
Femtoscopic study of $N\Xi$ interaction and search for the H dibaryon state around the $N\Xi$ threshold

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in collaboration with

Y. Kamiya, K. Sasaki, T. Fukui, T. Hatsuda, T. Hyodo, K. Morita, K. Ogata
APFB2020, Mar. 1-5, 2021, Kanazawa / Online (Hybrid)

■ Introduction

- Femtoscopic diagnosis of bound state existence

■ $N\Xi$ - $\Lambda\Lambda$ interaction and $p\Xi^-$ and $\Lambda\Lambda$ correlation functions

- HAL QCD potential / Comparison with data and previous results

■ Summary

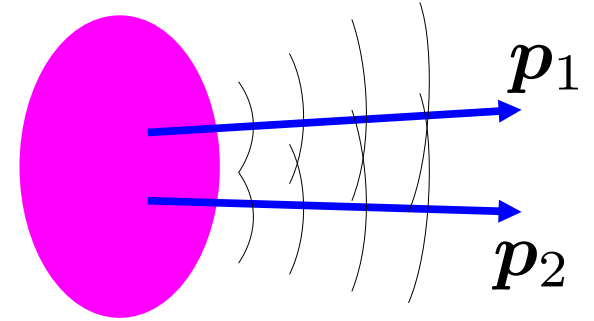
*Y. Kamiya, K. Sasaki, T. Fukui, T. Hatsuda,
T. Hyodo, K. Morita, K. Ogata, AO, in prep.*

Introduction

Femtoscopic study of hadron-hadron interaction

- **Koonin-Pratt formula**

Koonin('77), Pratt+('86), Lednicky+('82)



$$C(\mathbf{p}_1, \mathbf{p}_2) = \frac{N_{12}(\mathbf{p}_1, \mathbf{p}_2)}{N_1(\mathbf{p}_1)N_2(\mathbf{p}_2)} \simeq \int d\mathbf{r} \underbrace{S_{12}(\mathbf{r})}_{\text{source fn.}} \underbrace{|\varphi_q(\mathbf{r})|^2}_{\text{relative w.f.}}$$

- **Standard usage = Source size using quantum statistics**

Hanbury Brown & Twiss ('56); Goldhaber, Goldhaber, Lee, Pais ('60)

- **Non-Std. usage = Hadron-hadron interaction**

CF shows how much $|\varphi|^2$ is enhanced $\rightarrow V_{\text{hh}}$ effects !

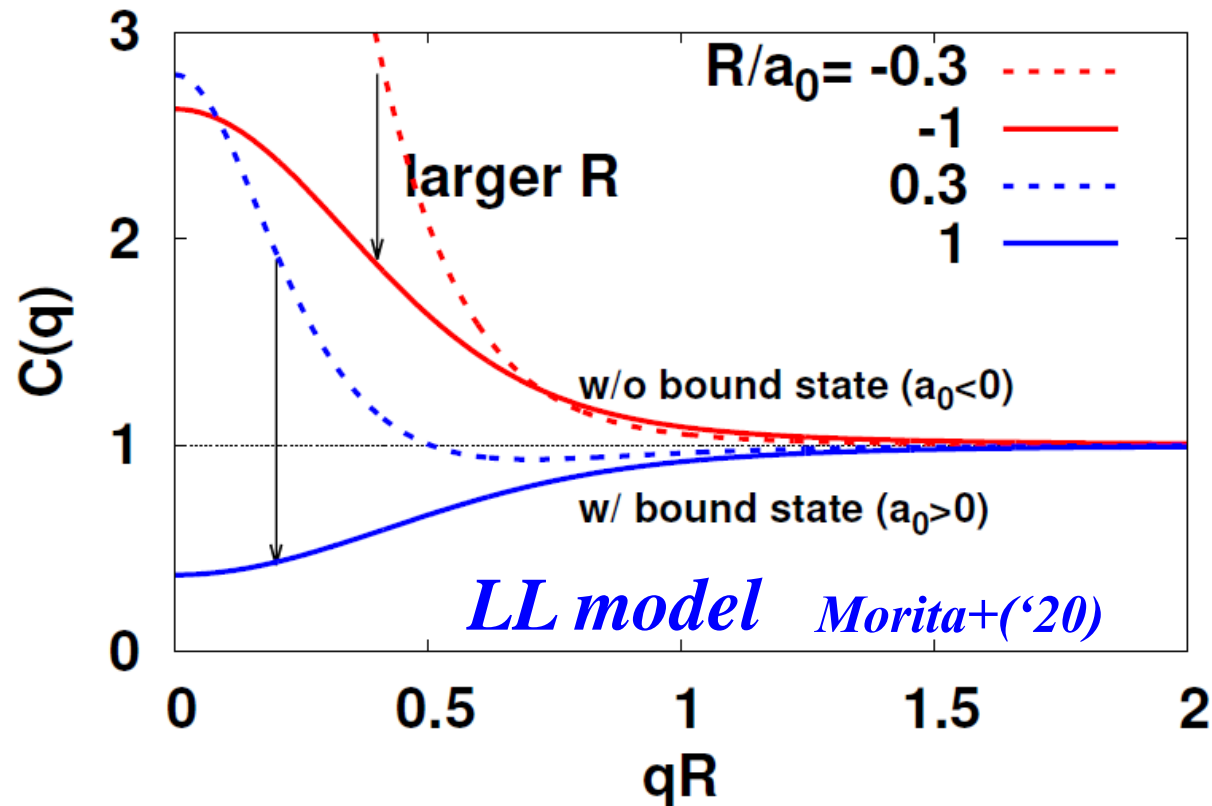
R. Lednicky, V. L. Lyuboshits ('82); K. Morita, T. Furumoto, AO (1408.6682)

$$C(\mathbf{q}) = 1 + \int d\mathbf{r} S(\mathbf{r}) \{ |\varphi_0(\mathbf{r})|^2 - |j_0(q\mathbf{r})|^2 \} \quad (\varphi_0 = \text{s-wave w.f.})$$

