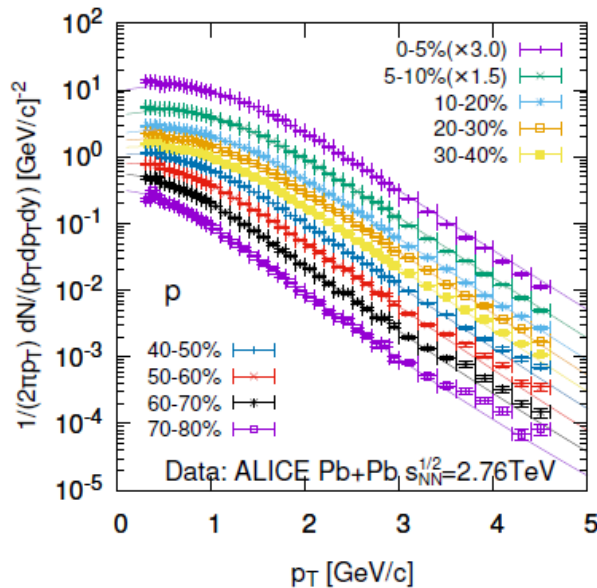
- Gaussian Source  $\propto \exp(-r^2/4R^2)$ ,  $R=(0.8-4)$  fm [Simple and convenient]
- Expanding source model [Reasonably realistic]

Flow velocity

$$u_\mu(x)$$

$$d^4x S_i(x, p) = \tau_0 d\eta_s d^2r_T \frac{d}{(2\pi)^3} \underbrace{n_f(u \cdot p, T)}_{\text{Fermi dist.}} \exp\left(-\frac{r_T^2}{2R_T^2}\right)$$

Fermi dist.

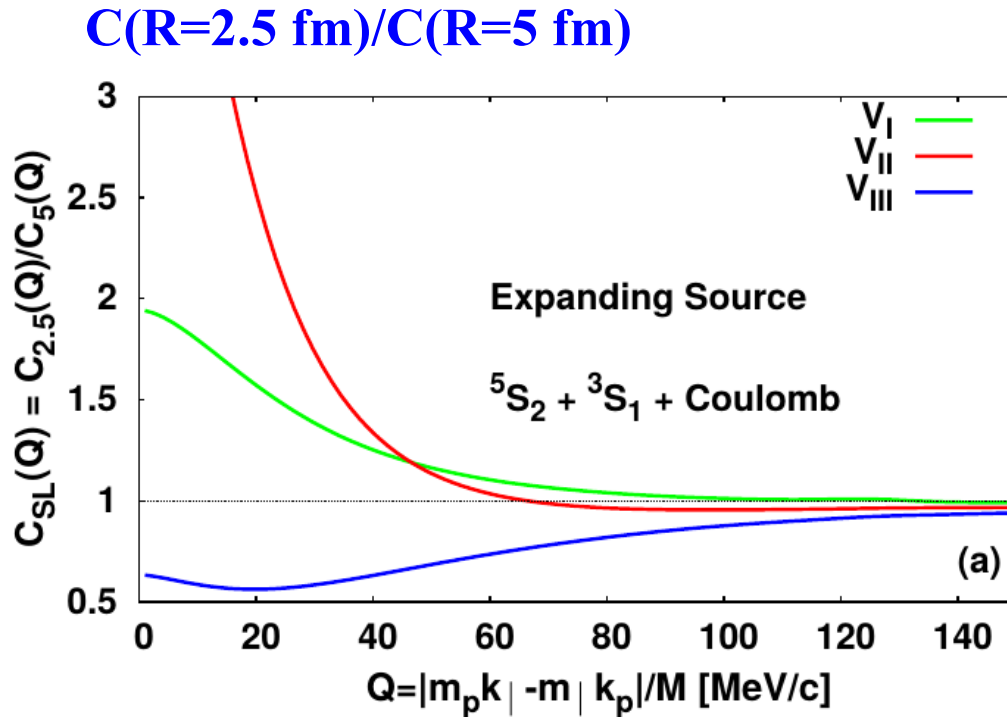


Morita+('19)

- Transport model result [should be realistic] → Future work

# Comparison with STAR data

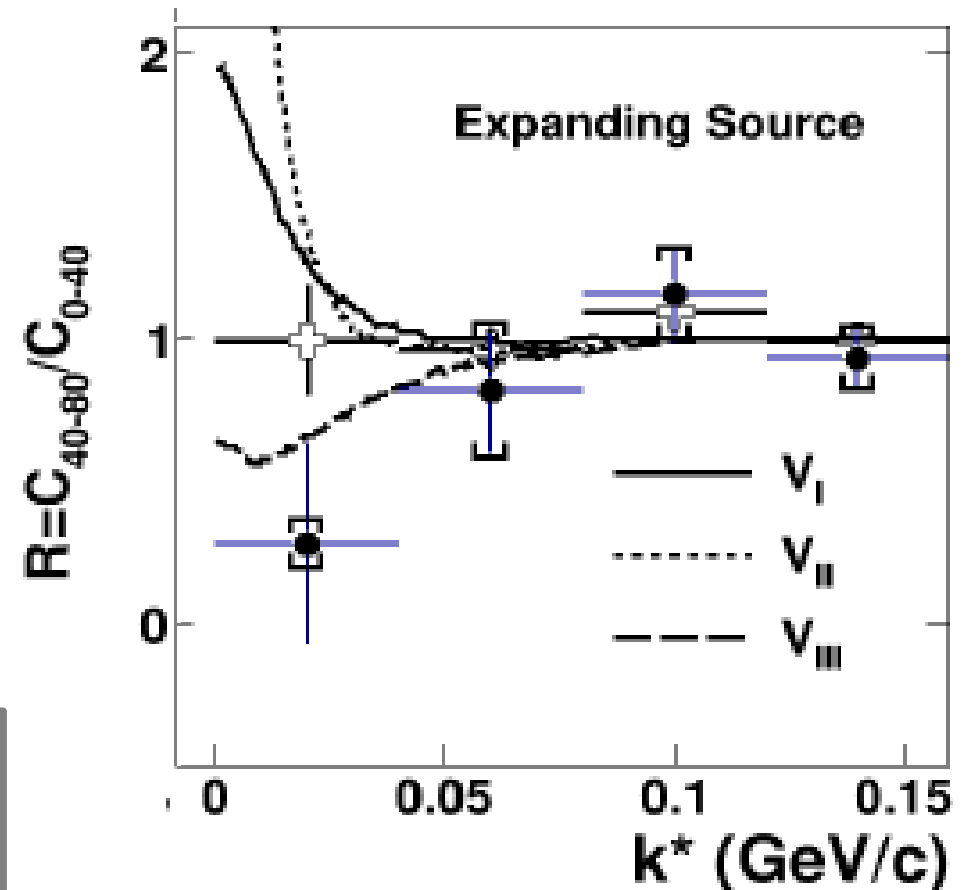
## ■ Dip structure in Small-Large ratio of CF



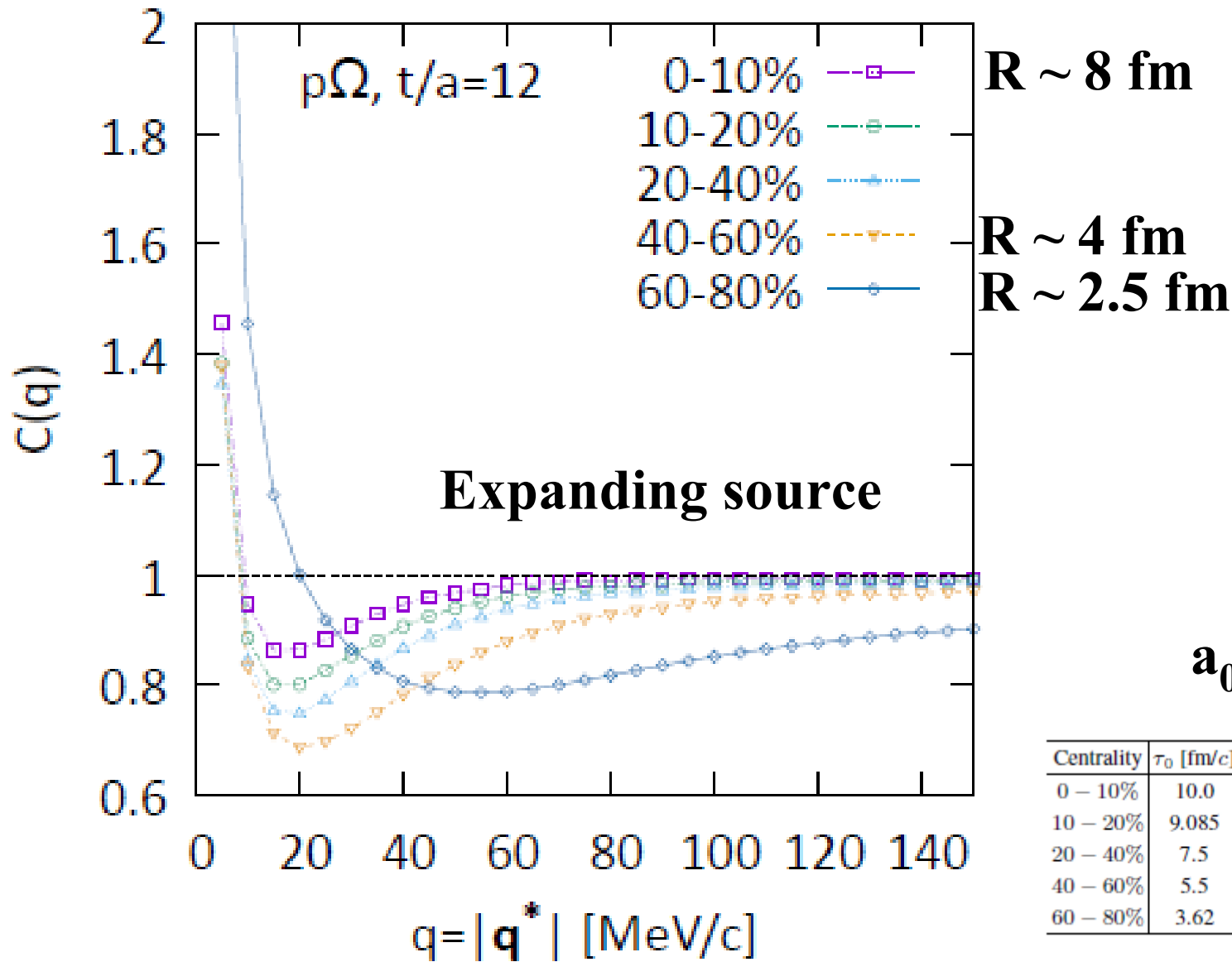
*Morita, AO, Etminan, Hatsuda ('16)*

*We may have a dibaryon state  
in  $\Omega N$  channel*

*STAR (1808.02511,  
PLB790 ('19) 490)*



# Source Size Dependence of Correlation Function



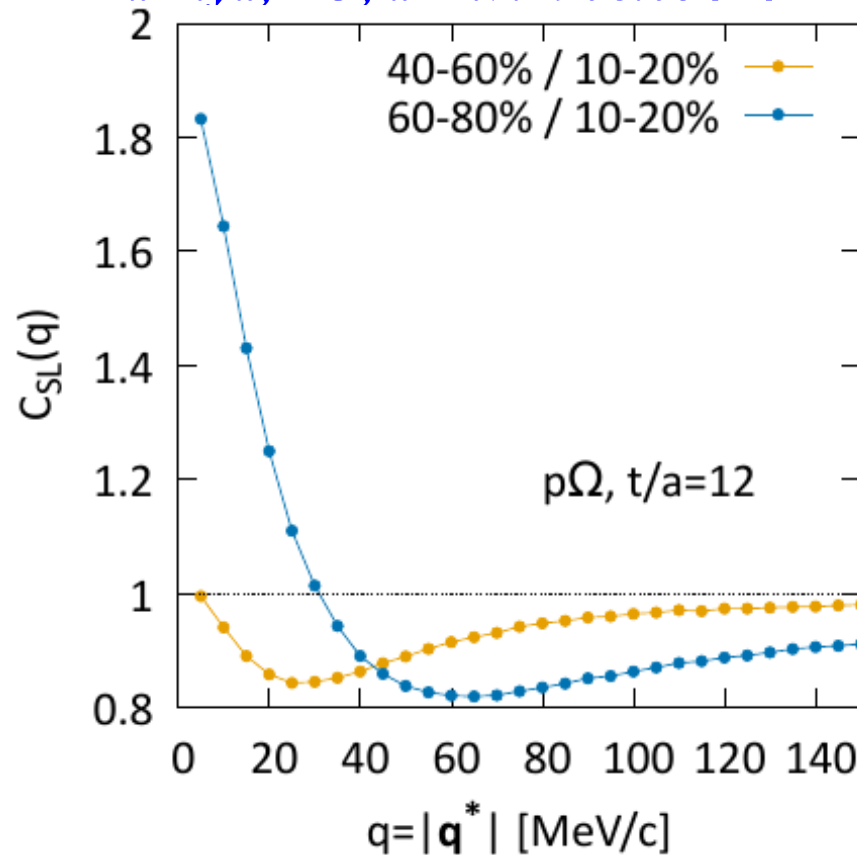
Centrality	$\tau_0$ [fm/c]	$R_T^\Omega$ [fm]	$R_T^p$	$\alpha^\Omega$	$\beta^\Omega$	$\alpha^p$	$\beta^p$
0 – 10%	10.0	8.0	6.8	0.584	0.628	0.759	0.421
10 – 20%	9.085	6.75	6.23	0.618	0.579	0.750	0.425
20 – 40%	7.5	5.88	5.2	0.546	0.692	0.707	0.466
40 – 60%	5.5	4.38	3.92	0.444	0.858	0.604	0.6
60 – 80%	3.62	2.12	2.66	0.456	0.812	0.456	0.82

*K. Morita, S. Gongyo, T. Hatsuda,  
T. Hyodo, Y. Kamiya, AO ('19)*

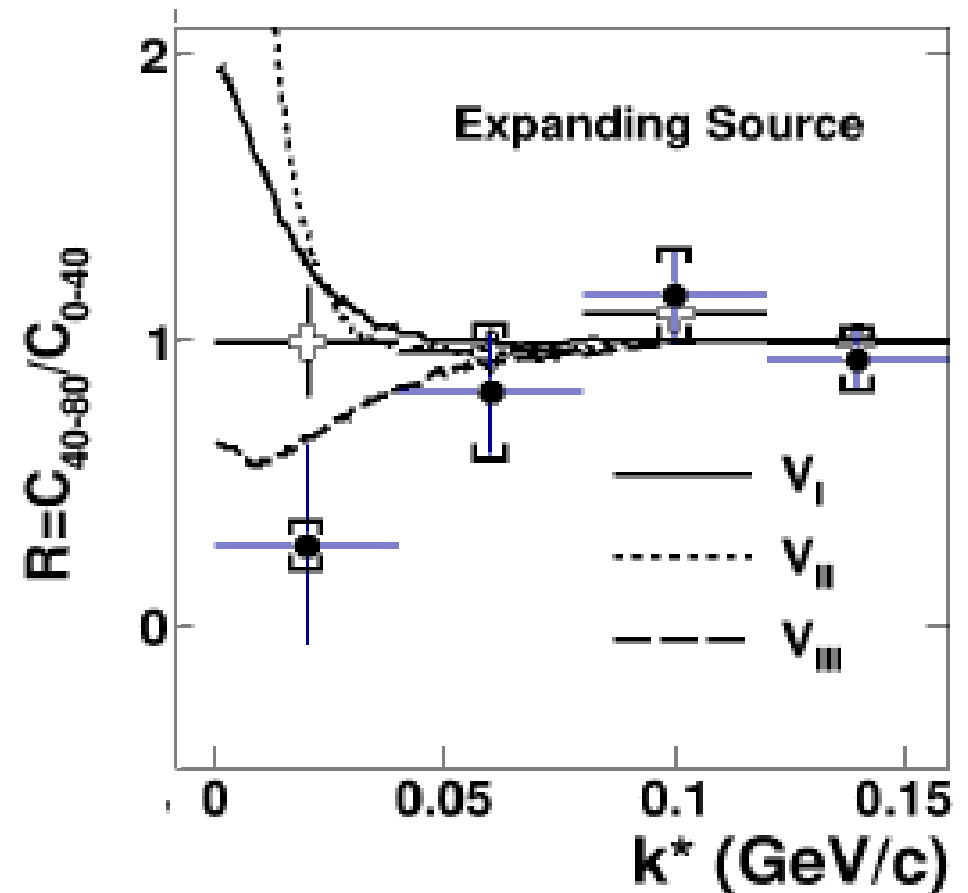
# Comparison with STAR data

- Results with potential at nearly physical quark mass (= between  $V_{II}$  and  $V_{III}$ )
- Dip is seen but is not deep enough to explain STAR data.

