Bound states in femtoscopy – To be (bound) or not to be (*), that is the question – (* Resonance, virtual pole, ..) Akira Ohnishi (YITP, Kyoto U.) TUM, Sep. 18 (Wed), 2019, Munich, Germany

K. Morita, T. Furumoto, AO, PRC91('15)024916 (ΛΛ)
AO, K. Morita, K. Miyahara, T. Hyodo, NPA954 ('16)294 (ΛΛ, K-p)
K. Morita, AO, F. Etminan, T. Hatsuda, PRC94('16), 031901(R) (ΩN)
S. Cho et al. (ExHIC Collab.), Prog.Part.Nucl.Phys.95('17)279 (ΛΛ, K-p)
T. Hatsuda, K. Morita, AO, K. Sasaki, NPA967('17), 856 (Ξ-p)
K. Morita, S. Gongyo, T. Hatsuda, T. Hyodo, Y. Kamiya, AO, arXiv:1908.05414 (NΩ, ΩΩ)





Coauthors of Morita+('19) [arXiv:1908.05414]











Y. Kamiya





Outline

- Introduction
 - Candidates of hadronic bound state
- Correlation function and interaction
 - Scattering length and size dependence of the correlation function
- Bound States in Femotoscopy
 - $\Lambda(1405) (\rightarrow \text{Kamiya})$
 - ΛΛ & ΞΝ
 - (N $\Omega \leftarrow$ FemTUM talk)
 - ΩΩ
- Summary



Exotic Hadrons

Exotic hadrons: Z, X, Y, Pc, (Θ⁺), ... at LEPS, Belle, BaBar, LHCb, BES, ...



- Various pictures
 - Di-quark component
 - Hadronic molecule
 - $Q\overline{Q}$ couples with $Q\overline{Q}$ $q\overline{q}$



- Hadronic bound state candidates
 - Λ(1405), Η, ΝΩ, ΩΩ,



Hadronic Molecule Candidates: Light flavors (u,d,s)

	Particle	m [MeV]	(I, J^p)	qā/qqq (L)	multiquark	Mol. (L)	$\omega_{\rm Mol}$ [MeV]
	$f_0(980)$	980	$(0, 0^+)$	$q\bar{q}\left(P ight)\left(s\bar{s}\left(P ight) ight)$	qsąs	$\overline{K}K(S)$	67.8(B)
	$a_0(980)$	980	$(1, 0^+)$	$q\bar{q}(P)$	qsąs	KK(S)	67.8(B)
	K(1460)	1460	$(1/2,0^{-})$		$qq\bar{q}\bar{s}(P)$	RKK(P)	69 O(R)
	Λ(1405)	1405	$(0, 1/2^{-})$	uds (P)	ud sq q	KN(S)	20.5(R)
	$\Delta\Delta$	2380	$(0, 3^+)$		q°		
1	$\Lambda\Lambda$ -N Ξ (H)	2245	(0,0+)		uuddss	$N\Xi(S)$	73.2(B)
	NΩ	2592	$(1/2, 2^+)$		uudsss		

 $\omega_{\rm mol} = 6$ B.E. or $3/2\mu \langle r^2 \rangle$

S. Cho, T. Hyodo, D. Jido, C. M. Ko, S. H. Lee, S. Maeda, K. Miyahara, K. Morita, M. Nielsen, A. Ohnishi, T. Sekihara, T. Song, S. Yasui, K. Yazaki (ExHIC Collaboration), Prog. Part. Nucl. Phys. 95 (2017), 279-322



Hadronic Molecule Candidates: Heavy flavors (c,b)

Particle	m [MeV]	(I, J^p)	qā/qqq (L)	multiquark	Mol. (L)	ω_{Mol} [MeV]
$D_s(2317)$	2317	(0,0+)	$c\bar{s}(P)$	csqq	DK(S)	273(B)
X(3872)	3872	(0,1+)	$c\bar{c}(P)$	$c\bar{c}q\bar{q}$	$D\bar{D}^{*}(S)$	3.6(B)
$Z_c(3900)$	3900	$(1, 1^+)$		ccud		
$Z_c(4430)$	4430	(1,1+)		ccuđ	$D_1 \bar{D}^*(S)$	13.5(B)
$Z_b(10610)$	10610	(1,1+)		bbud	_	
$Z_b(10650)$	10650	$(1, 1^+)$		bbud	_	
X(5568)	5568	(1,0+)		sbud	_	
$P_c(4380)$	4380	$(1/2, 3/2^{-})^{b}$		c c uud (S)	$D\Sigma_{c}^{*}(S)$	60(B)
$P_{c}(4450)$	4450	$(1/2, 5/2^+)^b$		ccuud (P)	_	—