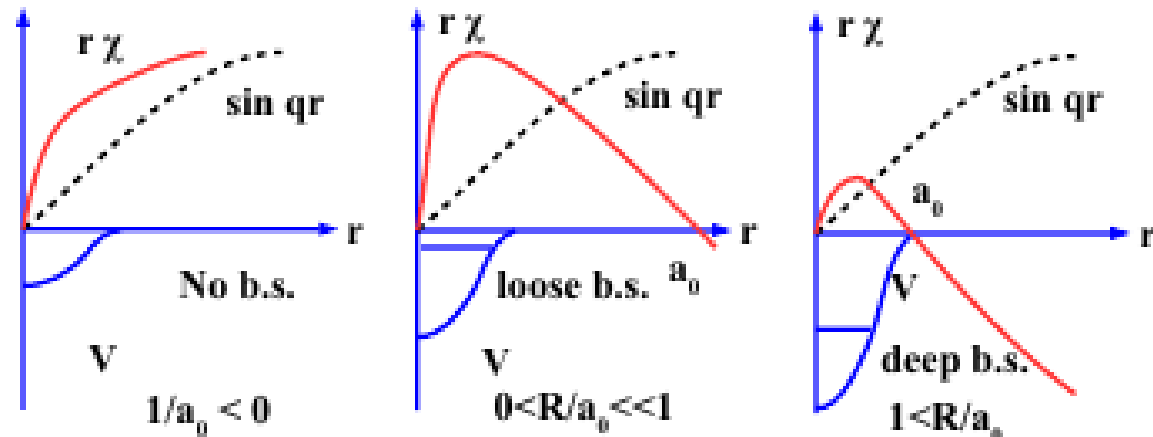
Bound state diagnosis by femtoscopy

■ Source size dep. of CF tells the sign of the scattering length (a_0).

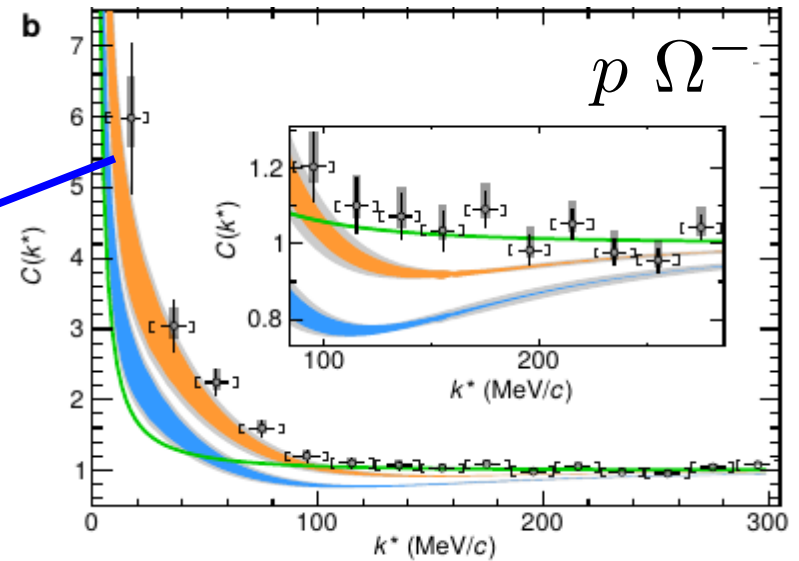
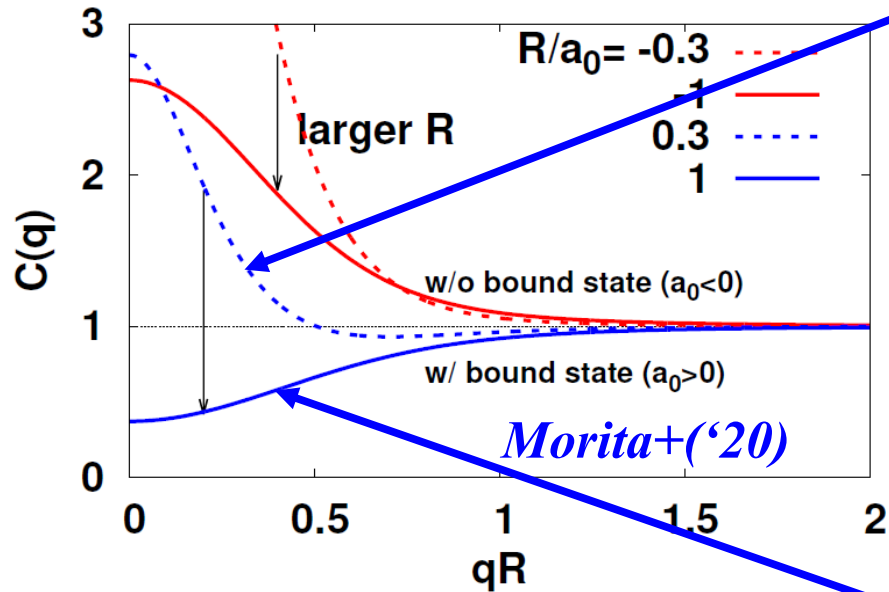
- With attraction, CF is enhanced at small R.
- With $a_0 > 0$, CF is suppressed at $R \sim a_0$



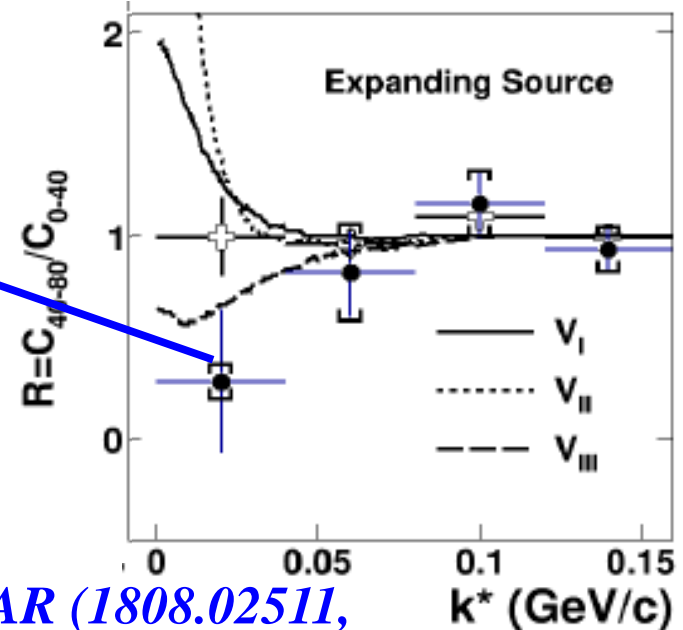
Source size dep. of CF
 → *To be bound,*
or not to be bound.