Morita, Gongyo, Hatsuda, Hyodo, Kamiya, AO, arXiv:1908.05414

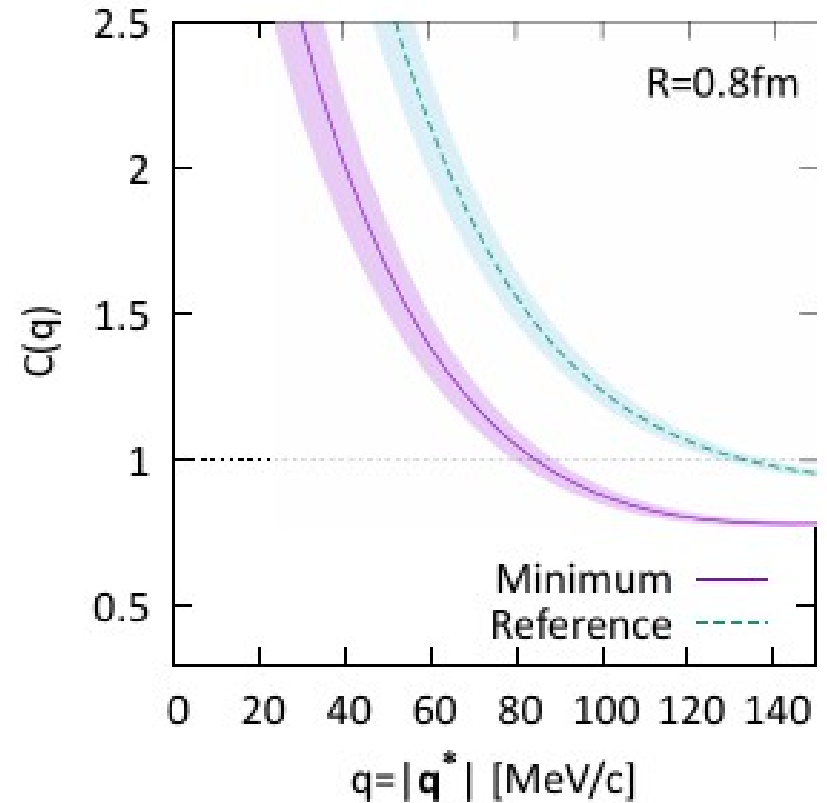
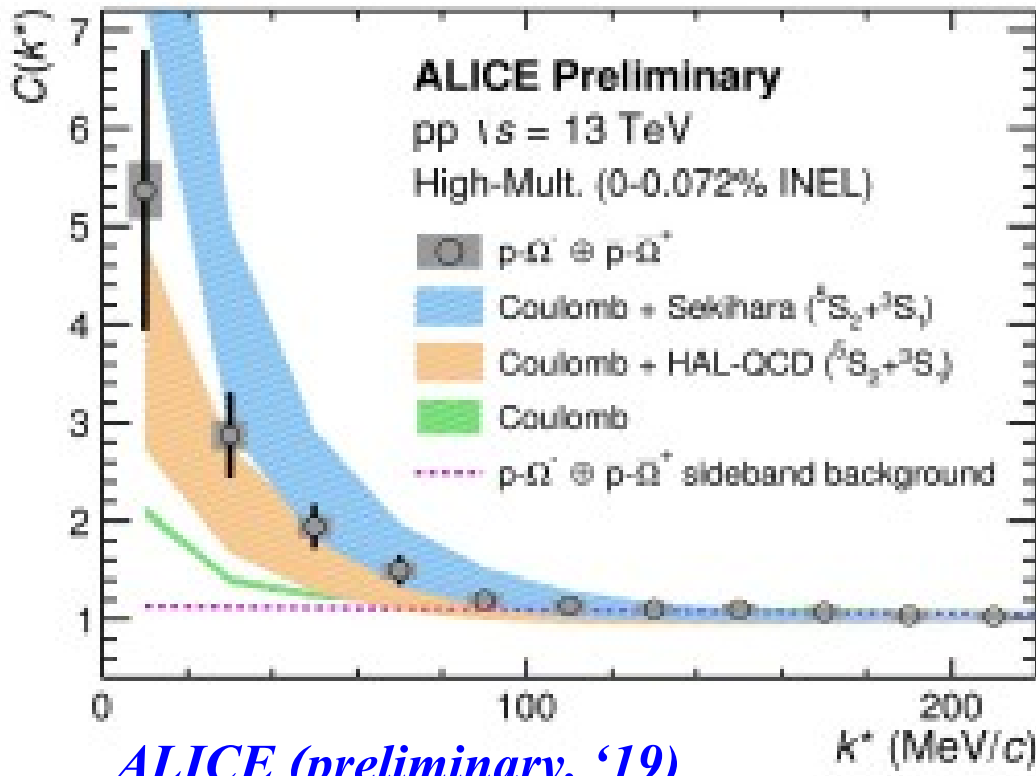


STAR (1808.02511, PLB790 ('19) 490)



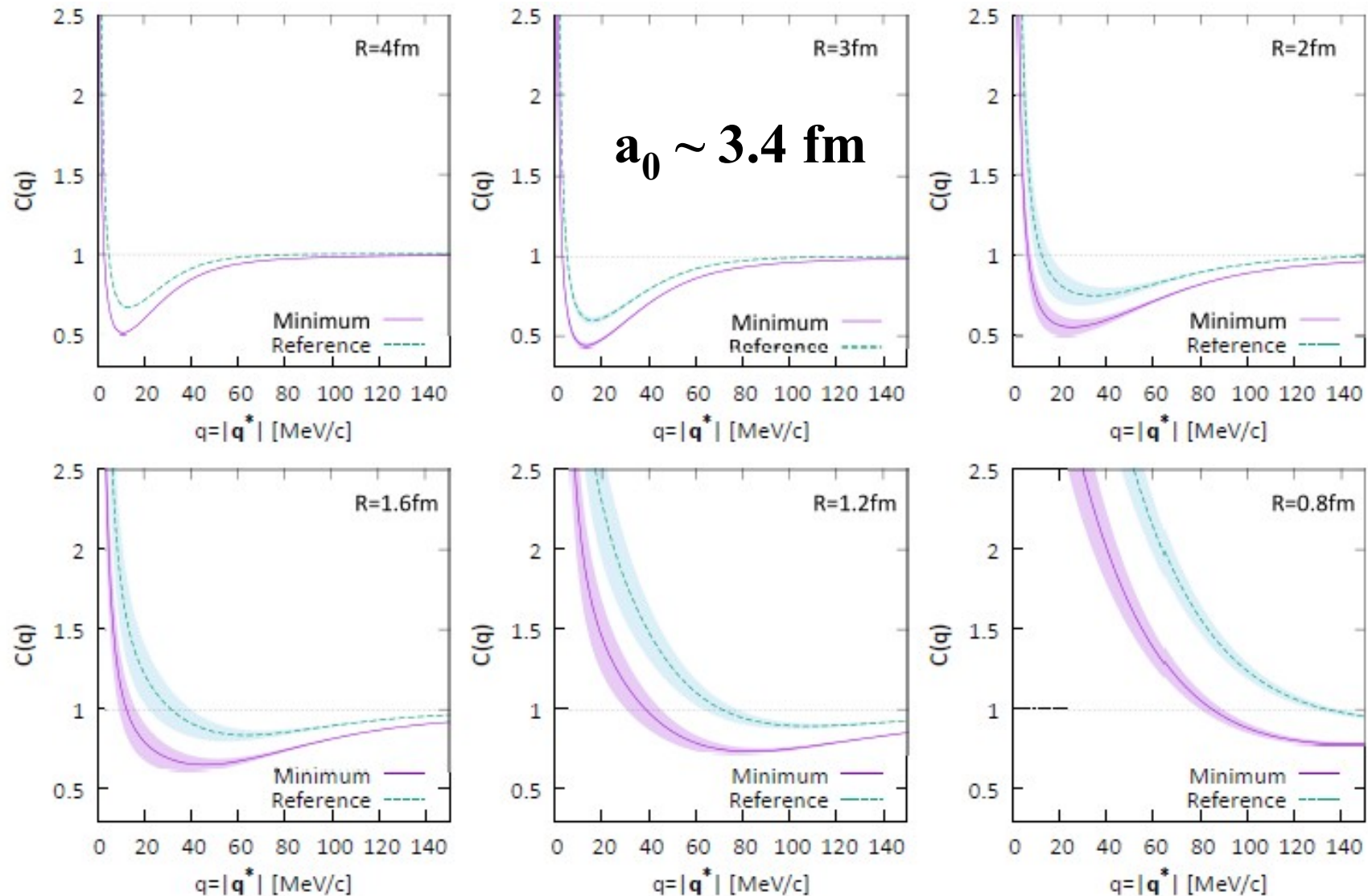
# ALICE Preliminary

- pp 13 TeV high-multiplicity events in ALICE  
 → Strong enhancement of CF at small  $q$   
*O. Vázquez Doce et al. (ALICE), Hadrons 2019*



*Morita+, arXiv:1908.05414*

# Source Size Dependence of Correlation Function



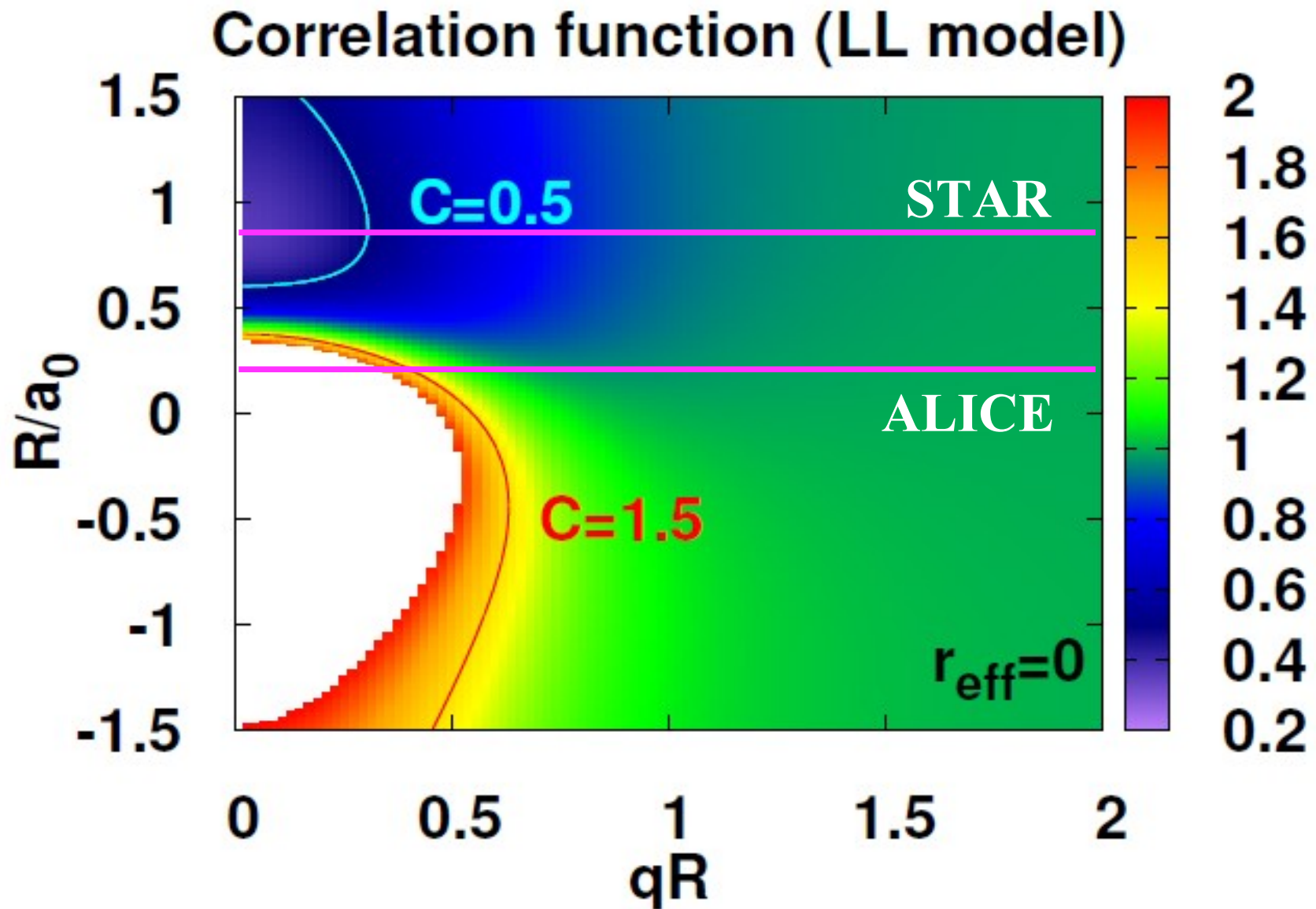
Gaussian Source

*K. Morita, S. Gongyo, T. Hatsuda,  
T. Hyodo, Y. Kamiya, AO ('19)*

*A. Ohnishi, FemTUM, Sep. 12, 2019* 24

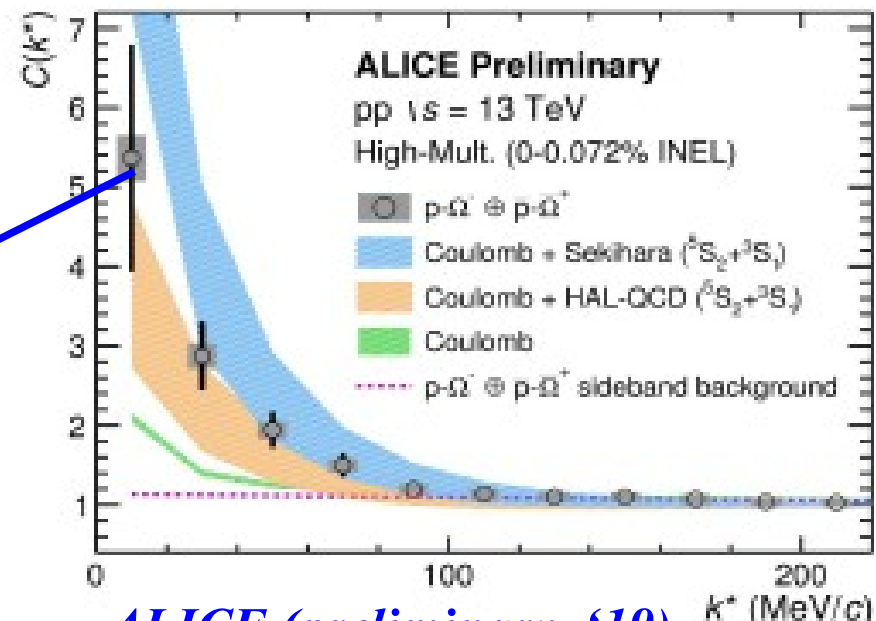
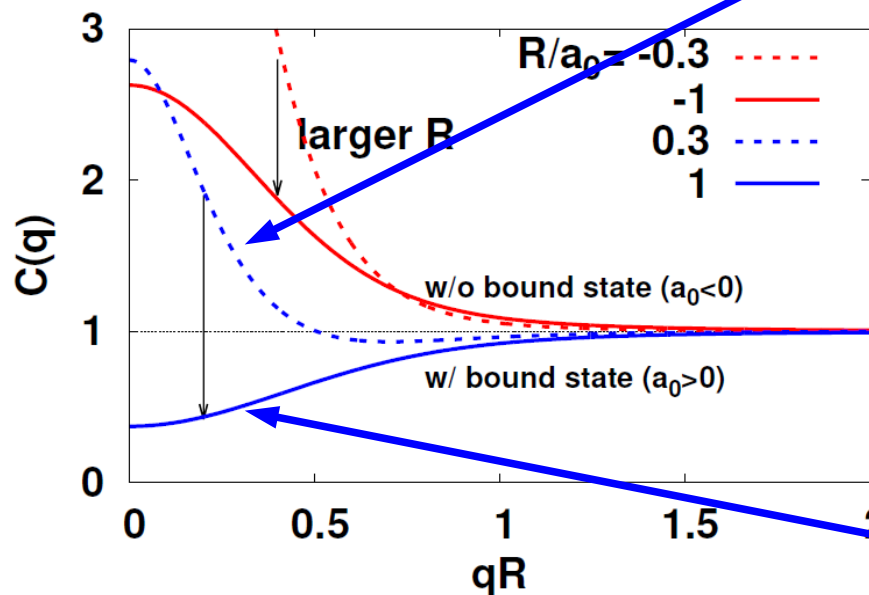


# Correlation Function in LL model

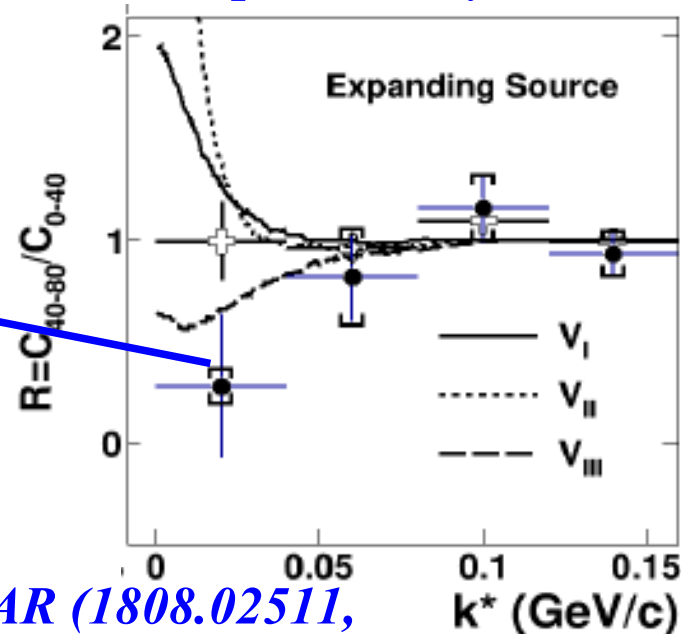


$a_0$  ( $p\Omega$ )  $\sim$  3.4 fm,  $R(\text{ALICE}) \sim 0.7$  fm,  $R(\text{STAR}) \sim 3$  fm

# STAR + ALICE = $N\Omega$ Dibaryon



ALICE (preliminary, '19)



STAR (1808.02511, PLB790 ('19) 490)

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*Do I have time ?*

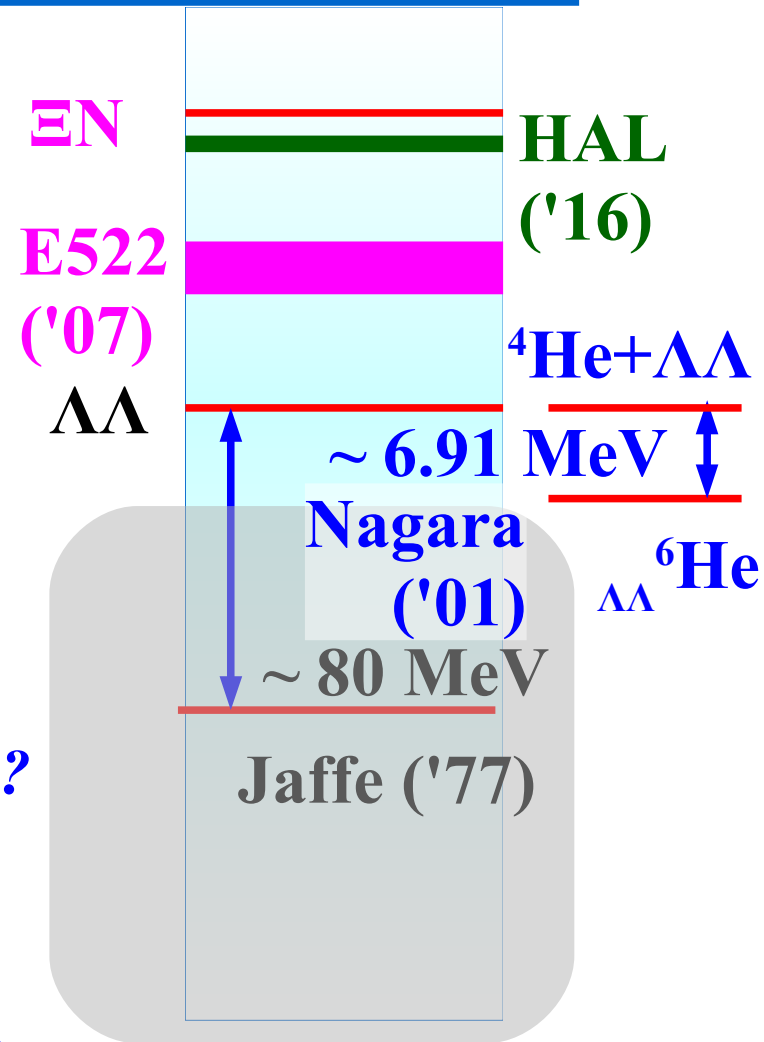
# Relevance of $\Lambda\Lambda$ interaction to physics