ExHIC('17)

Hadronic Molecule Candidates: Others

Particle	<i>m</i> [MeV]	(I, J^p)	qq/qqq	multiquark (L)	Mol. (L)	$\omega_{\rm Mol} [{ m MeV}]$
Θ(1530)	1530	$(0, 1/2^+)$		qqqqs (P)		
KKN	1920	$(1/2, 1/2^+)$		qqqss (P)	<i>KKN</i>	42(R)
KNN	2352	(1/2,0-)	$q^5s(P)$	$q^6 s \bar{q}(S)$	KNN	20.5(T)
ΩΩ	3228	(0,0+)	_	s ⁶		—
T^1_{cc}	3797	(0,1+)		udēē		
DN	2790	$(0, 1/2^{-})$		qqqqc	DN	6.48(R)
D^*N	2919	$(0, 3/2^{-})$		$qqqq\bar{c}(D)$	D^*N	6.48(R)
Θ_{cs}	2980	$(1/2, 1/2^+)$		$qqqs\overline{c}(P)$		
H_{c}^{++}	3377	$(1,0^{+})$		qqqqsc		
DNN	3734	$(1/2, 0^{-})$		$q^7 \bar{c}$	DNN	6.48(T)
$\Lambda_c N$	3225	(1/2, 1+)		cuduud	$\Lambda_c N$	4.24(R)
$\Lambda_c NN$	4164	$(0, 3/2^+)$		cuduududd	$\Lambda_c NN$	33.16(R)
T^0_{cb}	7123	$(0, 0^+)$		udēb		—

Many loosely bound hadronic molecules are expected with charm quarks.





ExHIC('17)

Production rate of hadronic molecules

RHIC (Scenario 1)



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Many of exotic hadrons may have hadronic molecule structure (hadronic bound states).

Femtoscopic correlation may reveal their bound state nature !







Correlation Function

Emitting source function

$$N_i(\boldsymbol{p}) = \int d^4x S_i(x, \boldsymbol{p})$$

Two-particle momentum dist.



Assumption: Two particles are produced independently, and the correlation is generated by the final state int.

Koonin('77), Pratt+('86), Lednicky+('82)

$$N_{12}(\boldsymbol{p}_1, \boldsymbol{p}_2) \simeq \int d^4x d^4y S_1(x, \boldsymbol{p}_1) S_2(y, \boldsymbol{p}_2) |\Psi_{\boldsymbol{p}_1, \boldsymbol{p}_2}(x, y)|^2 \\ \mathbf{two-body w.f.} \\ \simeq \int d^4x d^4y S_1(x, \boldsymbol{p}_1) S_2(y, \boldsymbol{p}_2) |\varphi_{\boldsymbol{q}}(\boldsymbol{r})|^2$$

Correlation function

relative w.f.

$$C(\mathbf{p}_1, \mathbf{p}_2) = \frac{N_{12}(\mathbf{p}_1, \mathbf{p}_2)}{N_1(\mathbf{p}_1) N_2(\mathbf{p}_2)} \simeq \int d\mathbf{r} S_{12}(\mathbf{r}) |\varphi_{\mathbf{q}}(\mathbf{r})|^2$$



Wave function around threshold (S-wave, attraction)

Low energy w.f. and phase shift

 $u(r) = qr\chi_q(r) \to \sin(qr + \delta(q)) \sim \sin(q(r - a_0))$ $q \cot \delta = -\frac{1}{a_0} + \frac{1}{2}r_{\text{eff}}q^2 + \mathcal{O}(q^4) \ (\delta \sim -a_0q)$

- Wave function grows rapidly at small r with attraction.
- With a bound state $(a_0 > 0)$, a node appears around $r=a_0$