ALICE+STAR = $N\Omega$ Dibaryon



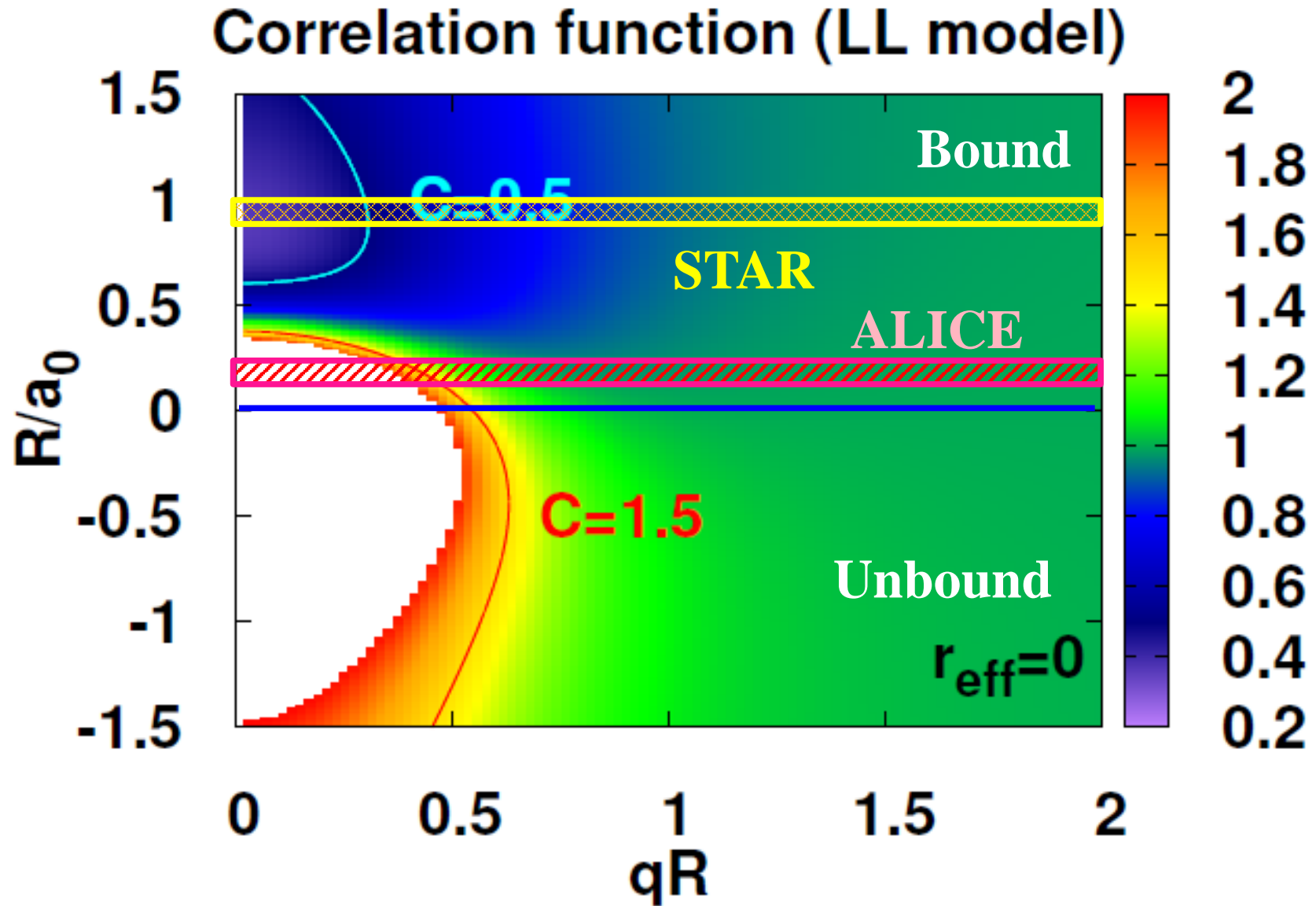
Acharya+(ALICE), Nature ('20)



STAR (1808.02511,
PLB790 ('19) 490)

c.f. Fabbietti's talk on Monday

Source Size Dependence of Correlation Function



LL model: R. Lednicky, V. L. Lyuboshits ('82)
AO, K.Morita, K.Miyahara, T.Hyodo (1605.06765)

To be bound or not to be bound, the question for H

- **H-dibaryon: 6-quark state (uuddss)**

- Prediction: *R.L.Jaffe, PRL38(1977)195*

- Ruled-out by double Λ hypernucleus
Takahashi et al., PRL87('01) 212502

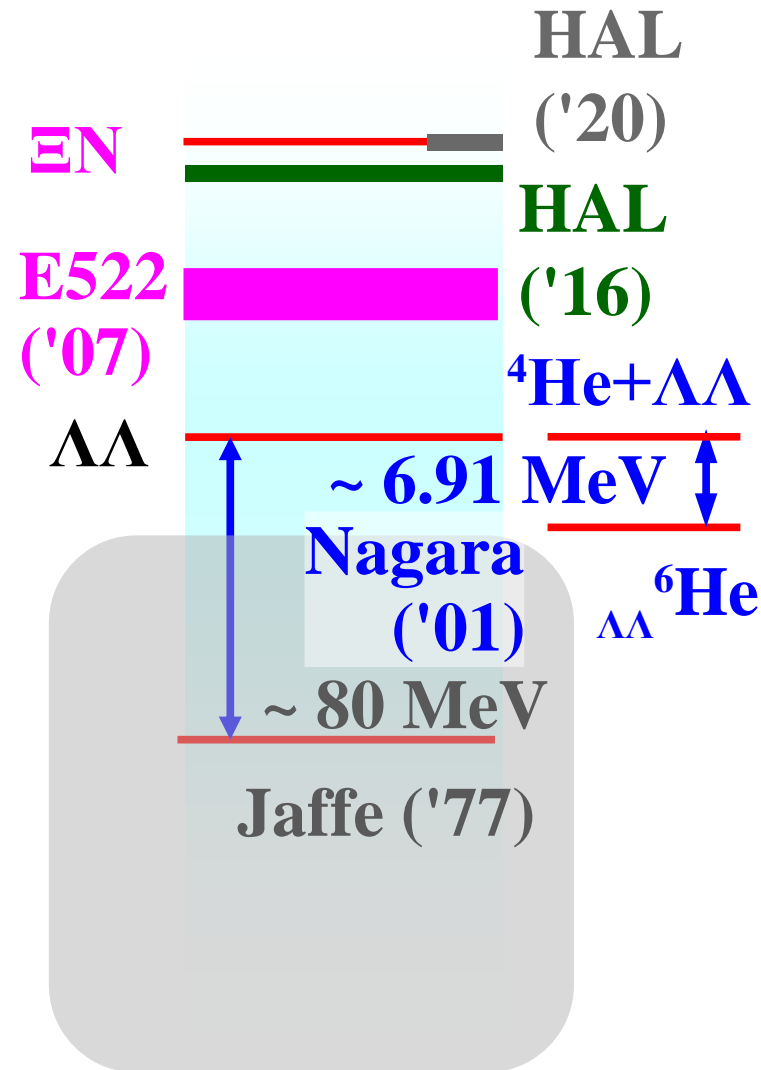
- Resonance or Bound “H” ?
Yoon et al.(KEK-E522)+AO ('07)

- **Lattice QCD results**

- Bound:
HALQCD('11), NPLQCD('11,'13), Mainz('19)
(heavier quark mass or SU(3) limit)

- Resonance (Bound state of $N\Xi$):
HAL QCD ('16,18) (heavier m_q)

- Virtual Pole (around $N\Xi$ threshold)
HAL QCD ('20) (almost physical m_q)



We examine LQCD $N\Xi$ - $\Lambda\Lambda$ potential and discuss H using CF !