## ■ H-particle: 6-quark state (uuddss)

- Prediction: *R.L.Jaffe, PRL38(1977)195*
- Ruled-out by double  $\Lambda$  hypernucleus  
*Takahashi et al., PRL87('01) 212502*
- Resonance or Bound “H” ?  
*Yoon et al. (KEK-E522)+AO ('07)*
- Lattice QCD  
*HAL QCD & NPLQCD ('11)*  
*HAL QCD ('16): H as a loosely bound  $\Xi N$  ?*

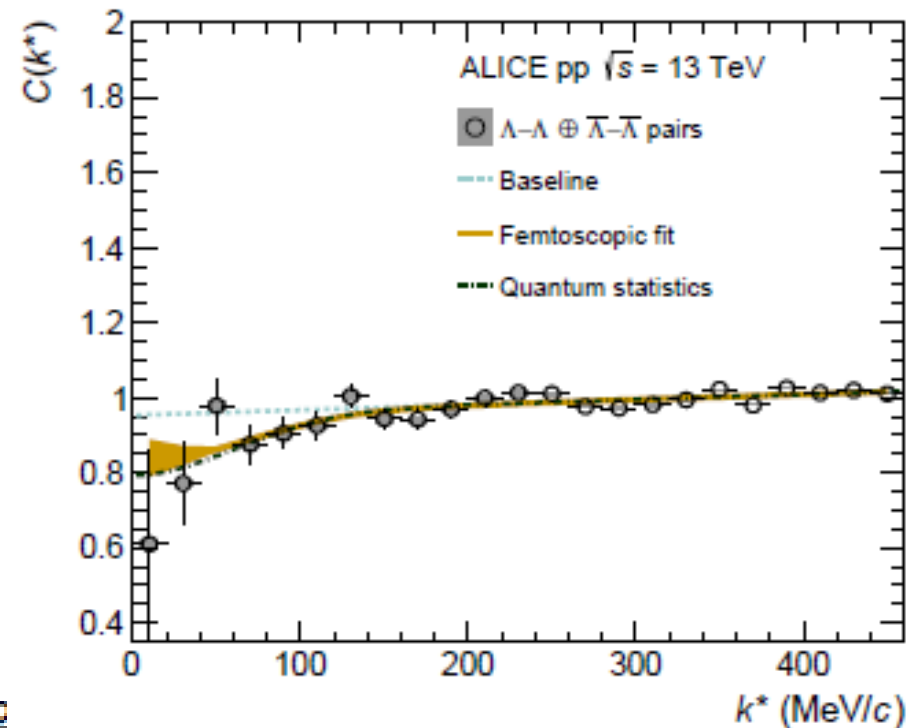
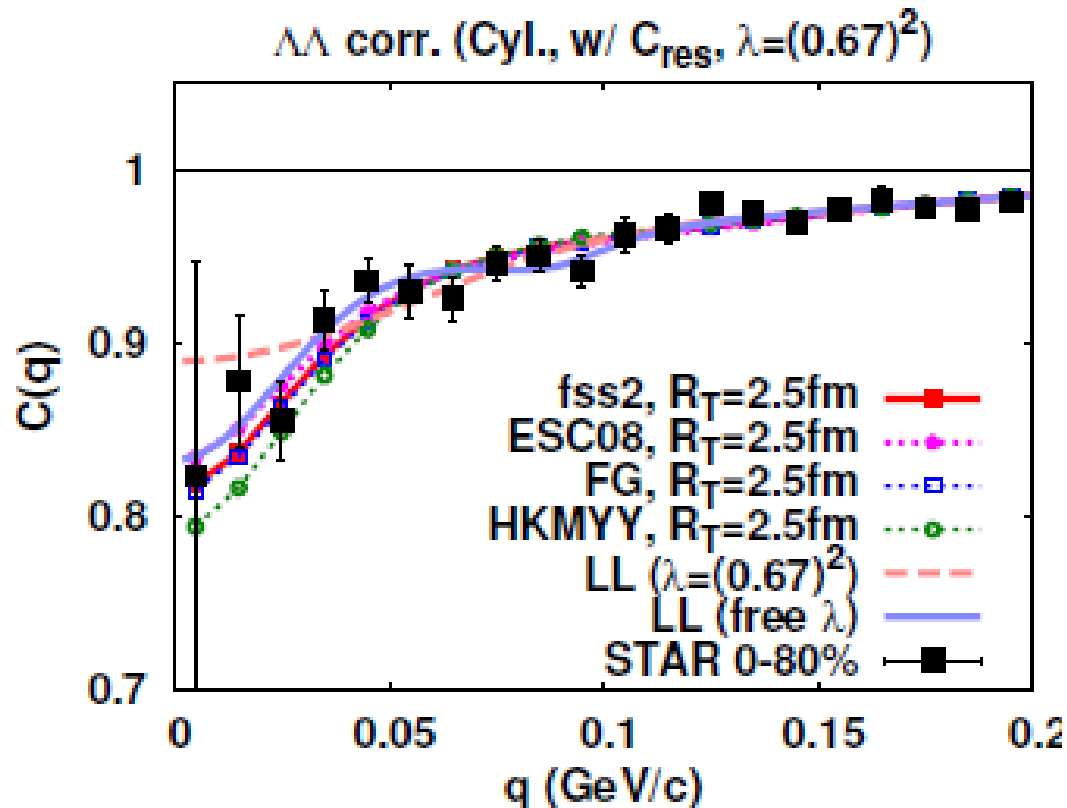
## ■ Neutron Star Matter EOS

- Hyperon Puzzle  
*Demorest et al. ('10), Antoniadis et al. ('13)*
- Cooling Puzzle ( $\Lambda\Lambda$  superfluidity)  
*T. Takatsuka, R. Tamagaki, PTP 112('04)37*



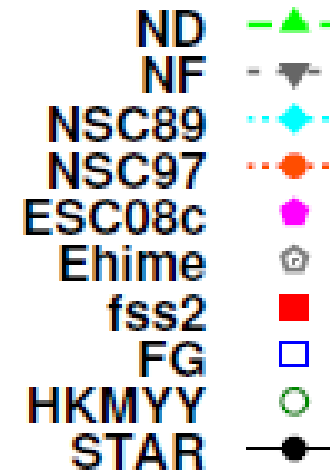
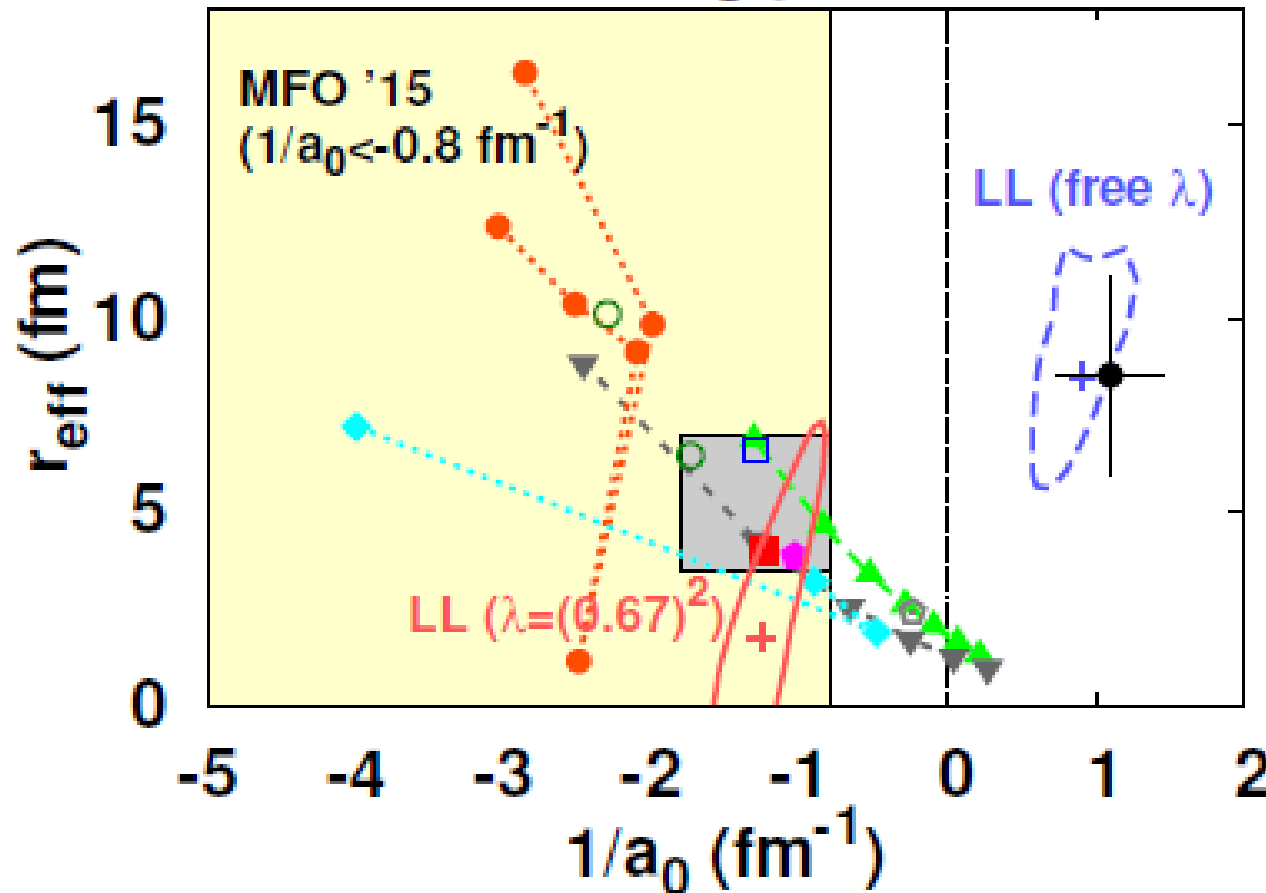
# $\Lambda\Lambda$ correlation at RHIC

- STAR collaboration at RHIC measured  $\Lambda\Lambda$  correlation !  
*Adamczyk et al. (STAR Collaboration), PRL 114 ('15) 022301.*
- Theoretical Analysis well explains the data  
*K.Morita et al., T.Furumoto, AO, PRC91('15)024916;*  
*AO, K.Morita, K.Miyahara, T.Hyodo, NPA954 ('16), 294.*
- New Data from ALICE  
*S. Achara+(ALICE), PRC99('19), 024001; arXiv:1905.07209*



# $\Lambda\Lambda$ interaction from $\Lambda\Lambda$ correlation

## $\Lambda\Lambda$ scattering parameters

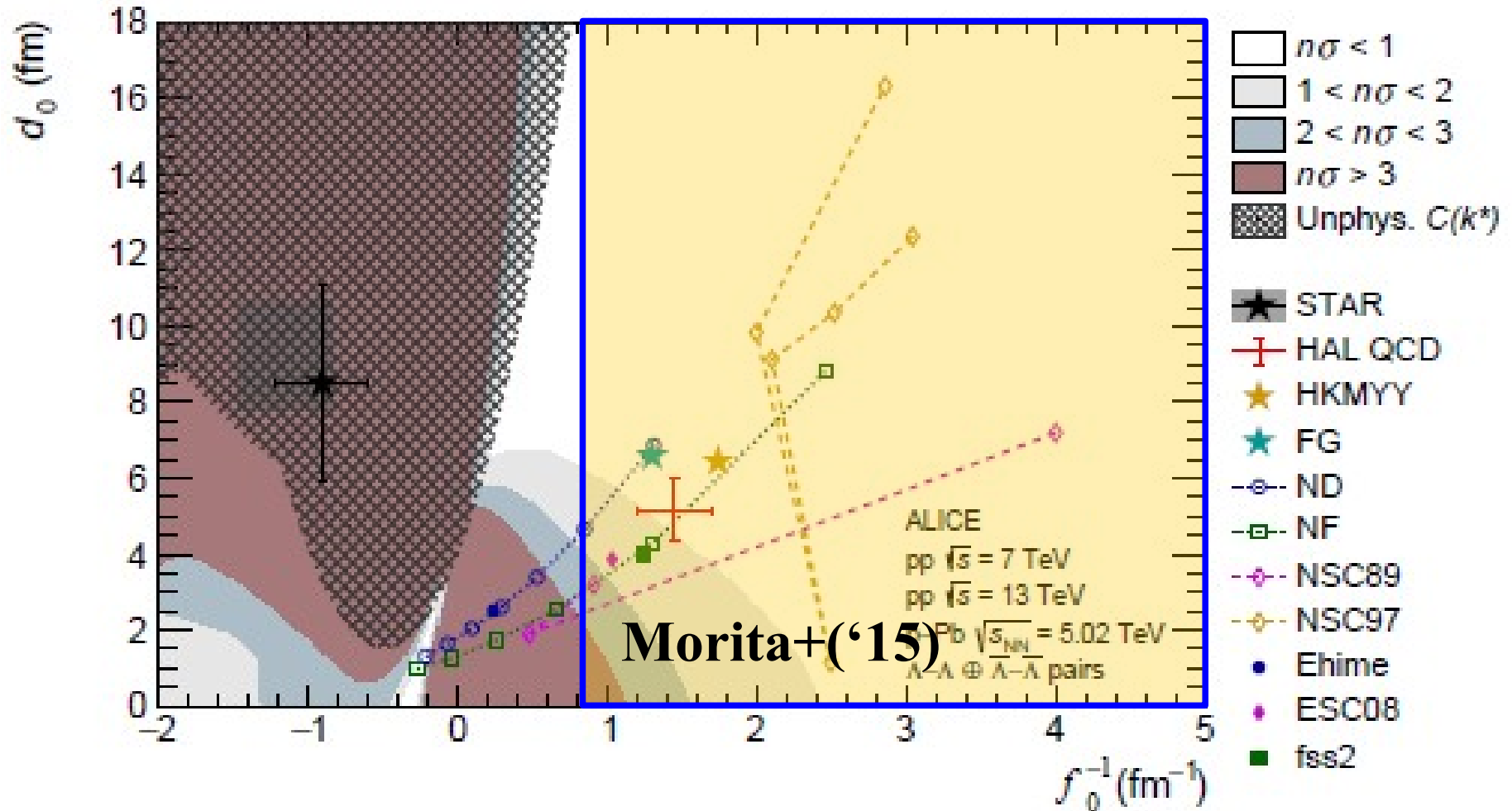


- Nijmegen potentials  
(ND, NF, NSC89, NSC97, ESC08)  
*Nagels+('77, '79), Maessen+('89),  
Rijken+('99, '10)*
- Ehime *Ueda et al. ('98)*
- Quark model interaction:  
fss2 *Fujiwara et al.('07)*
- Potential fitted to Nagara  
*Filikhin, Gal ('02) (FG),  
Hiyama et al. ('02, '10)(HKMYY)*

$$q \cot \delta = -1/a_0 + r_{\text{eff}} q^2/2 + \mathcal{O}(q^4)$$

# New Data from LHC-ALICE

ALICE (arXiv:1905.07209)



*Weakly attractive  $V_{\Lambda\Lambda}$*

*Large  $re_{\Lambda\Lambda}$   $\rightarrow$  Becomes repulsive at low relative density.*

# Relevance of $\Xi N$ interaction to physics

- H-particle: 6-quark state (uuddss) may be realized as a loosely bound state of  $\Xi N$  ( $I=0$ )

*K. Sasaki et al. (HAL QCD, '16, '17)*

- Repulsive  $\Xi N$  interaction ( $I=1$ ) may help to support  $2 M_{\odot}$  Neutron Star

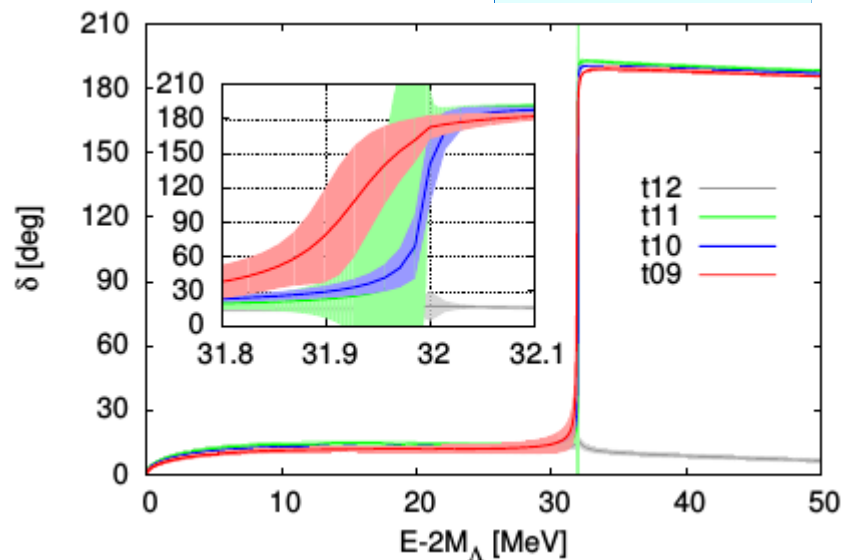
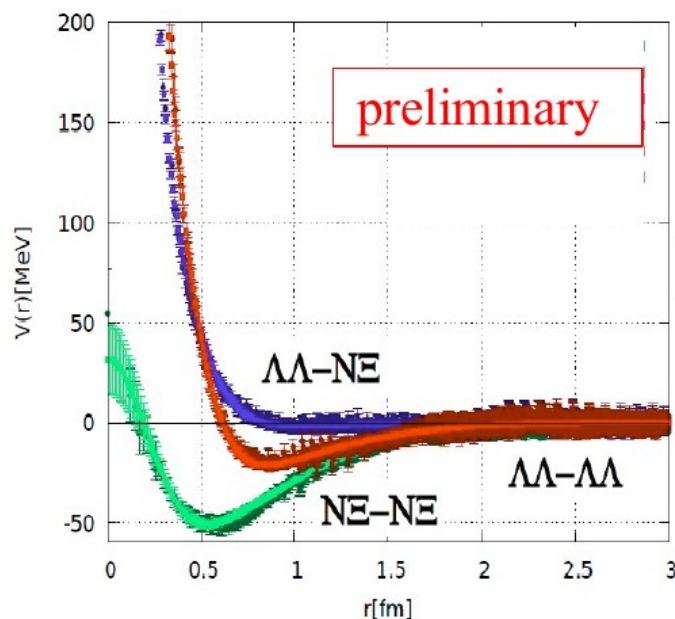
*Weissborn et al., NPA881 ('12) 62.*

$\Xi N$

E522 ('07)

$\Lambda\Lambda$

HAL ('16)