From correlation function to hadron-hadron interaction

Small q behavior rχ rχ rχ sin gr sin gr sin ar R • $\mathbf{R}/\mathbf{a}_0 < \mathbf{0} \rightarrow \mathbf{C}(\mathbf{q}) > \mathbf{1}$ • $-1 < R/a_0 < 0.3$ loose <u>b.s.</u> a₀ No b.s. \rightarrow C(q) >> 1 deep b.s. V $a_0 < 0 < R$ $\mathbf{a}_0 > \mathbf{R}$ $0 < a_0 < R$ • $0.5 < R/a_0 \rightarrow C(q) < 1$ 3 $R/a_0 = -0.3$ Loosely bound state (large a_0) larger R 0.3 2 \rightarrow C(q)>>1 at small R (d) (d) C(q)< 1 at large R w/o bound state (a₀<0) in the case w/o Coulomb, w/ bound state $(a_0 > 0)$ w/o coupled channel, LL model w/o quantum stat. 0 0.5 1.5 2 0 qR



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Lednicky-Lyuboshits (LL) model

Lednicky-Lyuboshits analytic model

• Asymp. w.f. + Eff. range corr. +
$$\psi^{(\cdot)} = [\psi^{(+)}]^*$$

 $\psi_0(r) \rightarrow \psi_{asy}(r) = \frac{e^{-i\delta}}{qr} \sin(qr+\delta) = S^{-1} \left[\frac{\sin qr}{qr} + f(q) \frac{e^{iqr}}{r} \right]$
 $\Delta C_{LL}(q) = \int d\mathbf{r} S_{12}(r) \left(|\psi_{asy}(r)|^2 - |j_0(qr)|^2 \right)$

$$= \frac{|f(q)|^2}{2R^2} F_3\left(\frac{r_{\text{eff}}}{R}\right) + \frac{2\text{Re}f(q)}{\sqrt{\pi}R}F_1(x) - \frac{\text{Im}f(q)}{R}F_2(x)$$

 $(x = 2qR, R = \text{Gaussian size}, F_1, F_2, F_3 : \text{Known functions})$

Phase shifts

$$q \cot \delta = -\frac{1}{a_0} + \frac{1}{2}r_{\text{eff}}q^2 + \mathcal{O}(q^4) \rightarrow \delta \simeq -a_0q + O(q^3)$$
$$\sin(qr + \delta) \simeq \sin(q(r - a_0) + \cdots) \qquad \begin{array}{l} \text{Node at } \mathbf{r} \sim \mathbf{a_0} \\ \text{for small } \mathbf{q} \end{array}$$



C(q) in the low momentum limit

• Correlation function at small q (and $r_{eff}=0$) \rightarrow $F_1=1$, $F_2=0$, $F_3=1$

$$\Delta C_{\rm LL}(q) \rightarrow \frac{|f(0)|^2}{2R^2} + \frac{2\text{Re}f(0)}{\sqrt{\pi}R} \quad (q \rightarrow 0)$$

$$f(q) = (q \cot \delta - iq)^{-1} \simeq \left(-\frac{1}{a_0} + \frac{1}{2}r_{\rm eff}q^2 - iq\right)^{-1} \rightarrow -a_0$$

$$C_{\rm LL}(q \rightarrow 0) = 1 + \frac{a_0^2}{2R^2} - \frac{2a_0}{\sqrt{\pi}R} = 1 - \frac{2}{\pi} + \frac{1}{2}\left(\frac{a_0}{R} - \frac{2}{\sqrt{\pi}}\right)^2$$

$$1 - 2/\pi \simeq 0.36, \quad \sqrt{\pi}/2 \simeq 0.89$$

 $C(q \rightarrow 0)$ takes a minimum of 0.36 at R/a₀ = 0.89 in the LL model.



Correlation Function in LL model



LL model: R. Lednicky, V. L. Lyuboshits ('82)



Correlation Function with Gaussian source



NΩ potential (J=2, HAL QCD) + Coulomb







Example 1: *A*(1405)

- A(1405) as a bound state of KN Dalitz, Wong ('67); Siegel, Weise ('88)
- How can we confirm it ?