*$N\bar{E}$ interaction
and $p\bar{E}^-$ correlation function*

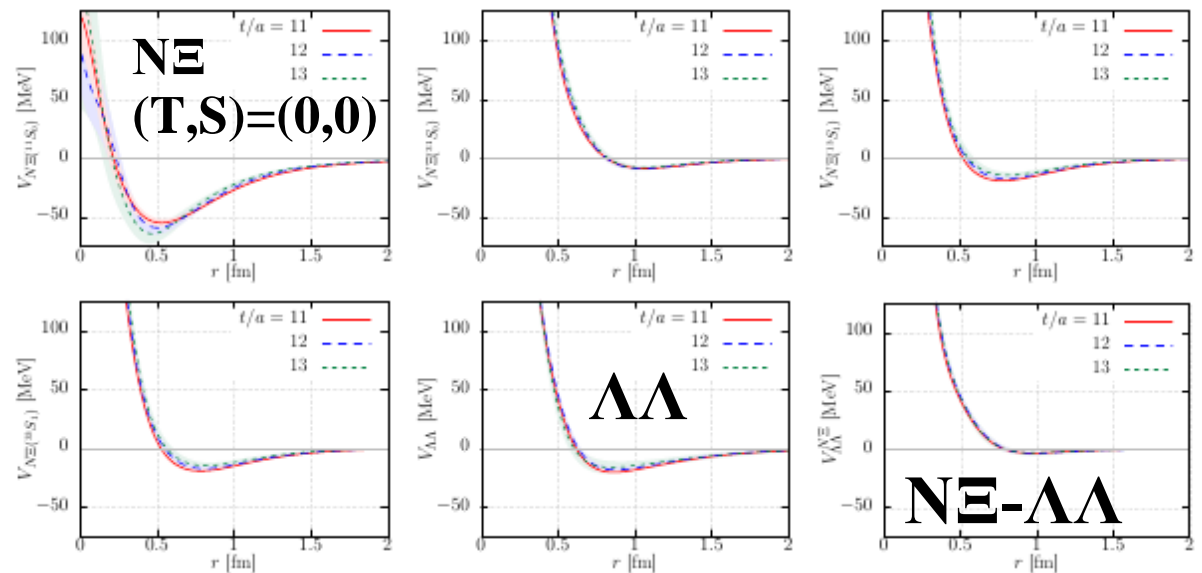
$N\Xi$ - $\Lambda\Lambda$ potential from Lattice QCD

- $N\Xi$ - $\Lambda\Lambda$ potential at almost physical quark mass ($m_\pi=146$ MeV)
K. Sasaki et al. (HAL QCD), NPA 998 ('20) 121737 (1912.08630)

- Strong attraction in $(T,S)=(0,0)$ $N\Xi$ channel

- Weak attraction in $\Lambda\Lambda$ → Coupling with $N\Xi$ causes $\Lambda\Lambda$ attraction.

- There is no bound state in $N\Xi$ - $\Lambda\Lambda$ system (except for Ξ^- atom), but there is a **virtual pole** around the $N\Xi$ threshold (3.93 MeV below $n\Xi^0$ threshold) on the irrelevant Riemann sheet, (+, -, +) [relevant=(-,+,+)]



$$E_{\text{pole}} = 2250.5 \pm i0.3 \text{ MeV}$$

Sign of $\text{Im}(\text{eigen momentum})$

Correlation Function with Coupled-Channel Effects

- Correlation function formula (Koonin-Pratt formula) with CC
 - sum of channel contributions to the relevant channel with outgoing momentum q

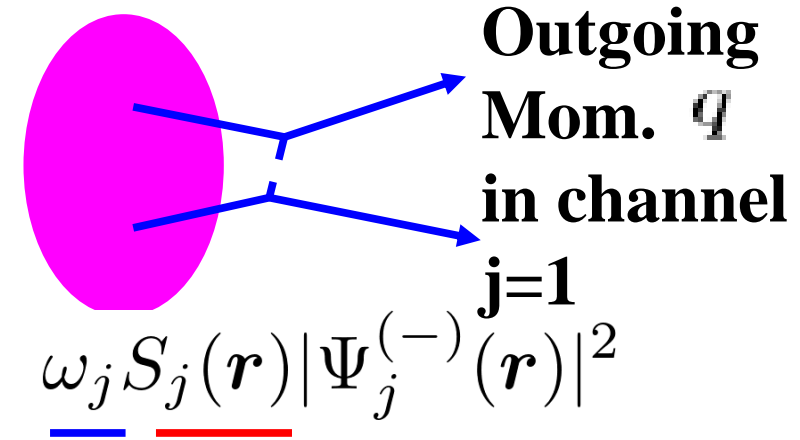
Lednicky, Lyuboshits ('81); Haudenbauer ('19)

$$C(q) = \sum_j \omega_j \int d\mathbf{r} S_j(\mathbf{r}) |\Psi_j^{(-)}(\mathbf{r})|^2$$

$$\Psi_j^{(-)}(\mathbf{r}) = [e^{i\mathbf{q}\cdot\mathbf{r}} - j_0(qr)]\delta_{1j} + \psi_j^{(-)}(r)$$

$$\psi_j^{(-)}(q) \propto e^{-iqr}/r \text{ or } e^{-\kappa r} \quad (r \rightarrow \infty)$$

(No Coulomb case)