*K. Sasaki et al. (HAL QCD Collab.), EPJ Web Conf. 175 ('18) 05010.*

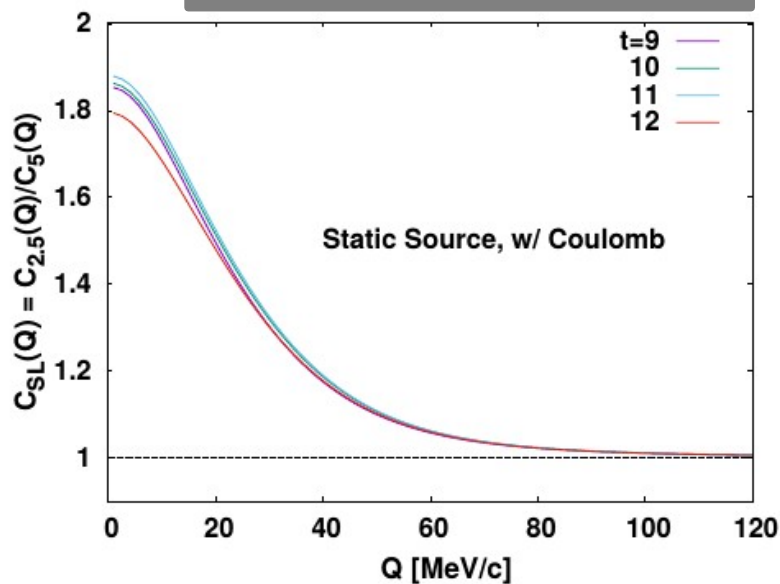


# $\Xi^-$ $p$ correlation

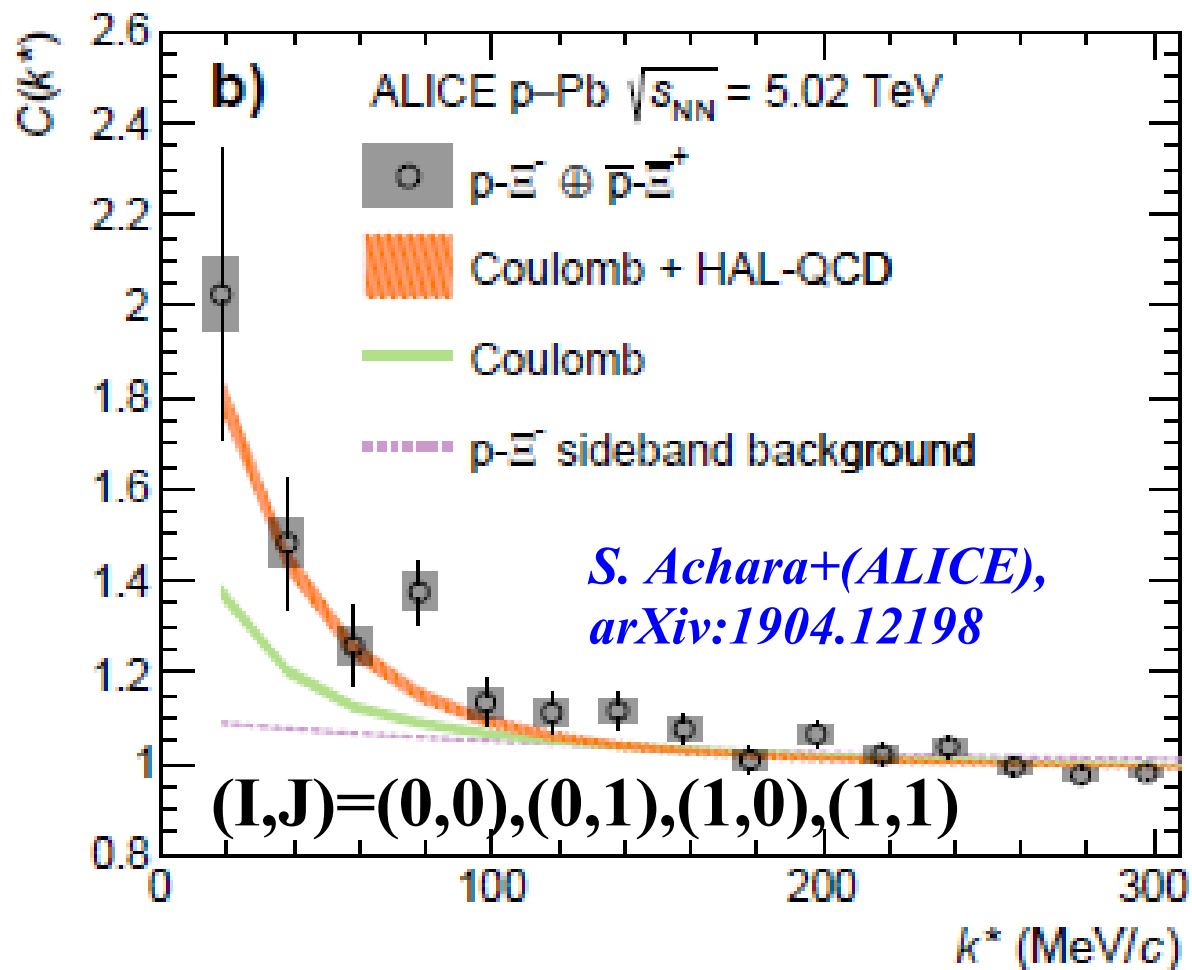
- Prediction of the correlation function by using  $\Xi N$  potential (HAL QCD Collab.) + Coulomb potential

$$|\psi|_{\text{spin av.}}^2 = \frac{1}{2} \sum_{I=0,1} \left[ \frac{1}{4} |\psi_I^{J=0}|^2 + \frac{3}{4} |\psi_I^{J=1}|^2 \right]$$

*HAL prediction is examined!*

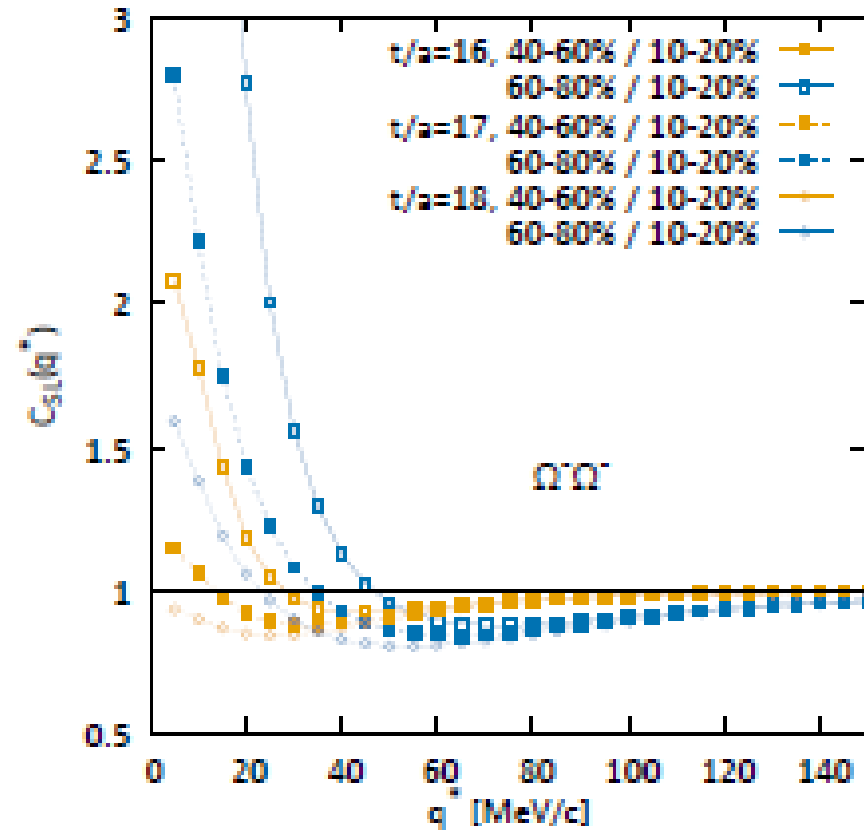
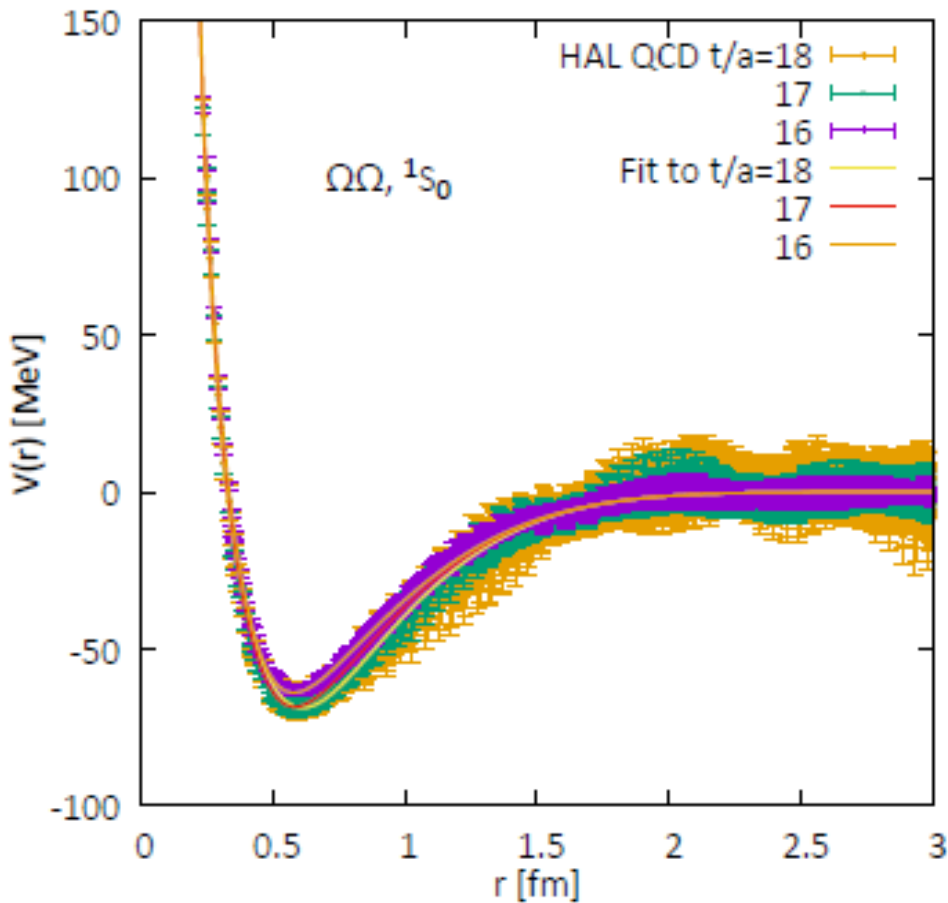


*T. Hatsuda, K. Morita, AO, K. Sasaki, NPA967('17), 856.*



# $\Omega\Omega$ correlation

$\Omega\Omega$  potential: S. Gongyo et al. (HAL QCD Collab),  
Phys. Rev. Lett. 120, 212001 (2017), 1709.00654.



K. Morita, S. Gongyo, T. Hatsuda, T. Hyodo, T. Iritani, AO, K. Sasaki, in prep.

# Summary

---

- High-energy nuclear collisions produce many and various hadrons in one event, and can be utilized as the hadron factory.
- Two-hadron correlation functions have been used to measure the source size. Once the properties of the source is known, we can utilize CF to get knowledge of hadron-hadron interaction.
- Large CF at small  $q$  implies large  $|a_0|/R$ , and the source size dependence of CF may shows the sign of  $a_0$ .
- It seems (to me) that there are at least two dibaryon states.
  - ALICE and STAR data strongly suggest the existence of  $S = -3$  dibaryon as a bound state of  $N\Omega$ .
  - ALICE data of  $\Xi^- p$  implies large  $|a_0|/R$ , and the existence of some kind of pole (b.s., res., virtual) around the threshold.
  - pp, pA and AA collision data will be helpful.

# To do (or Can do)

- Coupled channel effects → Talk by Heidenbauer, Kamiya
- $\Xi^-$  p correlation with updated HAL QCD potential (K. Sasaki et al.) → Sasaki's talk
- $\Lambda$ p correlation with various potentials ( $\chi$ EFT, Nijmegen, fss2, lattice) → Talk by Heidenbauer, Rijken
- $K^-$  p correlation with amplitude from chiral SU(3) dynamics (e.g. Ikeda, Hyodo, Weise) → Kamiya's talk
- J=1  $\Omega^-$  p potential → Hyodo's talk
- *Let us use deuteron !*  
 *$K^-$  d corr. (I=0 ampl.),  $\Lambda$ d corr. ( $^3_1$ H B.E., 3BF),  $\Xi^-$  d, ...*
- *Can we go to heavy-quarks ?*  
 *$c\tau(D) = 0.3 \text{ mm} \rightarrow \gamma c\tau(D) = c\tau(D) \cosh y \sim 15 \text{ cm} (y=7)$*   
*We may have enough D mesons at  $y=7$  in fixed target LHC.*  
*(N. Yamanaka)*

# Thank you for attention !

## Coauthors of arXiv:1908.05414

**K. Morita**



**S. Gongyo**



**T. Hatsuda**

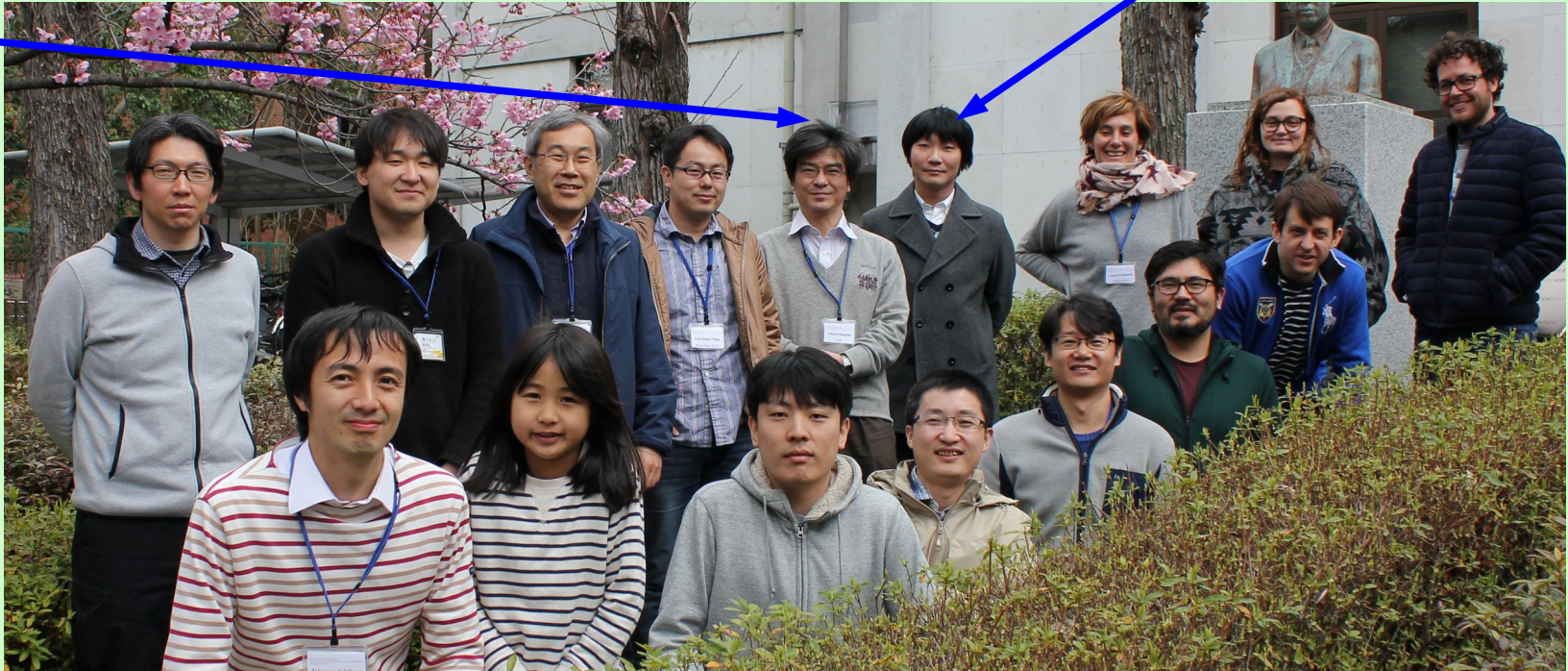


**T. Hyodo**



**Y. Kamiya**

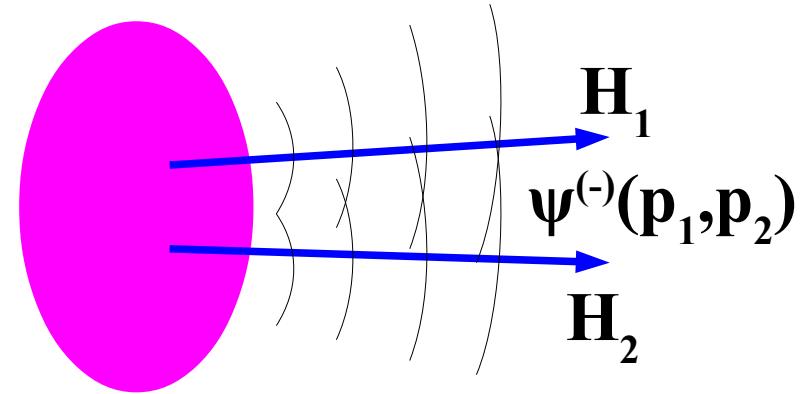
**AO**



# Hadron-Hadron Correlation in HIC

## ■ Hadron-Hadron Correlation Func. (Koonin-Pratt (KP) formula)

*S. E. Koonin, PLB 70 ('77) 43; S. Pratt, T. Csorgo and J. Zimanyi, PRC42 ('90) 2646; W. Bauer, C.-K. Gelbke, S. Pratt, Annu. Rev. Nucl. Part. Sci. 42 ('92)77; R. Lednicky, V. L. Lyuboshits, Sov. J. Nucl. Phys. 35 ('82) 770.*



$$C(\mathbf{q}) = \frac{E_1 E_2 dN_{12}/d\mathbf{p}_1 d\mathbf{p}_2}{(E_1 dN_1/d\mathbf{p}_1)(E_2 dN_2/d\mathbf{p}_2)} \simeq \int d\mathbf{r} \underbrace{S_{12}(\mathbf{r})}_{\text{Source fn.}} \underbrace{\left| \psi_{12}^{(-)}(\mathbf{r}, \mathbf{q}) \right|^2}_{\text{int.} \rightarrow \text{relative w.f.}}$$

$\mathbf{q}$ : rel. mom. (referred to also as  $\mathbf{k}^*$ )

## ■ Static sph. Gaussian source, Int. for s-wave, Identical fermions

$$C(q) \simeq 1 - \frac{1}{2} \exp(-4q^2 R^2) + \frac{1}{2} \int d^3 r \underbrace{S_{12}(r)}_{\text{Source}} \left[ \underbrace{|\psi_0(r)|^2}_{\text{s-wave w.f.}} - \underbrace{|j_0(qr)|^2}_{\text{free}} \right]$$