- Binding Energy

 → Deep (~ 30 MeV) or Shallow (~ 10 MeV) binding
 Akaishi, Yamazaki ('02); Jido, Oller, Oset, Ramos, Meissner ('03);
 Hyodo, Weise ('08); Shevchenko, Gal, Mares ('07); Ikeda, Sato ('07), ...
- Medium effects → Upward mass shift or No change Koch ('94); Waas, Kaiser, Weise ('96); Lutz ('98)
- In-medium Branching Ratio Change → Indirect (?) Vander Velde-Wilquet et al. ('77); AO, Nara, Koch ('98)
- Form factor \rightarrow Not measured yet

Sekihara, Hyodo, Jido ('08, '11)

Correlation Function
 AO+('16), S. Cho et al.(ExHIC)('17), S. Acharya et al. (ALICE), arXiv:1905.13470,
 Y. Kamiya, T. Hyodo, K. Morita, AO, in prep.



K-*p* correlation function

Strong source size dependence seems to be observed in pA collisions, which may signal the bound state nature of Λ(1405)



S. Acharya et al. (ALICE), arXiv:1905.13470



Kamiya, Morita, Hyodo, AO (in prep.)

Example 2: H particle

- H-particle: 6-quark state (uuddss)
 - Prediction: R.L.Jaffe, PRL38(1977)195
 - Ruled-out by double Λ hypernucleus Takahashi et al., PRL87('01) 212502
 - Resonance or Bound "H" ? Yoon et al.(KEK-E522)+AO ('07)
 - Lattice QCD
 HAL QCD & NPLQCD ('11)
 K. Sasaki et al. (HAL QCD, '16-'18)





AA correlation in HIC

- Resonance H particle search via femtoscopy C. Greiner, B. Muller, PLB219('89),199
- Interaction study via ΛΛ correlation in HIC AO, Y. Hirata, Y. Nara, S. Shinmura, Y. Akaishi, NPA670 ('00), 297c



AO, Hirata, Nara, Shinmura, Akaishi, NPA670('00)297c

$\Lambda\Lambda$ correlation from (K⁻,K⁺ $\Lambda\Lambda$) reaction





A. Ohnishi, TUM, Sep. 18, 2019 23

AA correlation at RHIC

- STAR collaboration at RHIC measured ΛΛ correlation ! Adamczyk et al. (STAR Collaboration), PRL 114 ('15) 022301.
 - RHIC, Au+Au ($\sqrt{s_{NN}}$ =200 GeV), Weak decay vertex analysis.
- Theoretical Analysis well explains the data K.Morita et al., T.Furumoto, AO, PRC91('15)024916; AO, K.Morita, K.Miyahara, T.Hyodo, NPA954 ('16), 294.



AA interaction

- Propsed ΛΛ interactions
 - Meson Ex. models: Nijmegen model D, F, Soft Core (89, 97), ESC08 Nagels, Rijken, de Swart ('77, '79), Maessen, Rijken, de Swart ('89), Rijken, Stoks, Yamamoto ('99); Rijken, Nagels, Yamamoto ('10).
 - Quark cluster model interaction: fss2 Fujiwara, Fujita, Kohno, Nakamoto, Suzuki ('00)
 - Phenomenological model: Ehime T. Ueda et al. ('99).
- Two (or three) range gaussian fit results are used in the analysis.



ΛΛ interaction from ΛΛ correlation



2019 26

Feed-Down Effects & Residual Source

Correlation Fn. w/ Feed-down & Residual source effects.





New Data from LHC-ALICE

S. Acharya+ (ALICE), PLB 797('19), 134822 [1905.07209]



Weakly attractive V_{14} .

Large reff \rightarrow Becomes repulieve at relatively low density.





Time dependence of AA interaction





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Detecting Dibaryon State from Invariant Mass Spectrum ?