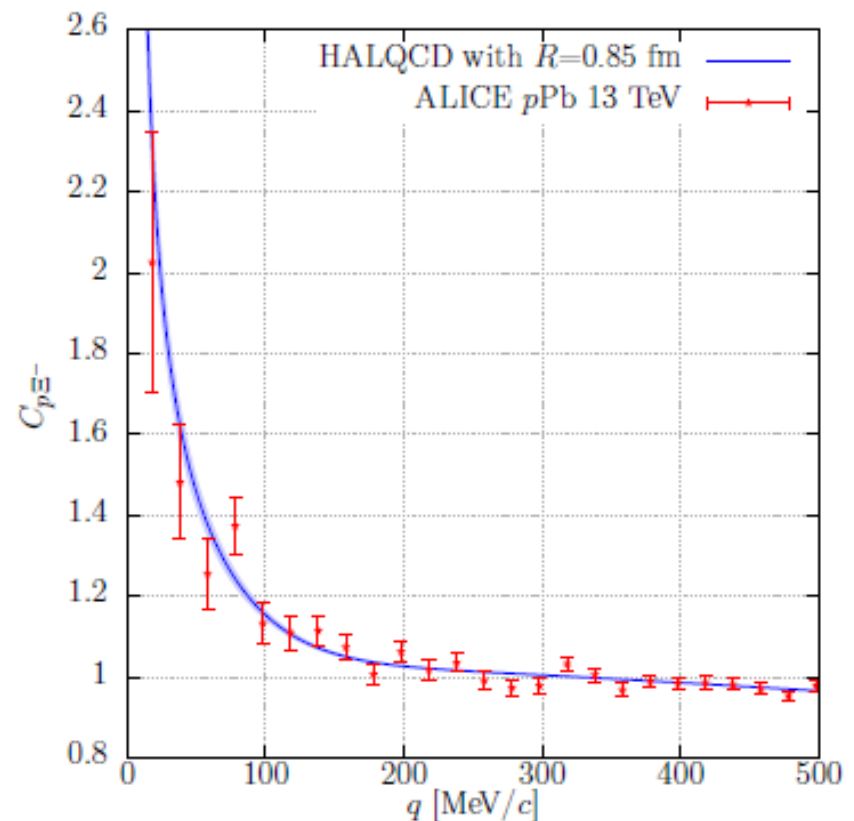
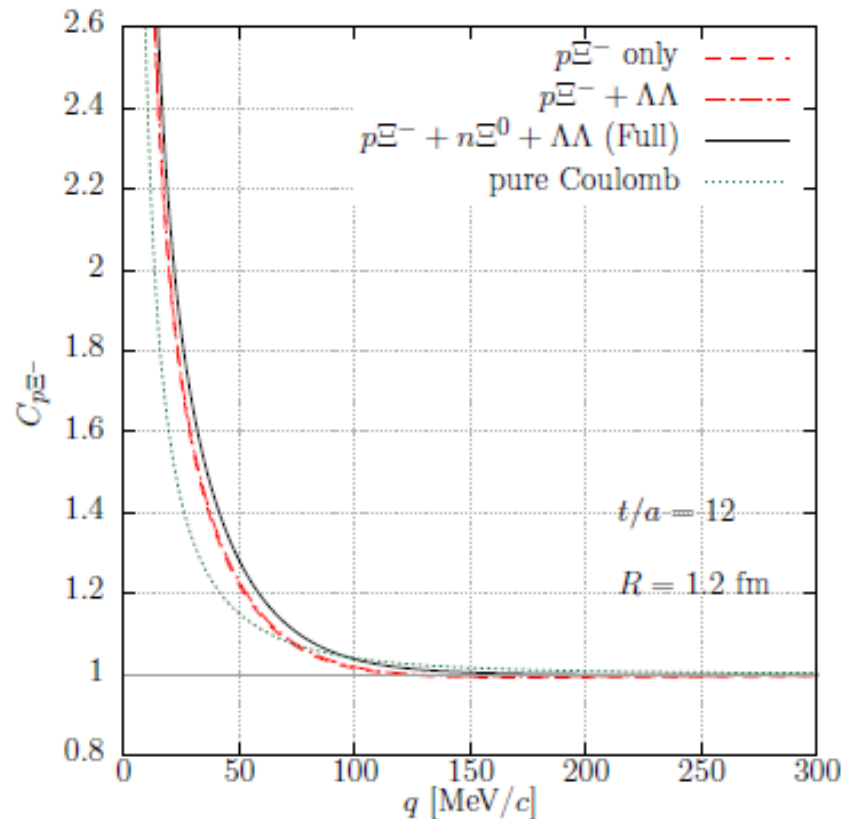


Source weight Normalized Source fn.

- Effects of coupled-channel, strong & Coulomb pot., and threshold difference are taken into account in the charge base, $p\Xi^-$, $n\Xi^0$, $\Lambda\Lambda$.
Y. Kamiya+, PRL('20, $K^- p$)
- Source weights (ω_j) are chosen to be unity.
- Source size R is taken as the parameter.

$p\Xi^-$ correlation function

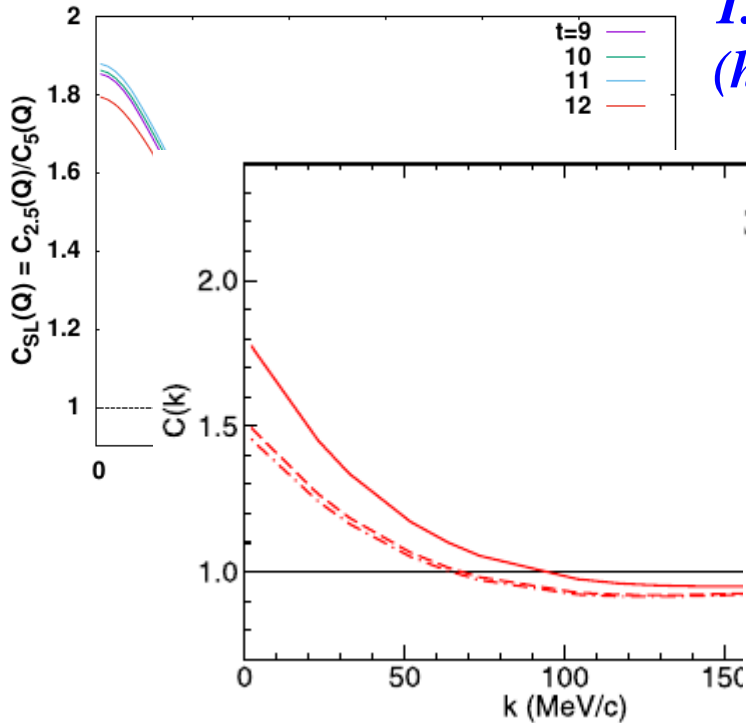
- $p\Xi^-$ correlation function data implies attractive $N\Xi$ interaction.
 - Calculated CF agrees with ALICE data.
 - Strong enhancement from pure Coulomb CF
 - $\Lambda\Lambda$ source effect is negligible. $n\Xi^0$ source effect is visible.