↖ **Fermion (Quant. Stat.)**

# Lednický-Lyuboshits (LL) model

## ■ Lednický-Lyuboshits analytic model

- **Asymp. w.f. + Eff. range corr. +  $\psi^{(-)} = [\psi^{(+)}]^*$**

$$\psi_0(r) \rightarrow \psi_{\text{asy}}(r) = \frac{e^{-i\delta}}{qr} \sin(qr + \delta) = \mathcal{S}^{-1} \left[ \frac{\sin qr}{qr} + f(q) \frac{e^{iqr}}{r} \right]$$

$$\begin{aligned} \Delta C_{\text{LL}}(q) &= \int dr S_{12}(r) (|\psi_{\text{asy}}(r)|^2 - |j_0(qr)|^2) \\ &= \frac{|f(q)|^2}{2R^2} F_3\left(\frac{r_{\text{eff}}}{R}\right) + \frac{2\text{Re}f(q)}{\sqrt{\pi}R} F_1(x) - \frac{\text{Im}f(q)}{R} F_2(x) \end{aligned}$$

( $x = 2qR$ ,  $R = \text{Gaussian size}$ ,  $F_1, F_2, F_3$  : Known functions)

## ■ Phase shifts

$$q \cot \delta = -\frac{1}{a_0} + \frac{1}{2} r_{\text{eff}} q^2 + \mathcal{O}(q^4) \rightarrow \delta \simeq -a_0 q + \mathcal{O}(q^3)$$

$$\sin(qr + \delta) \simeq \sin(q(r - a_0) + \dots) \quad \text{Node at } \mathbf{r} \sim \mathbf{a}_0$$

**for small  $q$**

# Binding Energy, Scattering Length, Effective Range ( $p\Omega$ )

- $m_\pi = 875$  MeV, B.E.  $\sim 19$  MeV

*F. Etminan et al. (HAL QCD Collab.), NPA928('14)89.*

- $m_\pi = 146$  MeV, B.E.  $\sim 2.2$  MeV

*T. Iritani et al. (HAL QCD Collab.), PLB 792('19)284.*

Spin-2 $N\Omega$ Potentials		$V_I$	$V_{II}$	$V_{III}$
without Coulomb	$E_B$ [MeV]	–	0.05	24.8
	$a_0$ [fm]	–1.0	23.1	1.60
	$r_{\text{eff}}$ [fm]	1.15	0.95	0.65
with Coulomb	$E_B$ [MeV]	–	6.3	26.9
	$a_0$ [fm]	–1.12	5.79	1.29
	$r_{\text{eff}}$ [fm]	1.16	0.96	0.65

$t/a$	$a_0$ [fm]	$r_{\text{eff}}$ [fm]	$E_B$ [MeV]
11	3.45	1.33	2.15
12	3.38	1.31	2.27
13	3.49	1.31	2.08
14	3.40	1.33	2.24

0.63

( $m_\pi = 146$  MeV,  $m_N = 955$  MeV and  $m_\Omega = 1712$  MeV). By using the same parameter set for  $t/a = 12$  in Table 1 with  $m_\pi = 146$  MeV kept fixed but with physical baryon masses ( $m_p = 938$  MeV and  $m_{\Omega^-} = 1672$  MeV), we find less binding than Eq. (10) as expected:  $B_{p\Omega^-} \simeq 2.18(32)$  MeV and  $\sqrt{\langle r^2 \rangle}_{p\Omega^-} \simeq 3.45(22)$  fm. On the other hand, if we additionally em-



# Correlation Function with Coulomb potential

- Correlation Function w/o Coulomb potential (spherical source)

$$C_i(\mathbf{q}) = 1 - \int d^3r S^i(\mathbf{r}) |j_0(q_i r)|^2 + \sum_j \omega_j \int d^3r S^j(\mathbf{r}) |\psi_{ij}(r)|^2$$

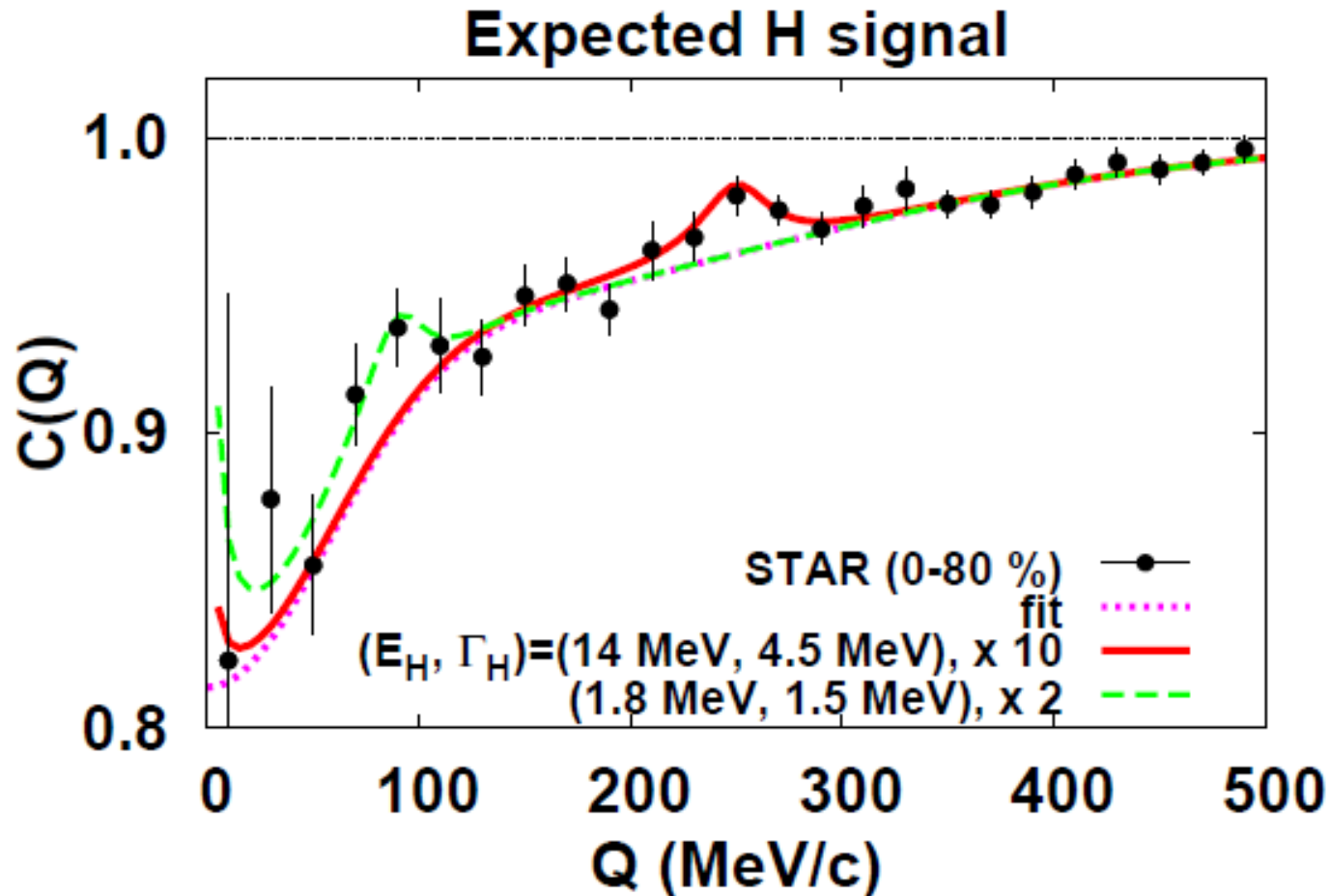
- Correlation Function with Coulomb potential (spherical source)

$$C_i(\mathbf{q}) = \int d^3r S^i(\mathbf{r}) \underbrace{|\psi^C(\mathbf{r}, \mathbf{q})|^2}_{\text{Coulomb wf}} - \int d^3r S^i(\mathbf{r}) \underbrace{|\psi_0^C(q_i r)|^2}_{\text{s-wave Coulomb wf}} + \sum_j \omega_j \int d^3r S^j(\mathbf{r}) \underbrace{|\psi_{ij}(r)|^2}_{\text{s-wave wf, } j \rightarrow i}$$

First two terms are large !

# Detecting Dibaryon State from Invariant Mass Spectrum ?

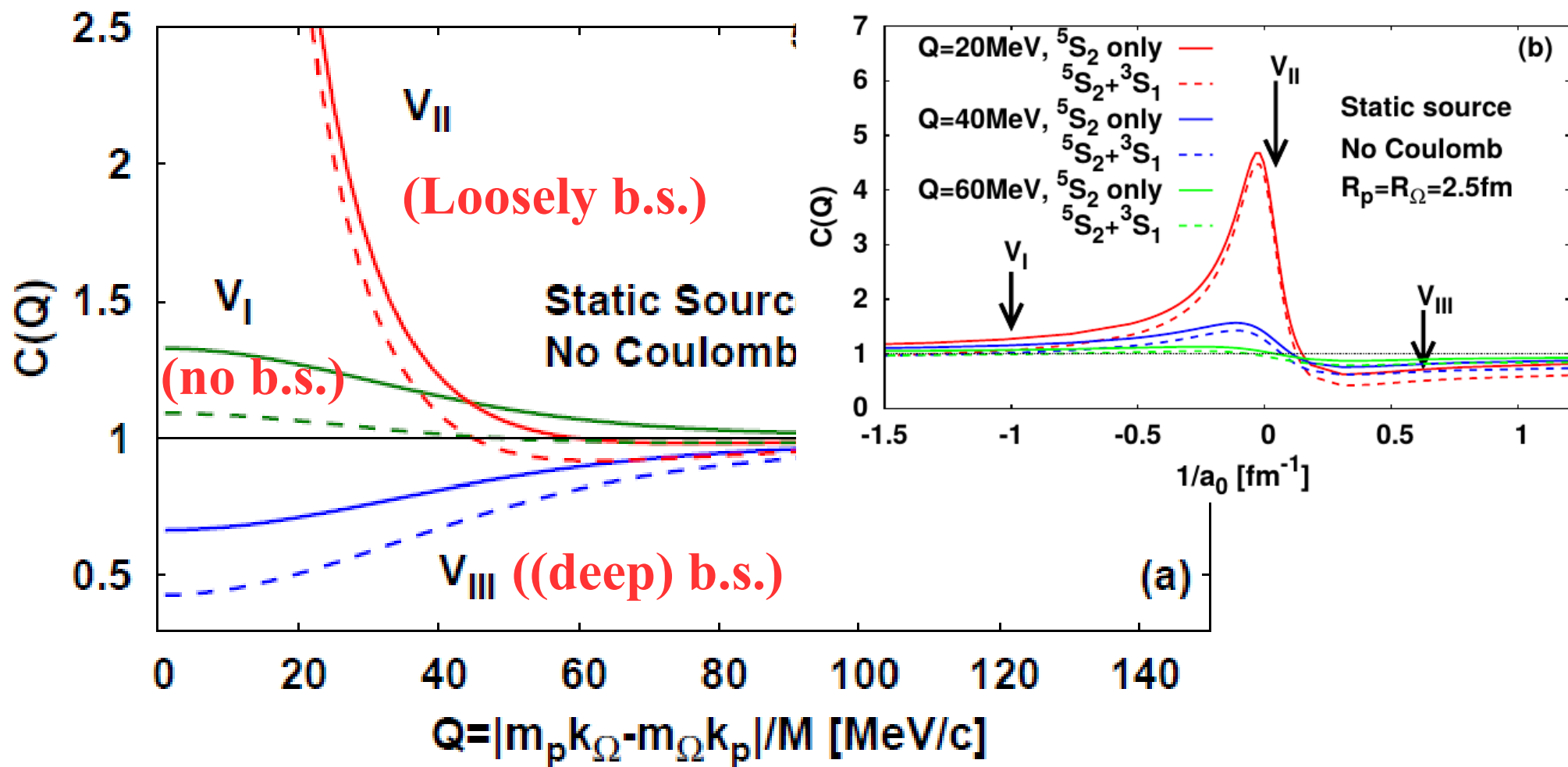
- Strong signal of dibaryon(s) in correlation function
  - How about invariant mass spectrum
  - Needs much more statistics



*Morita+ ('15)*

**J. Haidenbauer will make an objection ...**

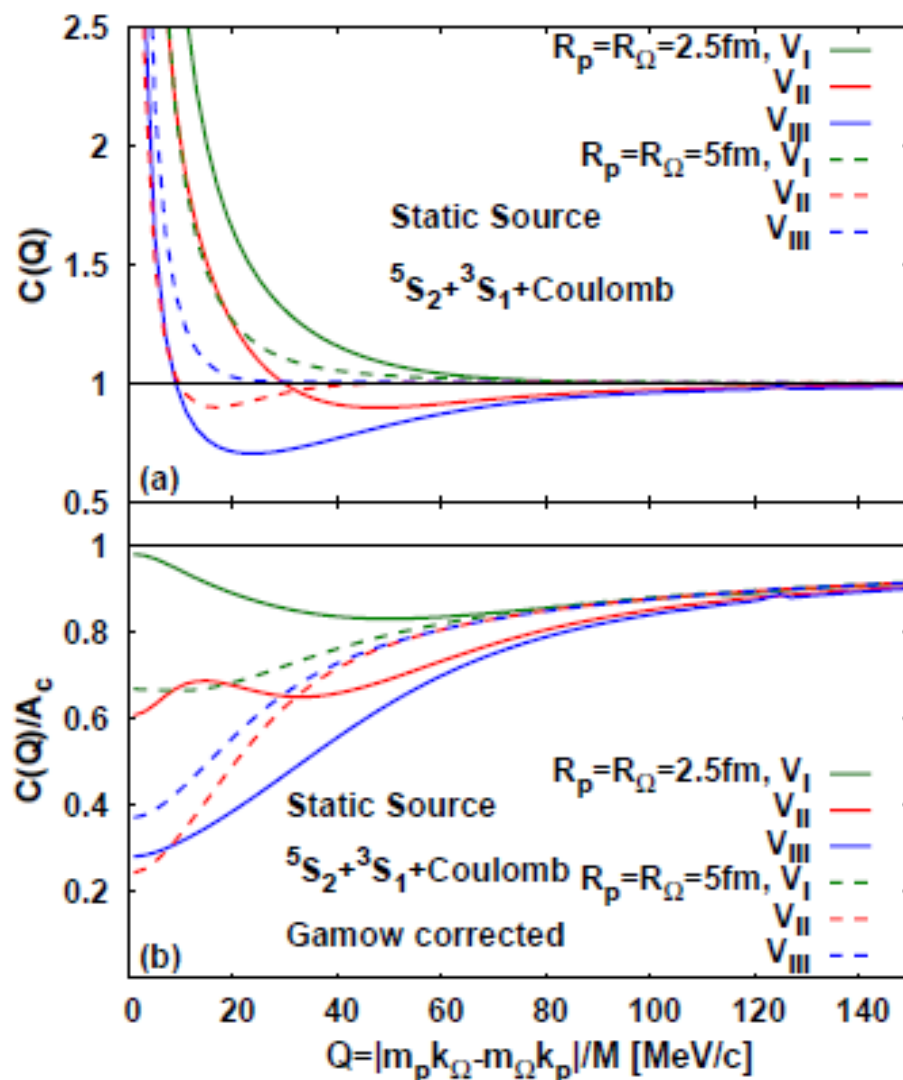
# $\Omega$ - $p$ correlation



(w/o Coulomb, Strong absorption at  $r < 2$  fm in  ${}^3S_1$  (decay to 8-8 in S-wave))

*K. Morita, AO, F. Etminan, T. Hatsuda, PRC94('16)031901(R) [arXiv:1605.06765 [hep-ph]]*

# $\Omega$ - $p$ correlation w/ Coulomb

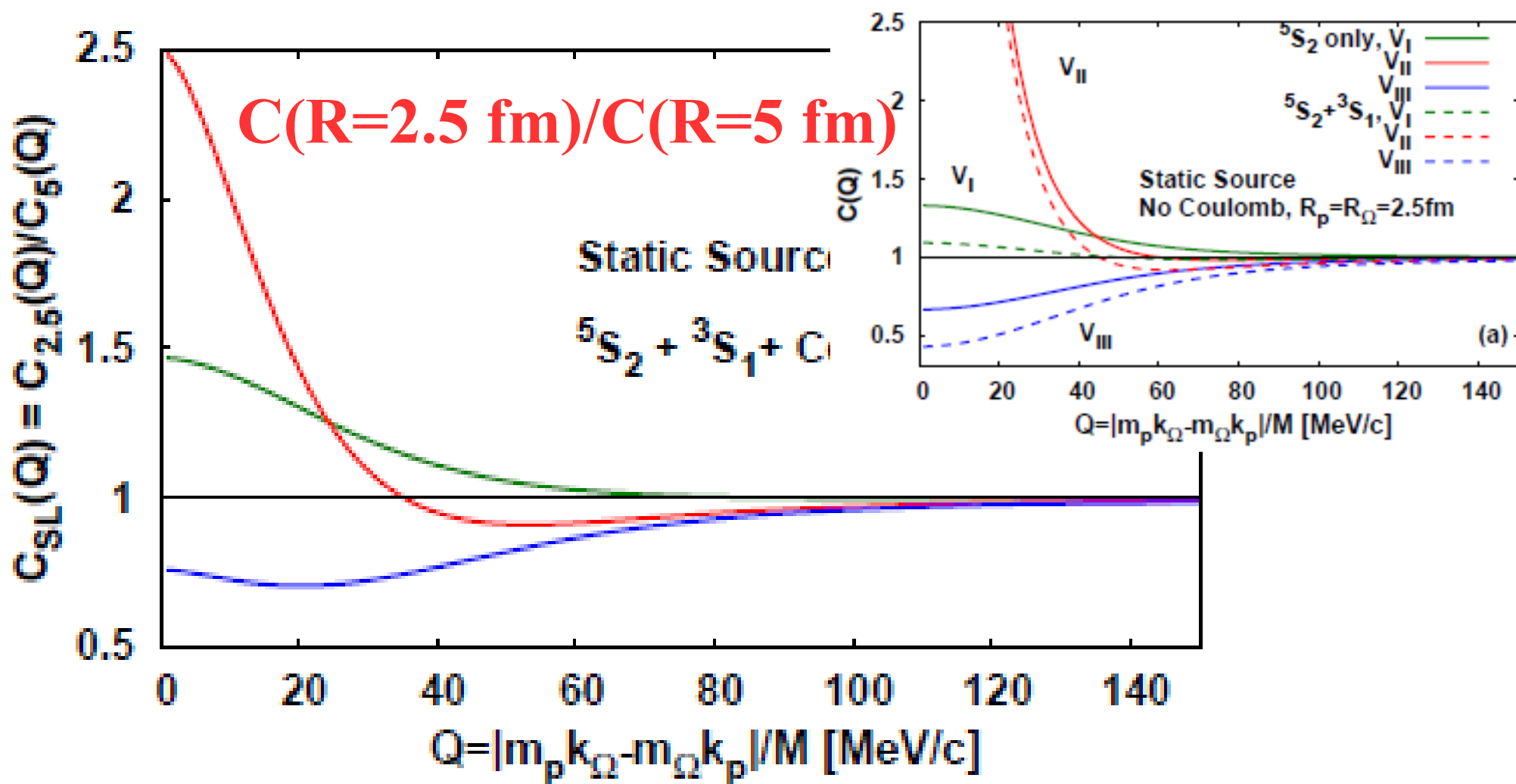


**With Coulomb**

**Coulomb + Gamow corr.**

*Coulomb potential washes out the features of  $V_I$ ,  $V_{II}$ ,  $V_{III}$ , and Gamow correction is not enough.*

# $\Omega$ - $p$ correlation: Small / Large Ratio



*By taking small ( $R=2.5 \text{ fm}$ ) / large ( $R=5 \text{ fm}$ ) ratio, we approximately see the corr. fn. w/o Coulomb !*

# $\Lambda\Lambda$ correlation in HIC

## ■ Merit of HIC to measure $\Lambda\Lambda$ correlation

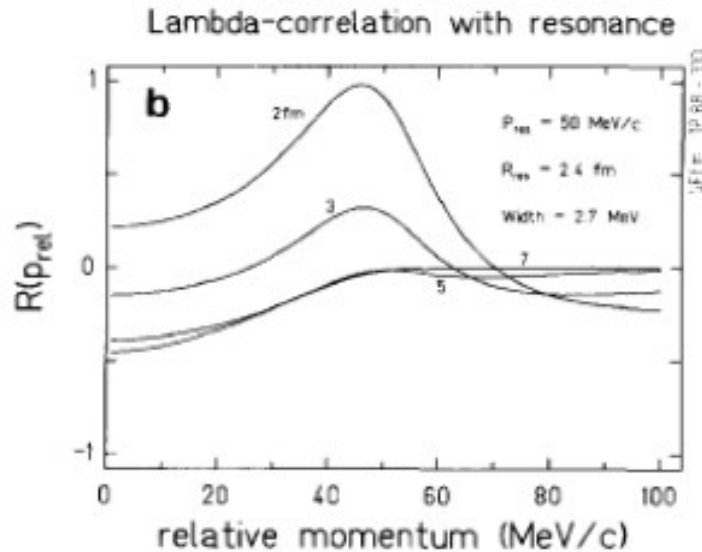
- Source is “Simple and Clean” !  
T,  $\mu$ , flow, size, ... are well-analyzed.