- Strong signal of dibaryon(s) in correlation function
 - → How about invariant mass spectrum
 - \rightarrow Needs much more statistics



J. Haidenbauer will make an objection ...

Morita+ ('15)



Ξ^{-} p correlation

Prediction of the correlation function by using EN potential (HAL QCD Collab.) + Coulomb potential





Example 3: Ω N dibaryon

- Ω : sss, J π =3/2+, M=1672 MeV
- Is there an ΩN bound state (S= -3 dibaryon)?
 - Predicted as a dibaryon candidate Goldman+ ('87), Oka ('88), Gal ('16)
 - Lattice QCD predicts a bound state with narrow width for J=2 (⁵S₂)
 Etminan+ (HAL QCD)('14), Iritani+ (HAL QCD) ('19)
 - Correlation function is measurable ! *Adam+ (STAR)('19), Oton (FemTUM), ALICE, in prep.*





ΩN potential from lattice QCD

- ΩN potential by HAL QCD Collab. (J=2)
 - m_π=875 MeV, B.E.~ 19 MeV
 F.Etminan et al. (HAL QCD Collab.), NPA928('14)89.
 - m_{π} =146 MeV, B.E.~ 2.2 MeV
 - T. Iritani et al. (HAL QCD Collab.), PLB 792('19)284.





Binding Energy, Scattering Length, Effecitve Range ($p\Omega$)

 m_{π} =875 MeV, B.E.~ 19 MeV

F.Etminan et al. (HAL QCD Collab.), NPA928('14)89.

m_{π} =146 MeV, B.E.~ 2.2 MeV

T. Iritani et al. (HAL QCD Collab.), PLB 792('19)284.

Spin-2 $N\Omega$ Potentials		V_{I}	$V_{\rm II}$	$V_{\rm III}$	t/a	a ₀ [fm]	reff [fm]	E_B [MeV]
	$E_{\rm B} [{\rm MeV}]$	-	0.05	24.8	11	3.45	1.33	2.15
without Coulomb	a_0 [fm] $r_{-\sigma}$ [fm]	-1.0 1.15	23.1	0.65	12	3.38	1.31	2.27
	$E_{\rm B}$ [MeV]	_	6.3	26.9	13	3.49	1.31	2.08
with Coulomb	a_0 [fm]	-1.12	5.79	1.29	14	2.40	1.22	2.00
	$r_{\rm eff}$ [fm]	1.16	0.96	0.65	14	5.40	1.55	2.24

 $(m_{\pi} = 146 \text{ MeV}, m_N = 955 \text{ MeV}$ and $m_{\Omega} = 1712 \text{ MeV})$. By using the same parameter set for t/a = 12 in Table 1 with $m_{\pi} = 146$ MeV kept fixed but with physical baryon masses $(m_p = 938 \text{ MeV} \text{ and } m_{\Omega^-} = 1672 \text{ MeV})$, we find less binding than Eq. (10) as expected: $B_{p\Omega^-} \simeq 2.18(32) \text{ MeV}$ and $\sqrt{\langle r^2 \rangle}_{p\Omega^-} \simeq 3.45(22)$ fm. On the other hand, if we additionally em-



Calculation Details

- NΩ potential from HAL QCD Collab. Etminan+(HAL QCD) ('14), Iritani+ (HAL QCD)('19)
 - J=1 potential is uncertain → Three models Strong abs. at r < r₀ (r₀ ~ 2 fm) (*Morita*+('16)) (Standard) Complete absorption χ (J=1) = 0 (Minimum) Same w.f. as that with J=2, χ (J=1) = χ (J=2) (Reference)
 - Statistical Error can be evaluated by using Jackknife potentials.
- Coulomb potential enhances CF even without strong int. → Small-Large ratio of CF (*Morita*+('16))
 - Large source → Coulomb force dominate
 Small source → Visible strong interaction effects
- Source function: Blast wave, Gaussian source



Emission Source Function

- Gaussian Source ∝ exp(-r²/4R²), R=(0.8-4) fm [Simple and convenient]
 Flow velocity
- **Expanding source model** [Reasonably realistic] $u_{\mu}(x)$

$$d^{4}xS_{i}(x, \mathbf{p}) = \tau_{0}d\eta_{s}d^{2}r_{T}\frac{d}{(2\pi)^{3}}\underline{n_{f}(u \cdot p, T)}\exp\left(-\frac{r_{T}^{2}}{2R_{T}^{2}}\right)$$

Fermi dist.