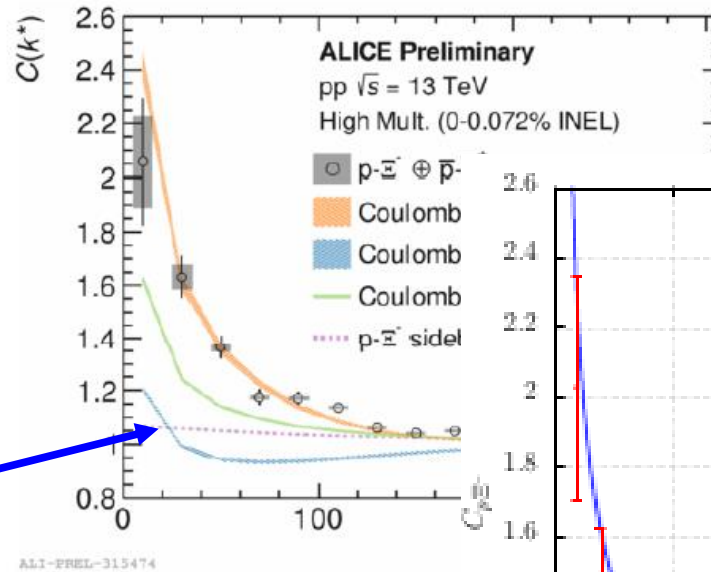
Kamiya+, in prep.; Acharya+(ALICE), Nature ('20)

Comparison with other results

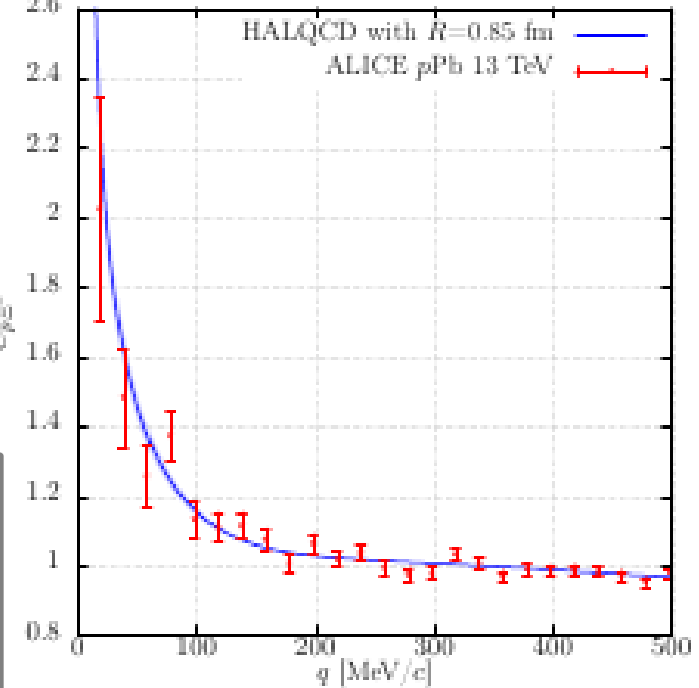
*T. Hatsuda, K. Morita, AO, K. Sasaki, NPA967('17)856.
(heavier quark mass)*



*J. Haidenbauer, NPA981('19)1.
(NLO(600), w/ CC effects, w/o Coulomb)
(w/ Coulomb, it will be comparable with data.)*



*D. L. Mihairov+[ALICE],
NPA1005('21)121760 (QM2019).
(Nijmegen potential does not
explain the data.)*

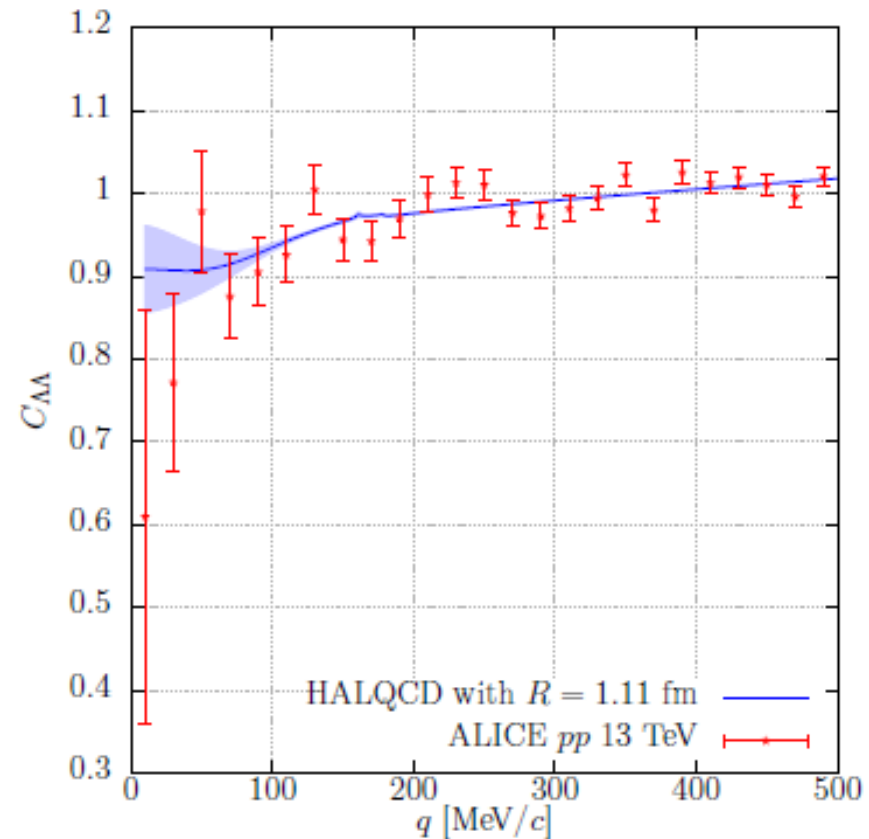
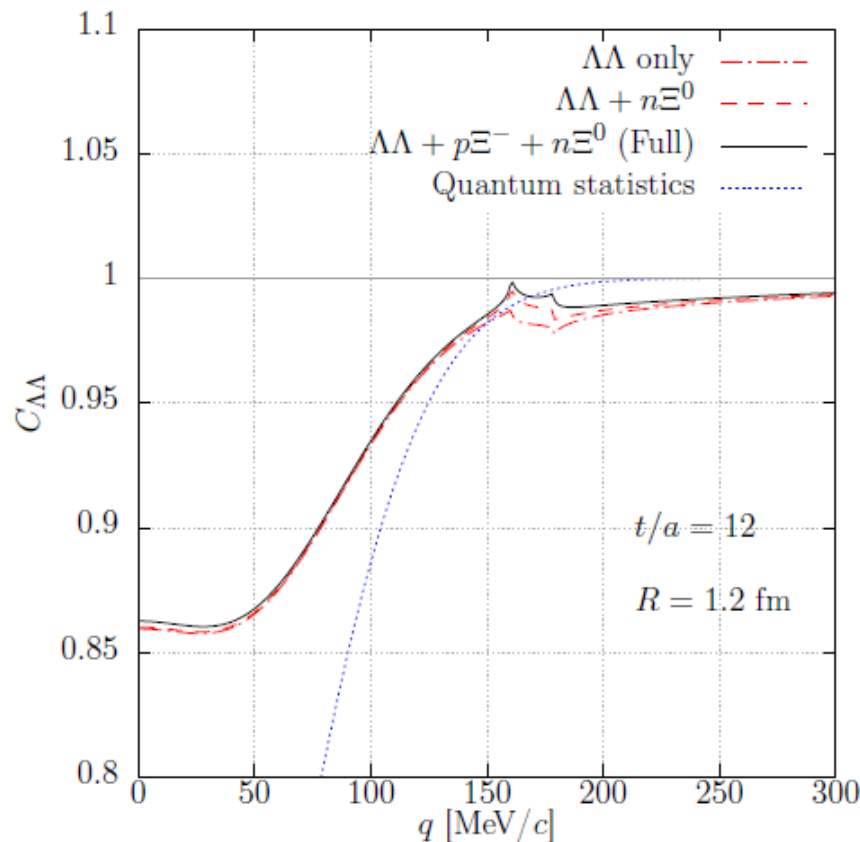