- Nearly Stat. prod.  
→ Many exotics will be produced.

*Schaffner-Bielich, Mattiello, Sorge ('00), Cho et al.(ExHIC Collab.) ('11)*

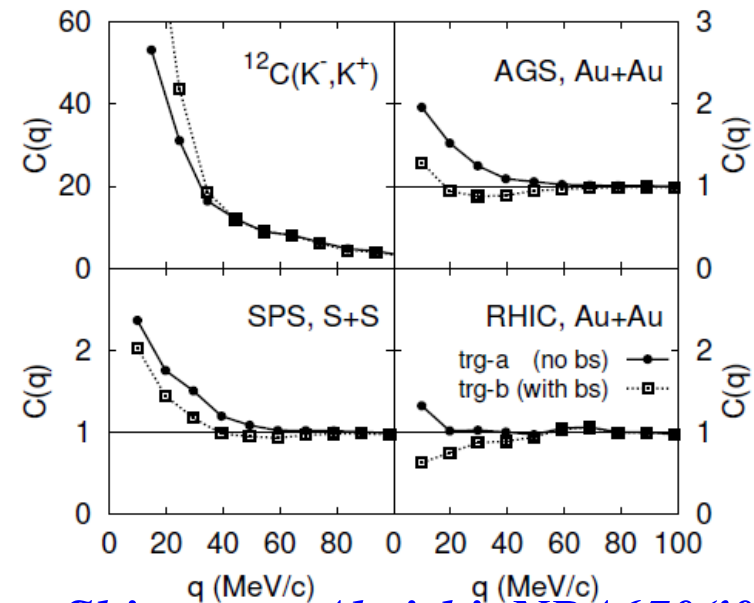
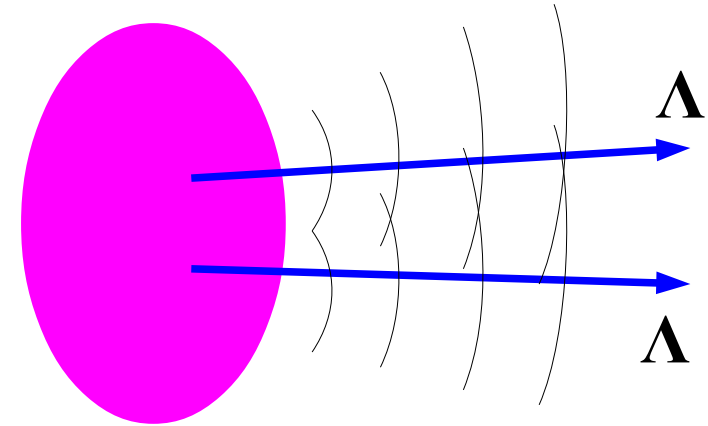
- Discovery of “H” and/or Constraint on  $\Lambda\Lambda$  int.

*Bound state exhaust the low q strength → suppressed  $C(q)$ .*

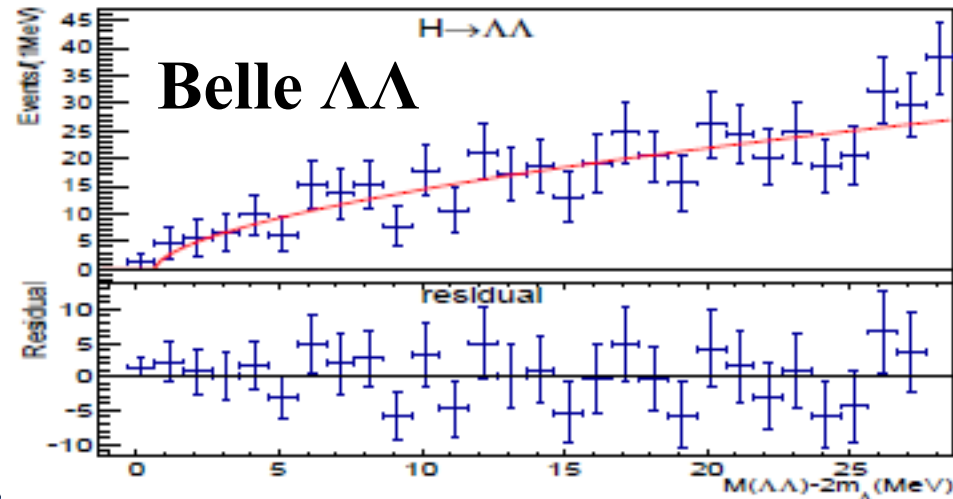
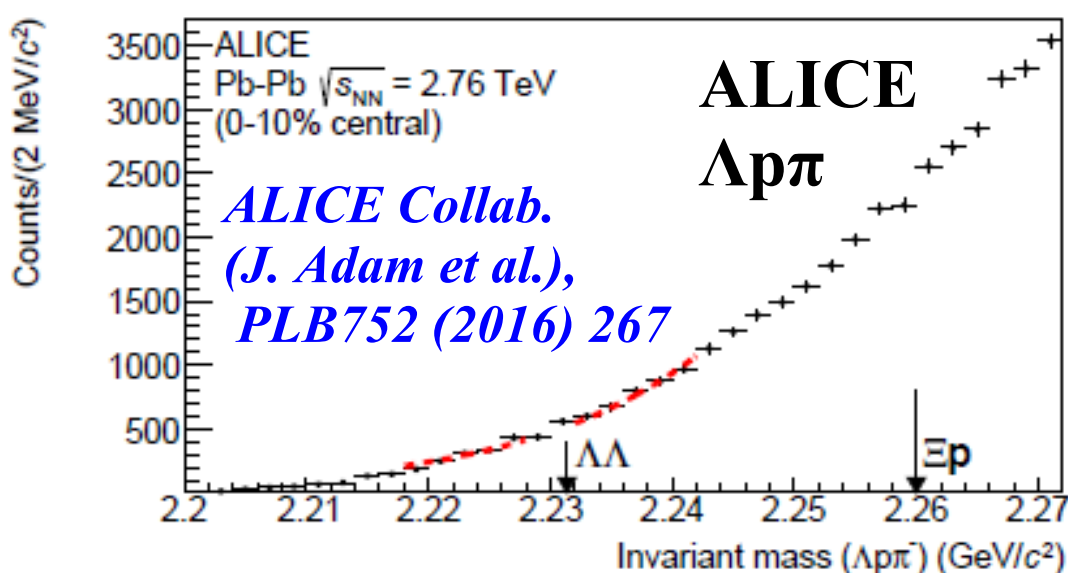


*C. Greiner, B. Muller, PLB219('89)199.*

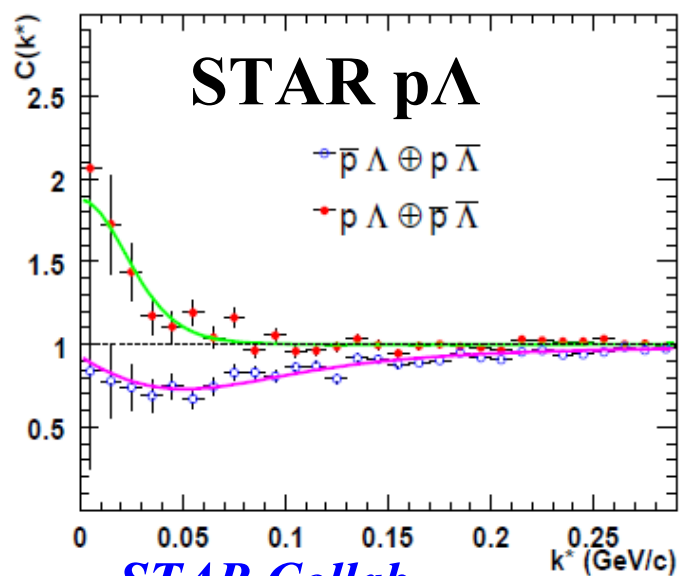
*AO, Hirata, Nara, Shinmura, Akaishi, NPA670('00)297c*



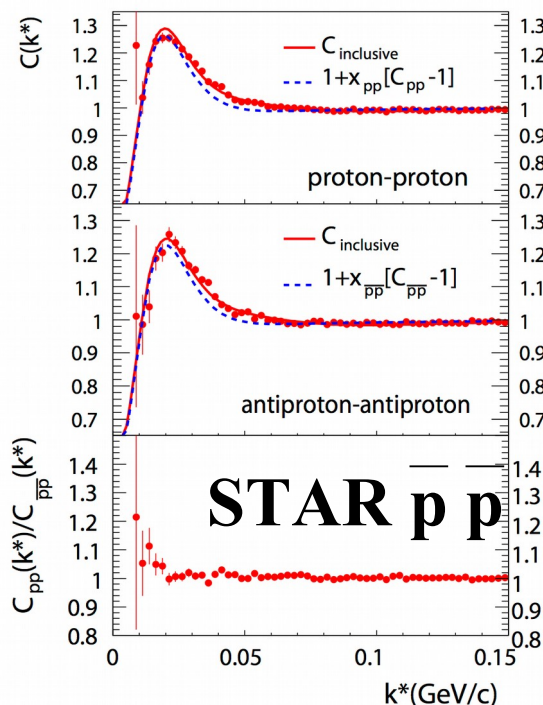
# $\Lambda\Lambda$ invariant mass / $BB$ correlation function (as of 2016)



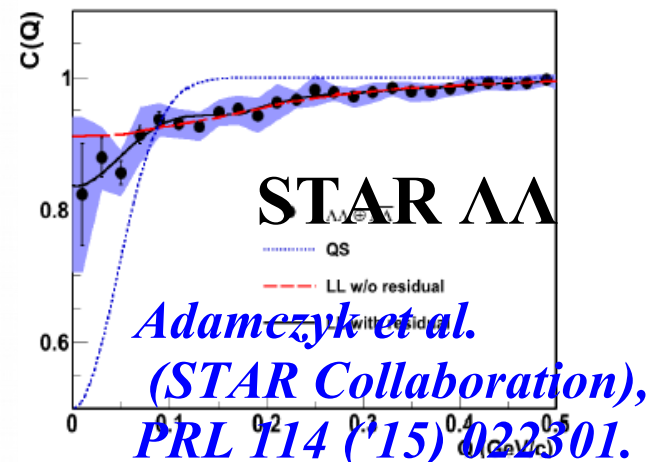
*Belle Collaboration (Kim, B.H. et al.),*  
*PRL110('13)222002.*



*STAR Collab.*  
*(J. Adams et al.),*  
*PRC74('06)064906.*



*STAR Collab.*  
*(J. Adams et al.)*  
*Nature 527('15)345*



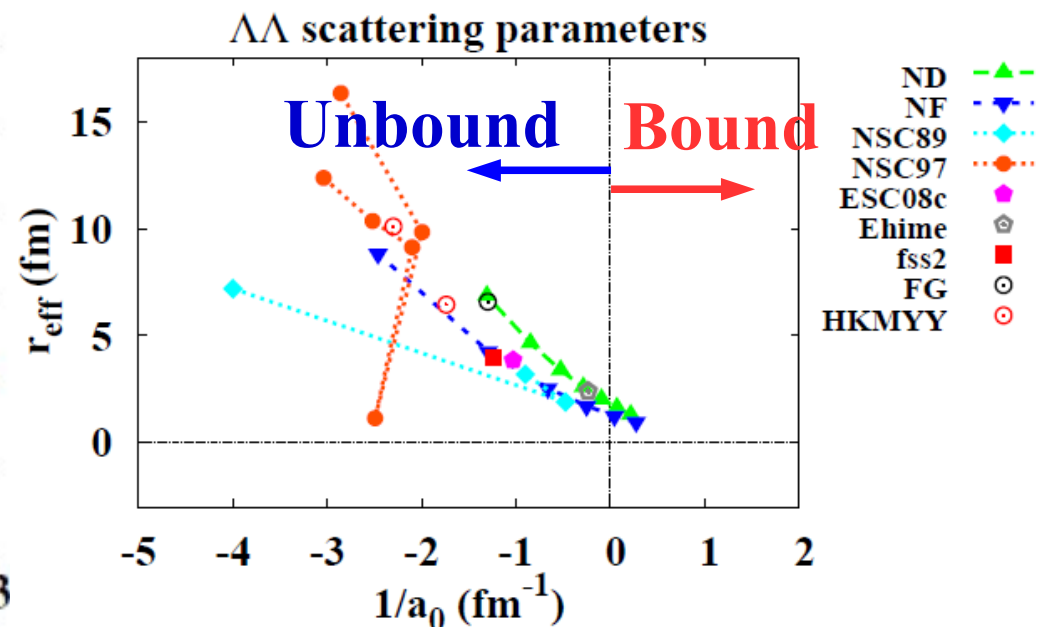
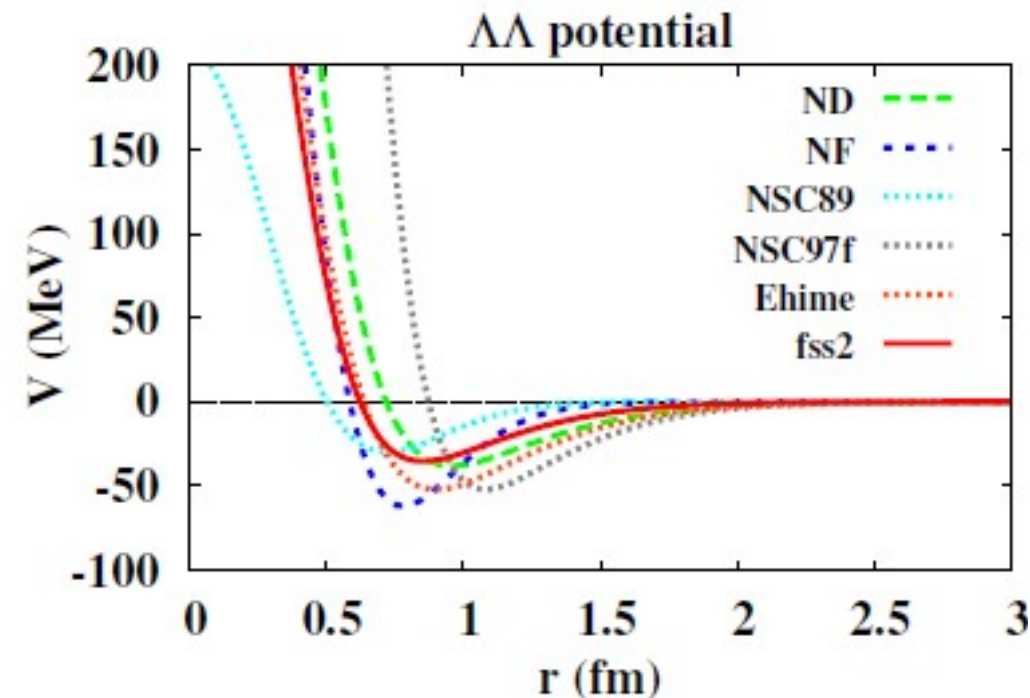
*Adamezyk et al.*  
*(STAR Collaboration),*  
*PRL 114 ('15) 022301.*

# $\Lambda\Lambda$ interaction

## Proposed $\Lambda\Lambda$ interactions

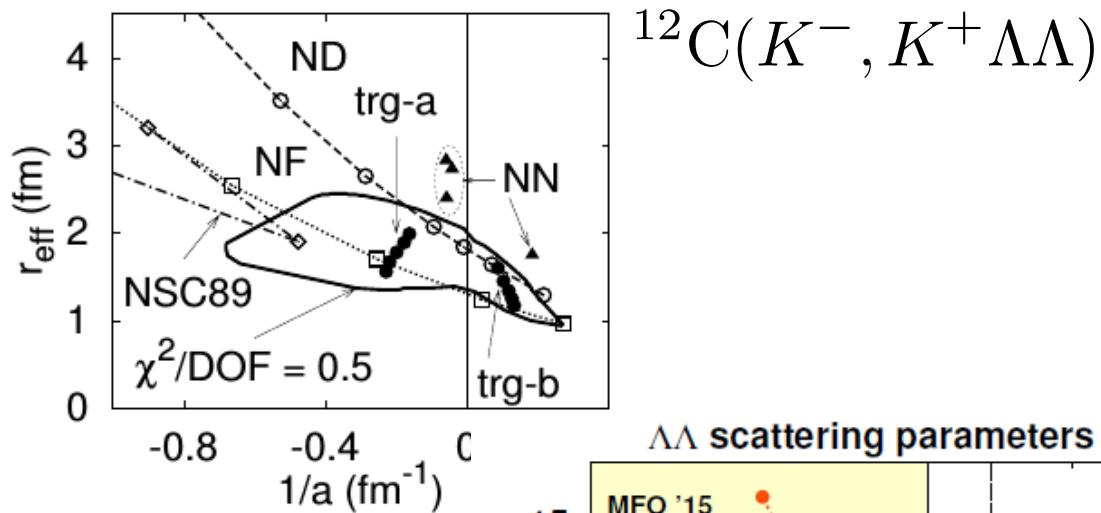
- **Meson Ex. models:** Nijmegen model D, F, Soft Core (89, 97), ESC08  
*Nagels, Rijken, de Swart ('77, '79), Maessen, Rijken, de Swart ('89), Rijken, Stoks, Yamamoto ('99); Rijken, Nagels, Yamamoto ('10).*
- **Quark cluster model interaction:** fss2  
*Fujiwara, Fujita, Kohno, Nakamoto, Suzuki ('00)*
- **Phenomenological model:** Ehime **T. Ueda et al. ('99).**