■ Transport model result [should be realistic] → Future work



$\Omega^{-}p$ correlation



(w/o Coulomb, Strong absorption at r< 2 fm in ³S₁ (decay to 8-8 in S-wave))

K. Morita, AO, F. Etminan, T. Hatsuda, PRC94('16)031901(R) [arXiv:1605.06765 [hep-ph]]



Ω^{-} p correlation w/ Coulomb



Coulomb potential washes out the features of V_I , V_{II} , V_{III} , and Gamow correction is too strong.





Ω⁻ p correlation: Small / Large Ratio



By taking small (R=2.5 fm) / large (R=5 fm) ratio, we approximately see the corr. fn. w/o Coulomb !

Comparison with STAR data

Dip structure in Small-Large ratio of CF \rightarrow We may have a dibaryon state in ΩN channel





Source Size Dependence of Correlation Function



K. Morita, S. Gongyo, T. Hatsuda, T. Hyodo, Y. Kamiya, AO ('19)

A. Ohnishi, TUM, Sep. 18, 2019 41

Comparison with STAR data

Results with potential at nearly physical quark mass (= between V_{II} and V_{III})

 \rightarrow Dip is seen but is not deep enough to explain STAR data.





Comparison with STAR data

Results with Gaussian source → Agrees with STAR data within 1σ

Detail discussion of source is also necessary....

Morita, Gongyo, Hatsuda, Hyodo, Kamiya, AO, arXiv:1908.05414 STAR (1808.02511, PLB790 ('19) 490)





ALICE Preliminary

■ pp 13 TeV high-multiplicity events in ALICE → Strong enhancement of CF at small q

O. Vázquez Doce et al. (ALICE), Hadrons 2019





Source Size Dependence of Correlation Function



Gaussian Source



K. Morita, S. Gongyo, T. Hatsuda, T. Hyodo, Y. Kamiya, AO ('19)

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Correlation Function in LL model



 a_0 (p Ω)~3.4 fm, R(ALICE)~0.7 fm, R(STAR)~3 fm

Correlation Function with Gaussian source



NΩ potential (J=2, HAL QCD, a₀=3.4 fm) + Coulomb



$STAR + ALICE = N\Omega$ Dibaryon





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Example 4: $\Omega\Omega$ dibaryon

- ΩΩ potential from lattice QCD (J=0)
 S. Gongyo et al. (HAL QCD Collab), Phys. Rev. Lett. 120, 212001 (2017).
 - $\Omega\Omega$ bounds for J=0 ! (Most strange dibaryon state)
 - B.E. is very small. B.E.=(0.1-1.0) MeV $\rightarrow a_0 > 10$ fm



$\Omega\Omega$ correlation

• $\Omega\Omega$ correlation

K. Morita, S. Gongyo, T. Hatsuda, T. Hyodo, Y. Kamiya, AO, arXiv:1908.05414

- $J \neq 0$ potentials are assumed to be zero.
- Large scattering length

 \rightarrow R/a₀ < 0.5 (unitary regime), but Coulomb suppressed





Detecting $\Omega\Omega$ dibaryon ?

- Large DCA Ω ? (J. Schaffner-Bielich)
 - Implies a weakly decaying ΩX bound state.
 - Fourier transf. of w.f. in $\Omega\Omega \to \Omega$ momentum dist.





Summary

- It is fun to find hadronic bound states.
- Bound states show characteristic features in C(q) at small q, Enhanced C(q) for R/a₀ < 0.4,
 Suppressed C(q) for 0.5 < R/a₀ < 2
- ALICE and STAR data strongly suggest the existence of S= -3 dibaryon as a bound state of NΩ.
 - Strong enhancement of pΩ corr. fn. found in ALICE data implies large |a₀|.
 - Coulomb potential enhances C(q).
 - Then I do not know the mechanism to show suppressed C(q) other than the existence of a bound state.
- Source size dep. of $p\Xi^-$ and pK^- corr. fn. would be interesting.
- Detecting diomega ? Fantastic.