*Kamiya+(in prep.).
(w/ Lattice BB pot. at phys. m_q
and CC effects with $\Lambda\Lambda$)*

$\Lambda\Lambda$ correlation function

■ $\Lambda\Lambda$ correlation function

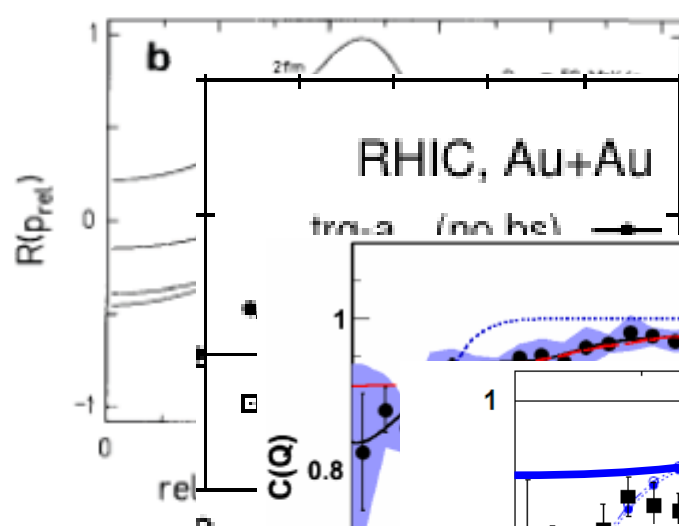
- Calculated CF agrees with ALICE data.
- $N\Xi$ source effect is visible only around threshold.



Kamiya+, in prep.; Acharya+(ALICE), Nature ('20)

Comparison with other results

Lambda-correlation with resonance



C. Greiner, B. Muller, PLB219('89)199.

(Assumed $\Lambda\Lambda$ resonance)

AO, Hirata, Nara, Shinmura, Akaishi, NPA670('00)297c

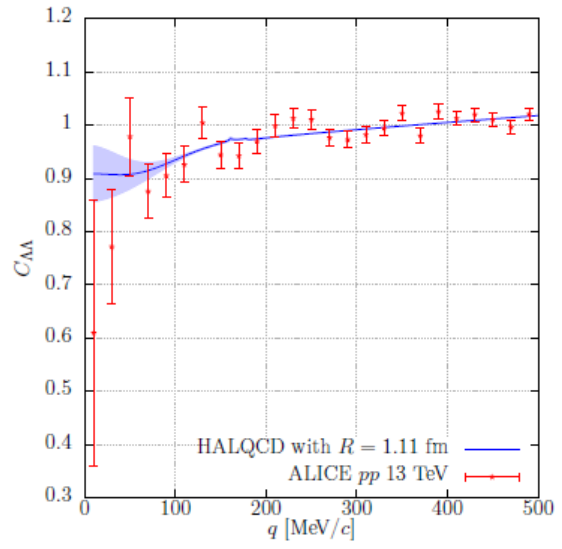
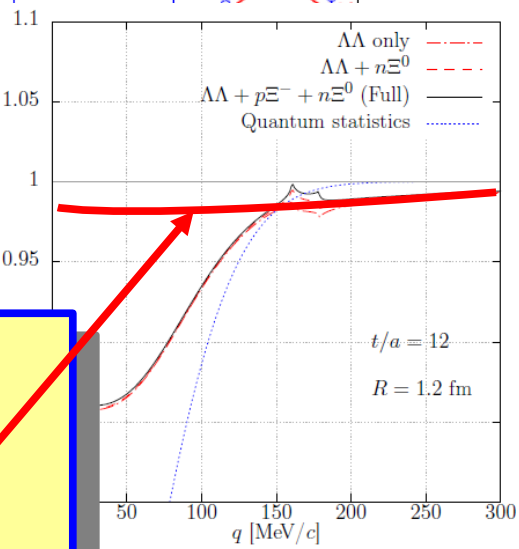
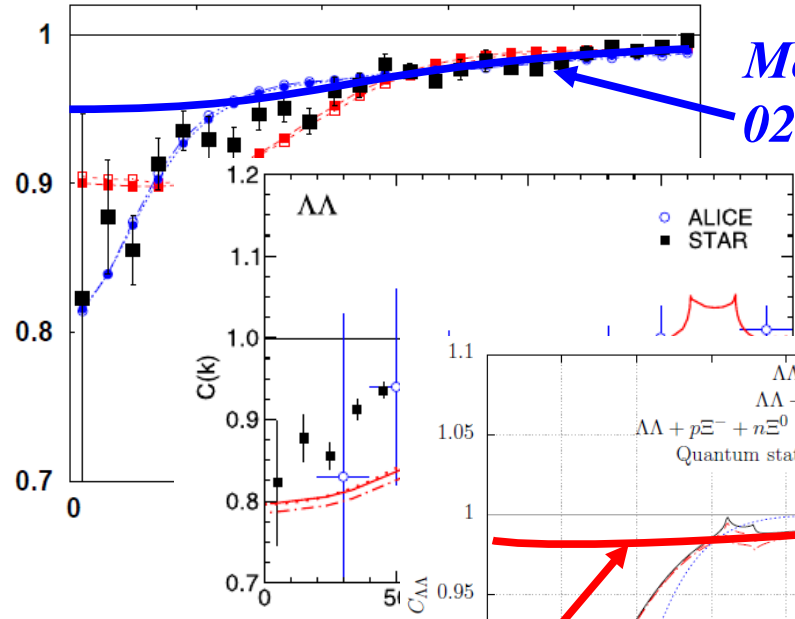
(Before NAGARA, interaction was too strong.)