## Two (or three) range gaussian fit results are used in the analysis.

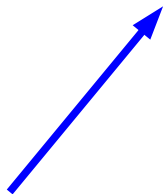




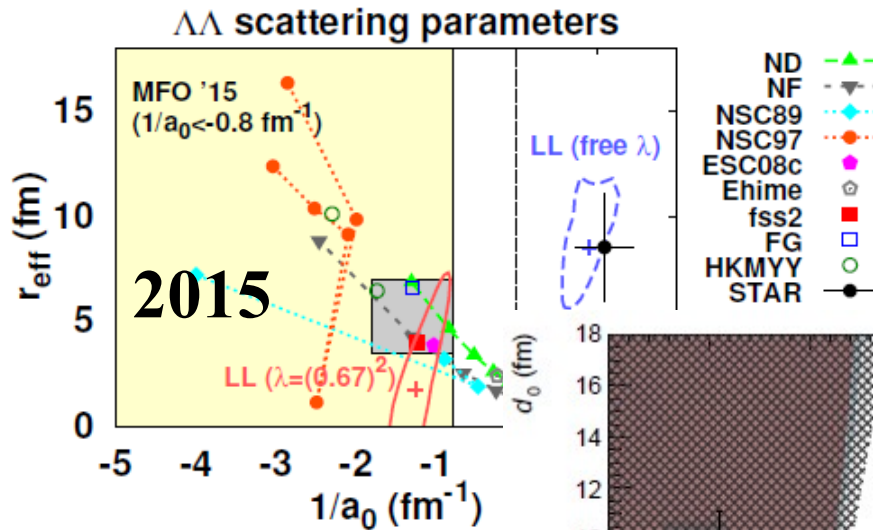
# Time dependence of $\Lambda\Lambda$ interaction



2000

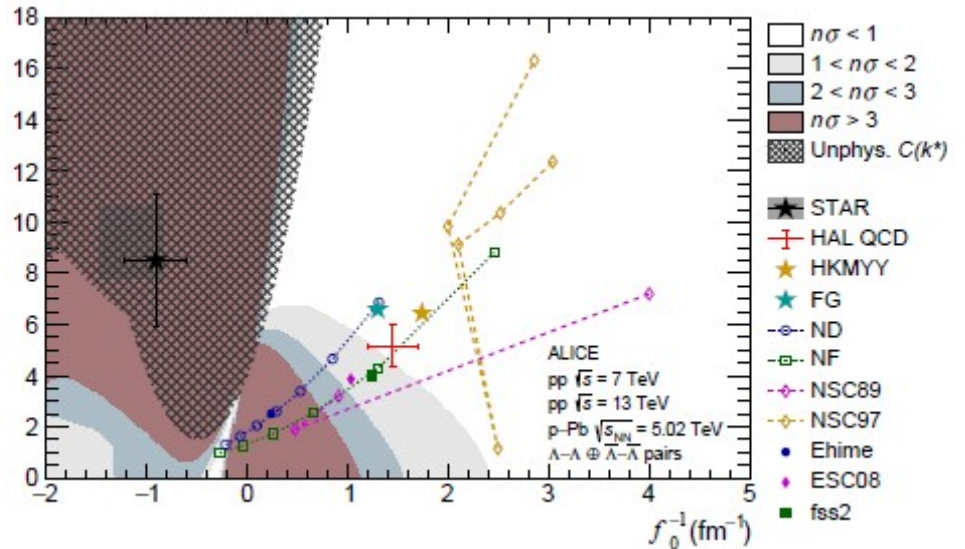


Nagara (2001)



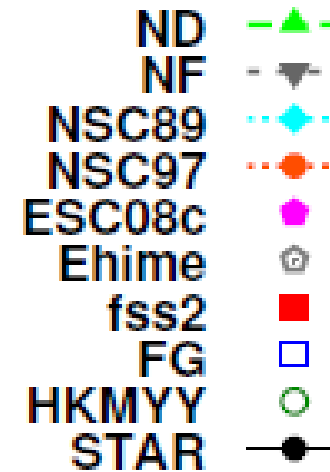
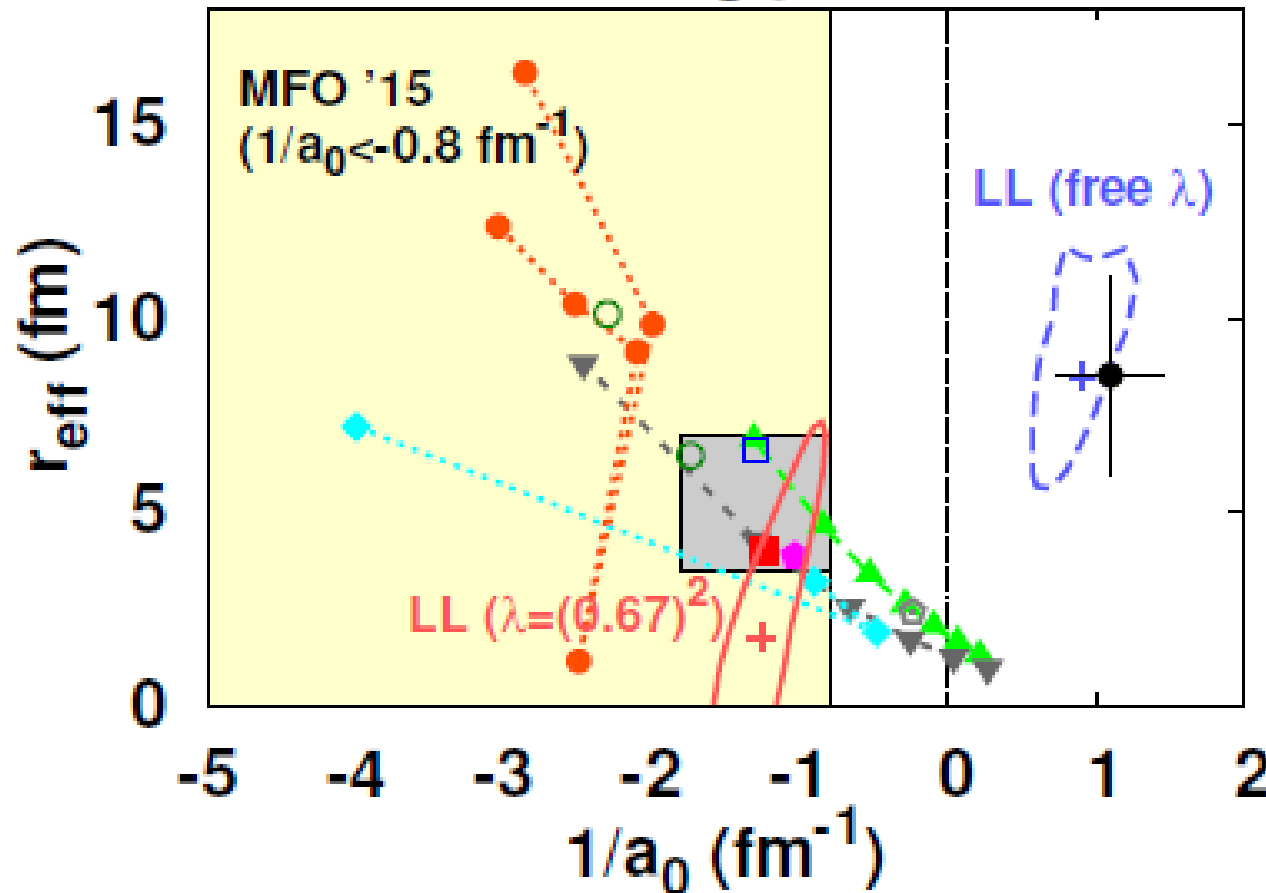
2015

2019



# $\Lambda\Lambda$ interaction from $\Lambda\Lambda$ correlation

## $\Lambda\Lambda$ scattering parameters



- Nijmegen potentials  
(ND, NF, NSC89, NSC97, ESC08)  
*Nagels+('77, '79), Maessen+('89),  
Rijken+('99, '10)*
- Ehime *Ueda et al. ('98)*
- Quark model interaction:  
fss2 *Fujiwara et al.('07)*
- Potential fitted to Nagara  
*Filikhin, Gal ('02) (FG),  
Hiyama et al. ('02, '10)(HKMYY)*

$$q \cot \delta = -1/a_0 + r_{\text{eff}} q^2/2 + \mathcal{O}(q^4)$$

*Positive  $a_0$  (STAR)  $\longleftrightarrow$  Negative  $a_0$  (MFO'15)*  
*Difference comes from the pair purity*

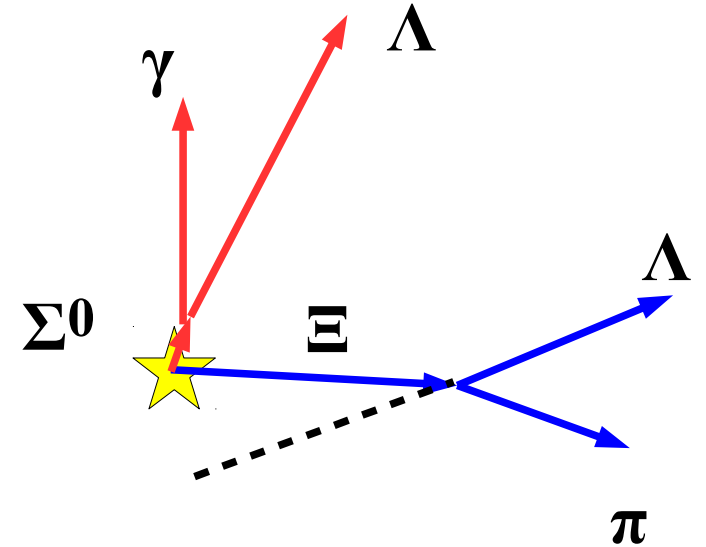
# Additional Source

## ■ Feed down effects

$$C_{\text{corr}}(Q) = 1 + \lambda(C_{\text{bare}}(Q) - 1)$$

$\lambda$  = Purity of  $\Lambda\Lambda$  pair

- Short-lived  $Y^*$   $\rightarrow$  mod. of source fn.
- $\Xi \rightarrow \Lambda\pi$  can be excluded ( $c\tau=8.71$  cm)
- $\Sigma^0 \rightarrow \Lambda\gamma$  is difficult to reject
- Data based purity  $\lambda=(0.67)^2$   
 $\Sigma^0/\Lambda=0.278$  (p+Be, 28.5 GeV/c) *Sullivan et al. ('87)*  
 $\Xi/\Lambda = 15\%$  (RHIC)



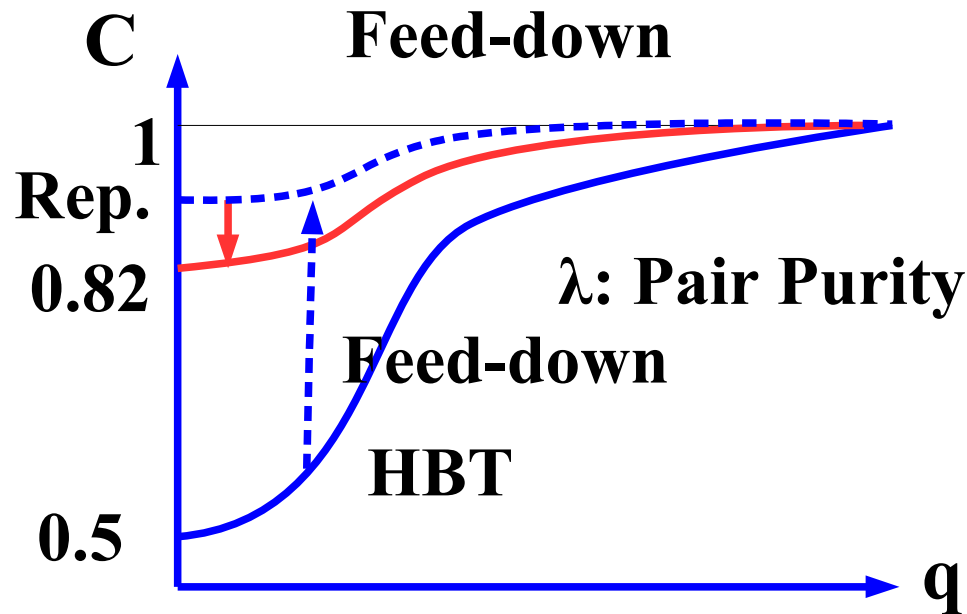
## ■ “Residual” source

- High-momentum tail  $\rightarrow R_{\text{res}} \sim 0.5$  fm (STAR collab.)

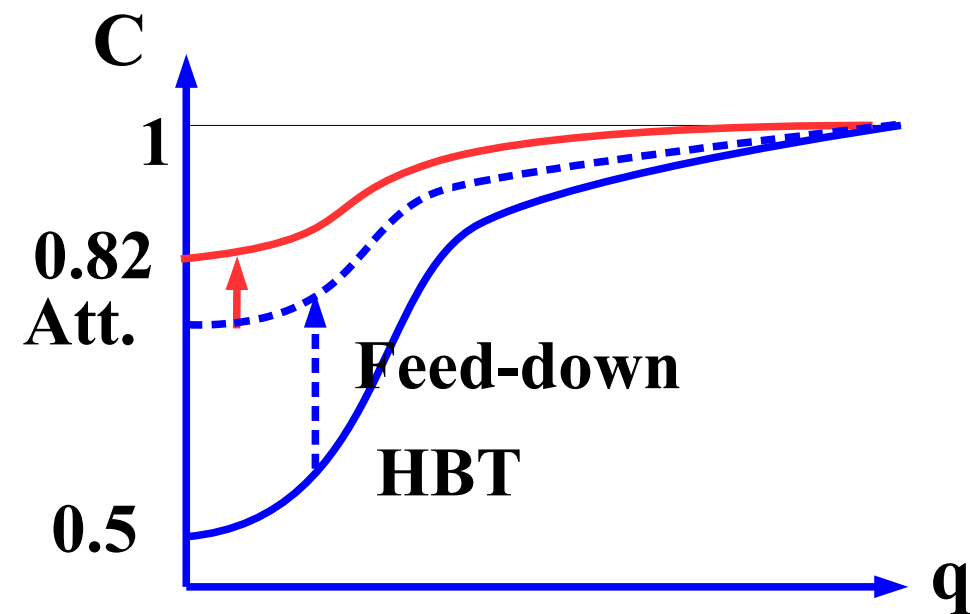
# Feed-Down Effects & Residual Source

- Correlation Fn. w/ Feed-down & Residual source effects.

$$C_{\text{corr}}(q) = 1 + \lambda(C_{\text{bare}}(q) - 1) + a_{\text{res}} \exp(-4r_{\text{res}}^2 q^2)$$



STAR:  
 $\lambda \sim 0.18$  (free para.)



Morita et al. (MFO15):  
 $\lambda \sim 0.45$

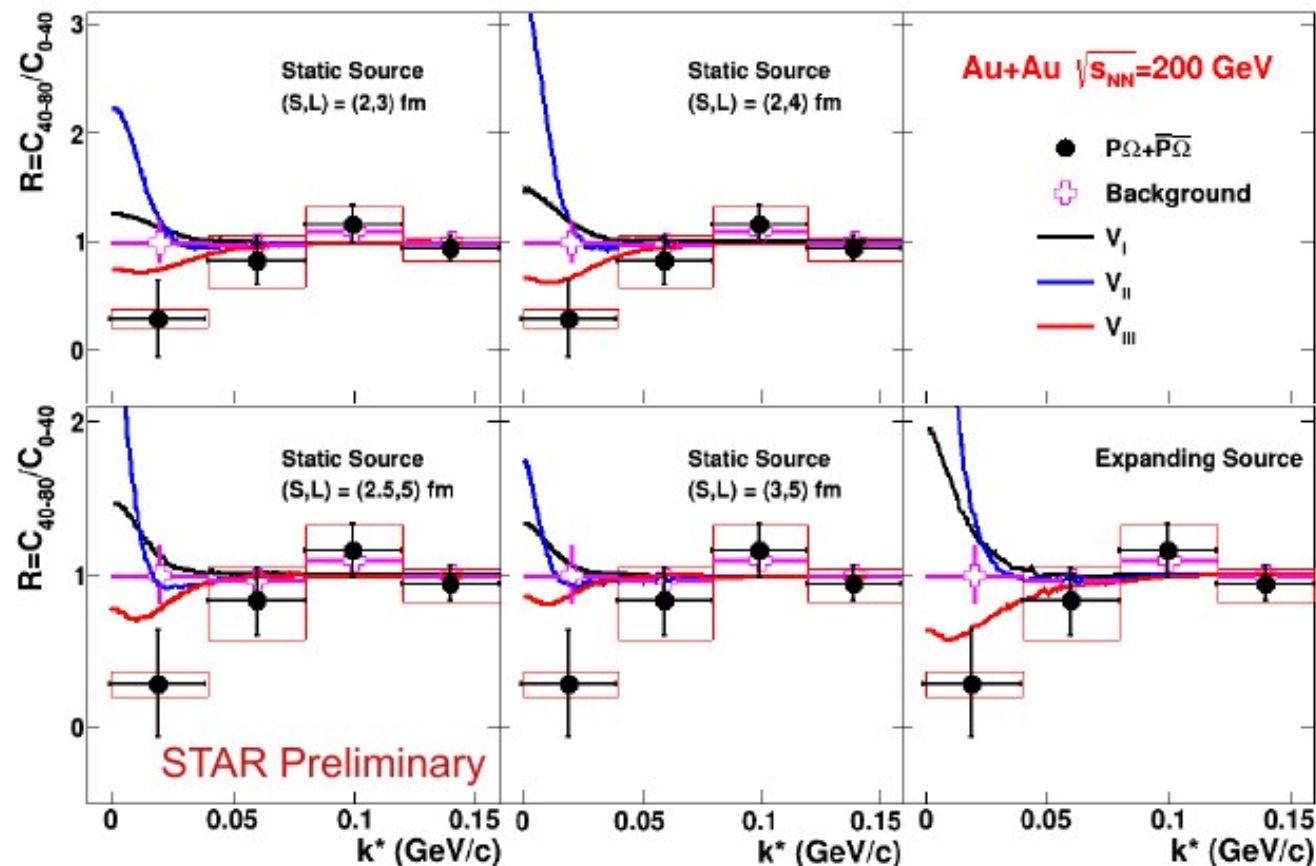
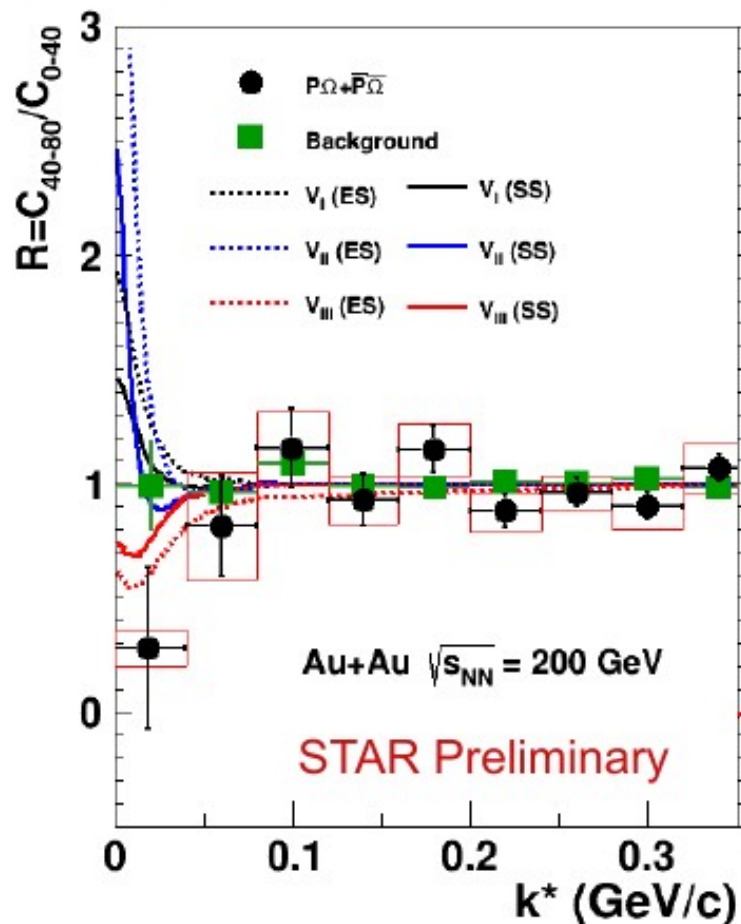
$\Sigma^0/\Lambda=0.278$  (p+Be, 28.5 GeV/c)  
*Sullivan et al. ('87)*  
 $\Xi/\Lambda = 15\%$  (RHIC)