To do (or Can do)

- Coupled channel effects → Talk by Haidenbauer, Kamiya
- Ξ⁻ p correlation with updated HAL QCD potential (K. Sasaki et al.) → Sasaki's talk
- Ap correlation with various potentials (χEFT, Nijmegen, fss2, lattice) → Talk by Heidenbauer, Rijken
- K⁻ p correlation with amplitude from chiral SU(3) dynamics (e.g. Ikeda, Hyodo, Weise) → Kamiya's talk
- **J=1** Ω^- p potential \rightarrow Hyodo's talk
- Let us use deuteron !
 K⁻ d corr. (I=0 ampl.), Ad corr. (³ H B.E., 3BF), Ξ⁻ d, ...
- Can we go to heavy-quarks ?
 cτ(D) = 0.3 mm → γcτ(D) = cτ(D) cosh y ~ 15 cm (y=7)
 We may have enough D mesons at y=7 in fixed target LHC.
 (N. Yamanaka)



Thank you for your attention !

Coauthors of Morita+('19) [arXiv:1908.05414] K. Morita S. Gongyo T. Hatsuda T. Hyodo









Y. Kamiya





AO

AA correlation at RHIC

- STAR collaboration at RHIC measured ΛΛ correlation ! Adamczyk et al. (STAR Collaboration), PRL 114 ('15) 022301.
- Theoretical Analysis well explains the data K.Morita et al., T.Furumoto, AO, PRC91('15)024916; AO, K.Morita, K.Miyahara, T.Hyodo, NPA954 ('16), 294.
- New Data from ALICE

S. Acharya+ (ALICE), PLB 797('19), 134822 [1905.07209]



AA interaction from **AA** correlation



 $q \cot \delta = -1/a_0 + r_{eff} q^2/2 + O(q^4)$

NF NSC89 NSC97 ESC08c Ehime fss2 FG нкмүү STAR

- Nijmegen potentials (ND, NF, NSC89, NSC97, ESC08) Nagels+('77, '79), Maessen+('89), *Rijken+('99,'10)*
- Ehime *Ueda et al. ('98)*
- **Quark model interaction:** fss2 Fujiwara et al.('07)
- Potential fitted to Nagara Filikhin, Gal ('02) (FG), *Hiyama et al. ('02, '10)(HKMYY)*



Additional Source

Feed down effects

$$C_{\text{corr}}(Q) = 1 + \lambda(C_{\text{bare}}(Q) - 1)$$

 $\lambda = Purity of \Lambda\Lambda pair$

- Short-lived $Y^* \rightarrow mod.$ of source fn.
- $\Xi \rightarrow \Lambda \pi$ can be excluded (c τ =8.71 cm)
- $\Sigma^0 \to \Lambda \gamma$ is difficult to reject
- Data based purity λ=(0.67)² Σ⁰/Λ=0.278 (p+Be, 28.5 GeV/c) Sullivan et al. ('87) Ξ/Λ = 15 % (RHIC)
- "Residual" source
 - High-momentum tail $\rightarrow R_{res} \sim 0.5$ fm (STAR collab.)





AA invariant mass / BB correlation function (as of 2016)



 (K^{-}, K^{+}) reaction

(K⁻, K⁺) reaction = doorway to produce S=-2 systems

• $K^- p \to K^+ \Xi^- \to \Xi$ nuclei, stopped Ξ to $\Lambda\Lambda$ nuclei, ...



Two A production in (K⁻, K⁺) reactions

- Experimentalists really measured two Λ emission ! J.K.Ahn et al. (KEK-PS E224 Collab.), PLB444('98)267.
 C. J. Yoon et al. (KEK-PS E522 Collab. +AO), PRC75('07)022201
- Invariant mass spectrum of ΛΛ is enhanced from our cascade calculation.
 - → FSI enhancement ? or H particle ?



From correlation to interaction

- Enhancement of ΛΛ may show ΛΛ interaction effects ! AO, Y. Hirata, Y. Nara, S. Shinmura, Y. Akaishi, NPA670('00)297c; NPA684('01)595; NPA691('01)242c
 - Enh. is roughly explained by $\Lambda\Lambda$ final state int.
 - It should be clearer to measure in heavy-ion collisions. Enh. w/o bound state, Suppression w/ bound state.
 - \rightarrow I asked Prof. Huan Z. Hunag (STAR) in ExHIC 2010 meeting and they measured it !

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A. Ohnishi, TUM, Sep. 18, 2019