*AO, Morita, Mihayara, Hyodo ('16)*

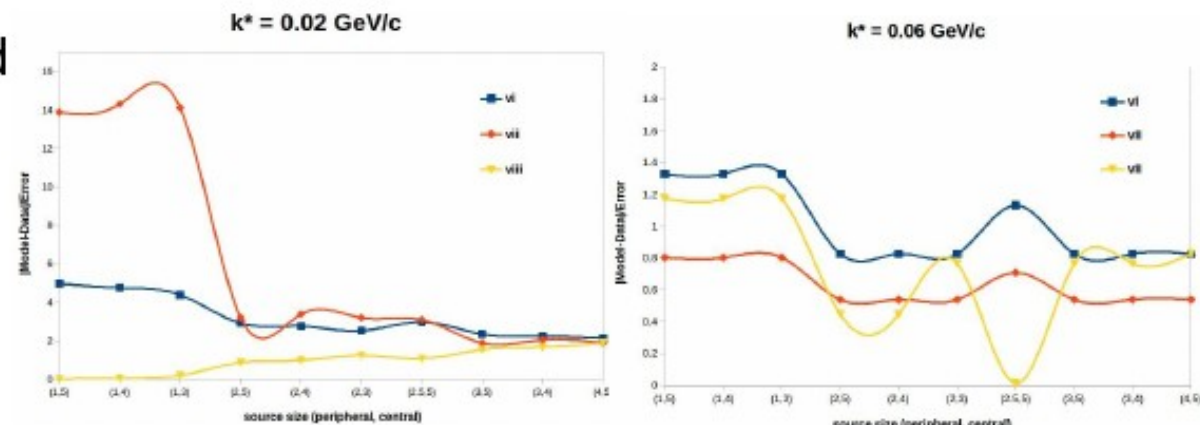
*Pair purity ( $\lambda$ ) should be determined experimentally !  
 Puzzle: Residual source*



## Source Size Analysis on $p\Omega$ Correlation Function



The source parameters are derived from other PID and centrality dependence measurements at RHIC



# $(K^-, K^+)$ reaction

- $(K^-, K^+)$  reaction = doorway to produce  $S=-2$  systems

- $K^- p \rightarrow K^+ \Xi^- \rightarrow \Xi$  nuclei, stopped  $\Xi$  to  $\Lambda\Lambda$  nuclei, ...

- What is the origin of large yield at smaller  $K^+$  momentum ?

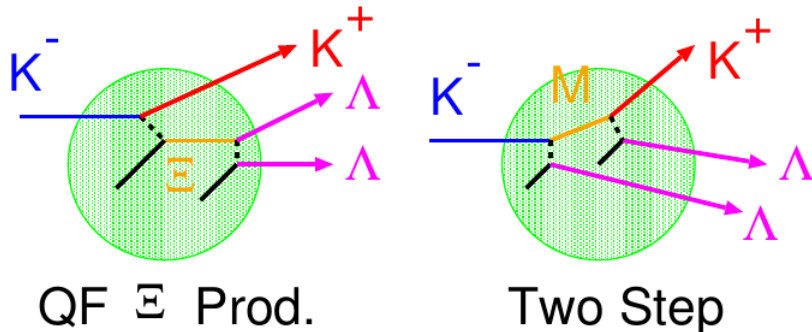
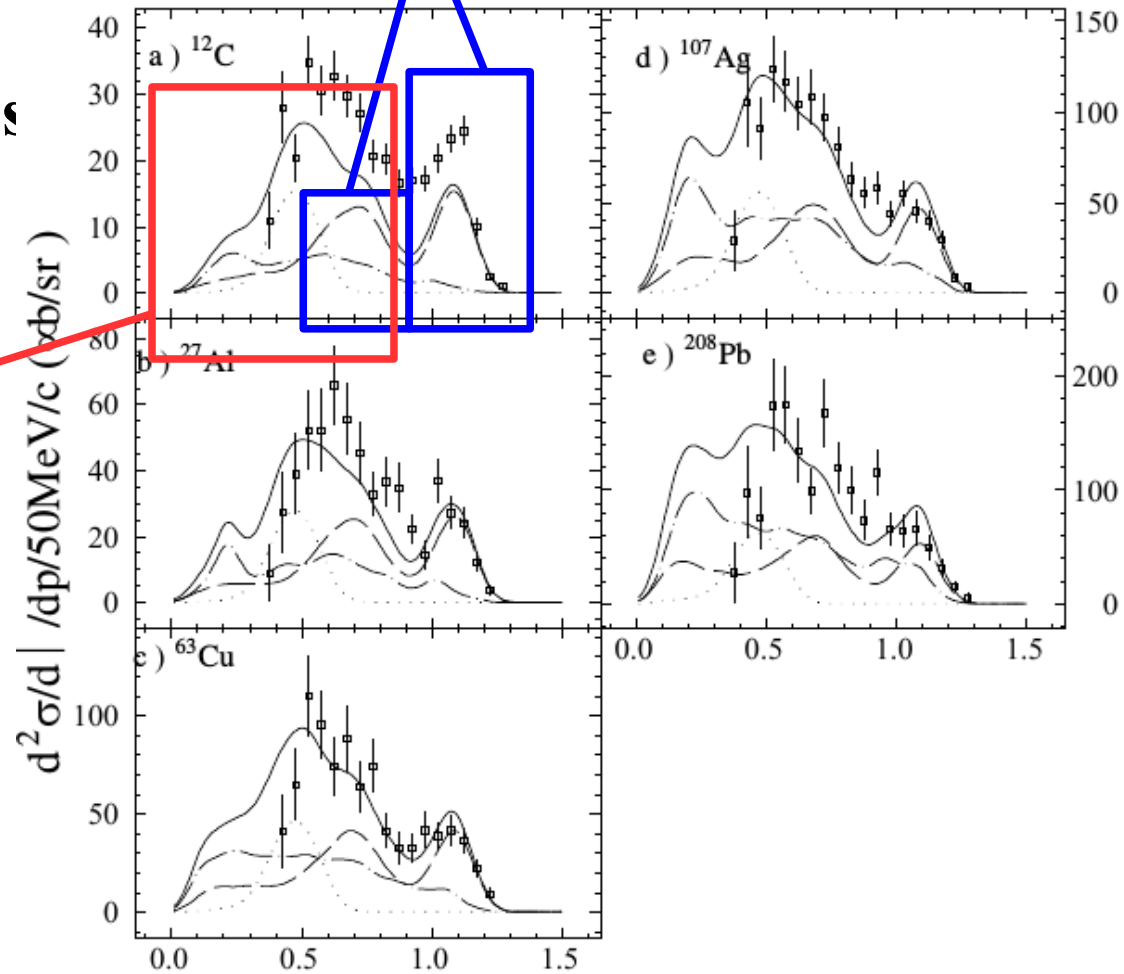
*T. Iijima et al., NPA546('92) 588.*

→ Various two step processes

*Y. Nara, AO, T. Harada, A. Engel, NPA614('97)433*

Quasi Free  $\Xi(\Xi^*)$

Multi-step



$K^+$  momentum (GeV/c)

# Two $\Lambda$ production in ( $K^-, K^+$ ) reactions

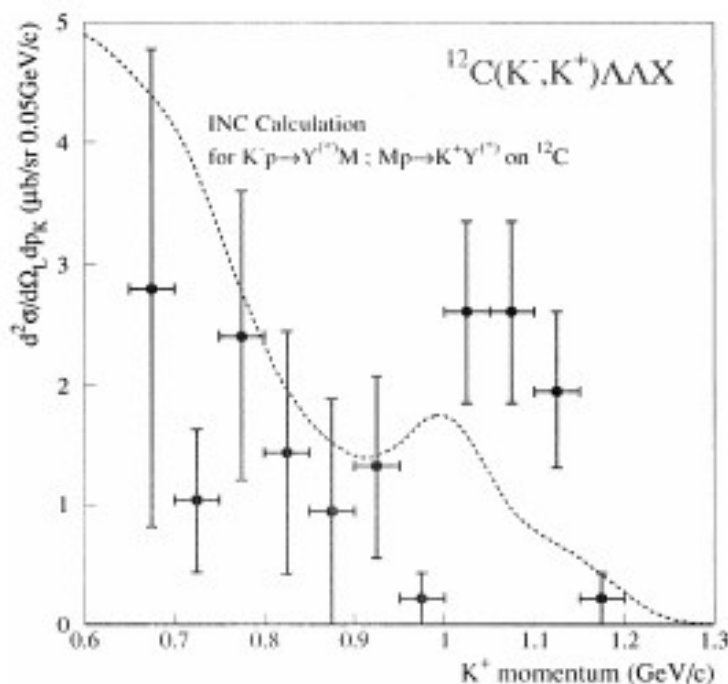
- Experimentalists really measured two  $\Lambda$  emission !

*J.K.Ahn et al. (KEK-PS E224 Collab.), PLB444('98)267.*

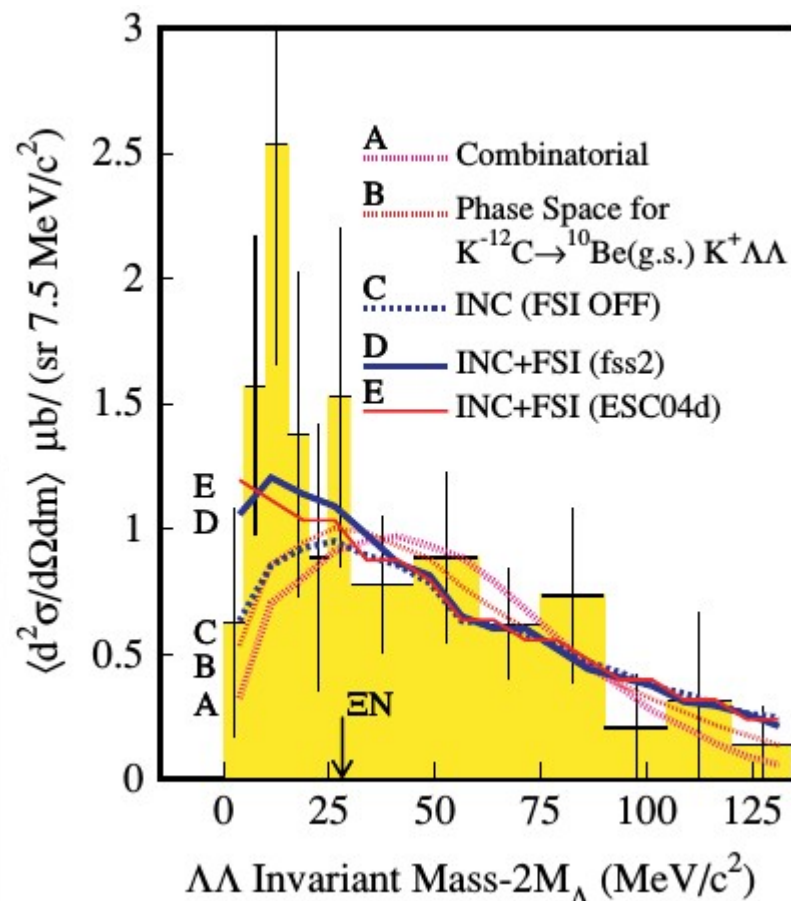
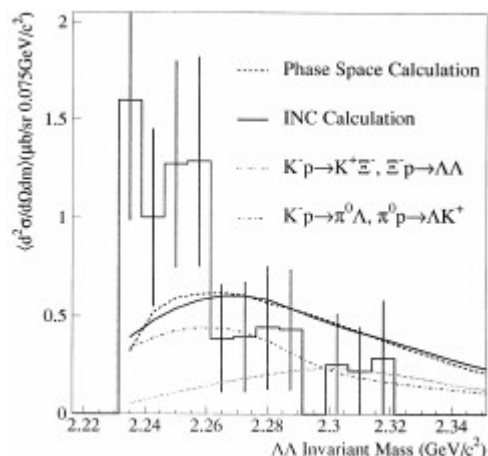
*C. J. Yoon et al. (KEK-PS E522 Collab. +AO), PRC75('07)022201*

- Invariant mass spectrum of  $\Lambda\Lambda$  is enhanced from our cascade calculation.

→ FSI enhancement ? or H particle ?



*Ahn et al. (E224)('98)*



*Yoon et al. (E522)('07)*

# From correlation to interaction

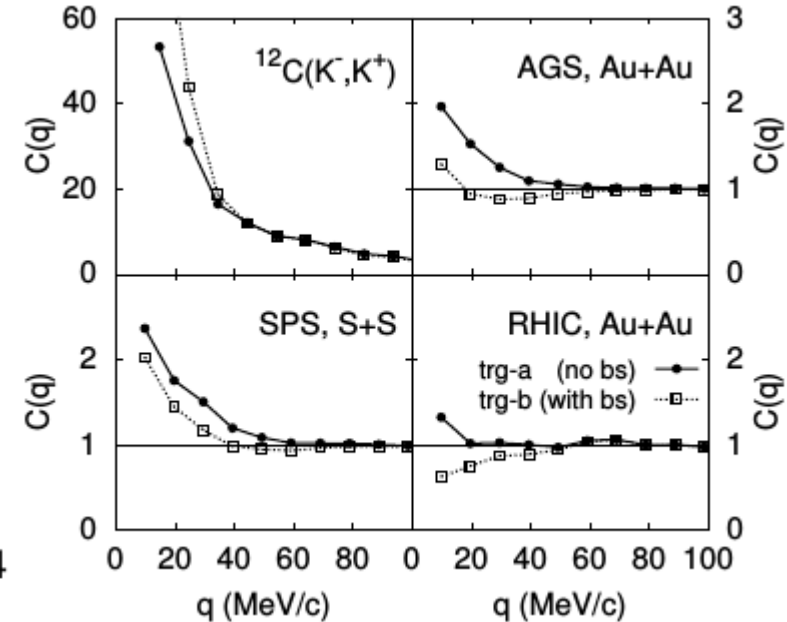
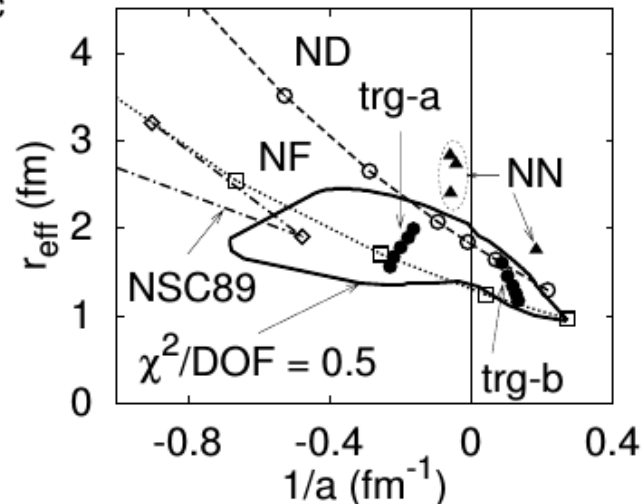
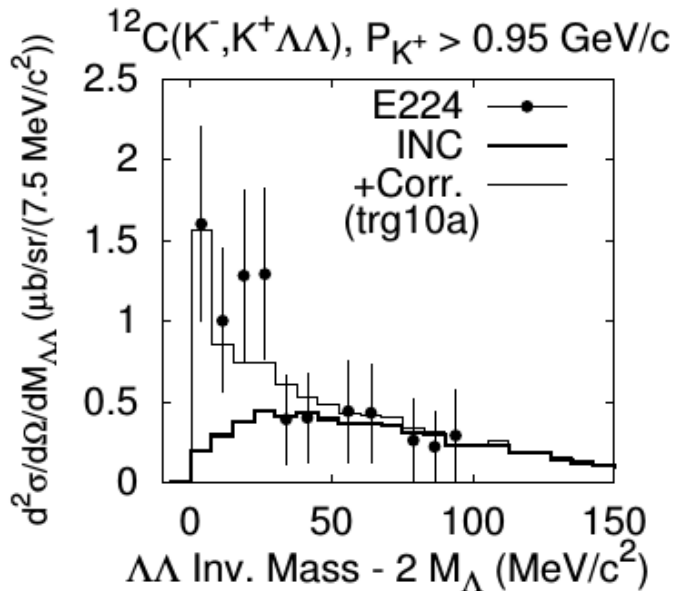
- Enhancement of  $\Lambda\Lambda$  may show  $\Lambda\Lambda$  interaction effects !

AO, Y. Hirata, Y. Nara, S. Shinmura, Y. Akaishi,  
*NPA670('00)297c; NPA684('01)595; NPA691('01)242c*

- Enh. is roughly explained by  $\Lambda\Lambda$  final state int.

- It should be clearer to measure in heavy-ion collisions.  
 Enh. w/o bound state, Suppression w/ bound state.

→ I asked Prof. Huan Z. Hunag (STAR) in **ExHIC 2010 meeting**  
 and they measured it !



AO, Hirata, Nara, Shinmura, Akaishi ('00)