Adamczyk+[STAR], PRL114('15)022301

(Residual source $R \sim 0.5$ fm was assumed.)

Morita, Furumoto, AO, PRC91('15)024916. (Res. Source + flow)

J. Haidenbauer, NPA981('19)1. (Larger cusp ?)



*Kamiya+(in prep.).
(Smaller cusp than χ EFT.
CC simulates res. source !)*

Summary

- **Correlation functions (CFs) can be utilized to constrain / examine the hadron-hadron interactions, as well as to deduce the pair is to be bound or not to be bound.**
- **We have calculated $p\Xi^-$ and $\Lambda\Lambda$ correlation functions by using lattice $N\Xi-\Lambda\Lambda$ coupled-channel (CC) potential with effects of CC, Coulomb potential, threshold difference.**
- **Recent ALICE data are consistent with the HAL QCD potential.**
 - **CC effect of $N\Xi-\Lambda\Lambda$ is not very big.**
 - **$\Lambda\Lambda$ CF shows cusps at $N\Xi$ thresholds, and CC effect simulates small residual source assumed to interpret heavy-ion data by STAR.**
 - **The data may suggest the existence of the H dibaryon state as the virtual pole around the $N\Xi$ threshold.**
- **Pole position is sensitive to the detail of interaction, so the source size dependence would be desirable to measure.**

Homeworks

- Pole position is sensitive to the detail of interaction, so the source size dependence would be desirable to measure.
 - Larger size (STAR seems to be analyzing $p\Xi^-$ correlation function data, private communication with Neha Shar)
 - Smaller size (J-PARC-E42 is measuring $p\pi\Lambda$, $\Lambda\Lambda$, $p\Xi^-$ invariant mass distribution, and effective source size is smaller than pp)
- Analytic Lednicky-Lyuboshits type formula should be developed / examined to confirm the pole position of the H dibaryon.
 - Explaining the data with a given potential is not enough !
Method to extract scattering parameters directly from data is desired.
- Further studies including charm, three-body, deuteron will enrich hadron interactions from femtoscopy.

Thank you for attention !

Coauthors of Y. Kamiya et al. ($p\bar{E}^-$), in prep.

Y. Kamiya



K. Sasaki



T. Fukui



T. Hatsuda



T. Hyodo



K. Morita



K. Ogata



AO



*Y. Kamiya, K. Sasaki, T. Fukui, T. Hatsuda,
T. Hyodo, K. Morita, K. Ogata, AO, in prep.*

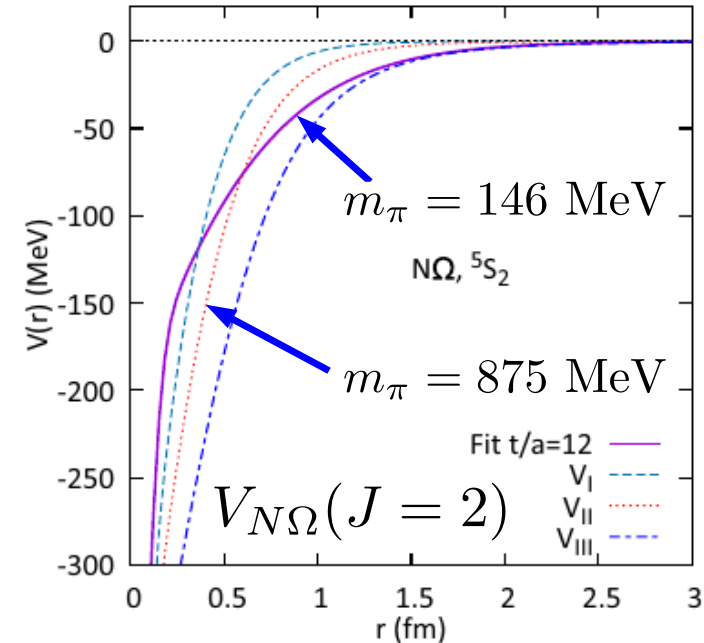
Modern Hadron-Hadron Interactions

Lattice QCD hh potential

- V_{hh} is obtained from the Schrödinger eq. for the Nambu-Bethe-Salpeter (NBS) amplitude.

N. Ishii, S. Aoki, T. Hatsuda, PRL99('07)022001.

→ $\Omega\Omega$, $N\Omega$, $\Lambda\Lambda$ - $N\Xi$ potentials at phys. quark mass are published



Chiral EFT / Chiral SU(3) dynamics

- V_{hh} at low E. can be expanded systematically in powers of Q/Λ_χ .

S. Weinberg ('79); R. Machleidt, F. Sammarruca ('16);

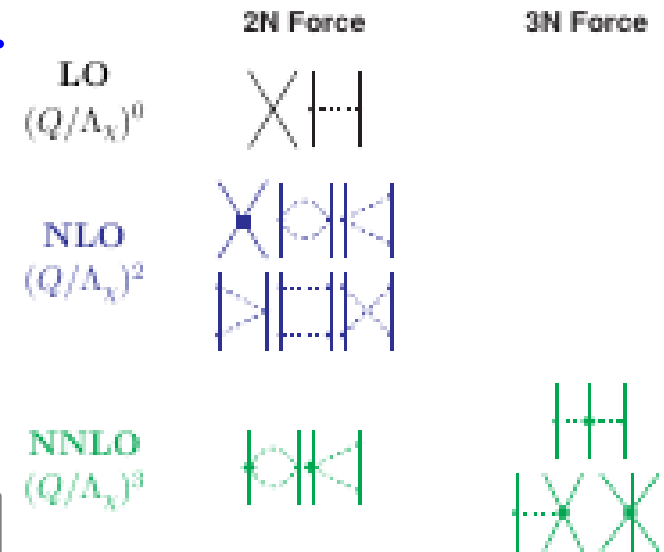
Y. Ikeda, T. Hyodo, W. Weise ('12).

→ NN , NY , YY , KN - $\pi\Sigma$ - $\pi\Lambda$, ...

Quark cluster models,

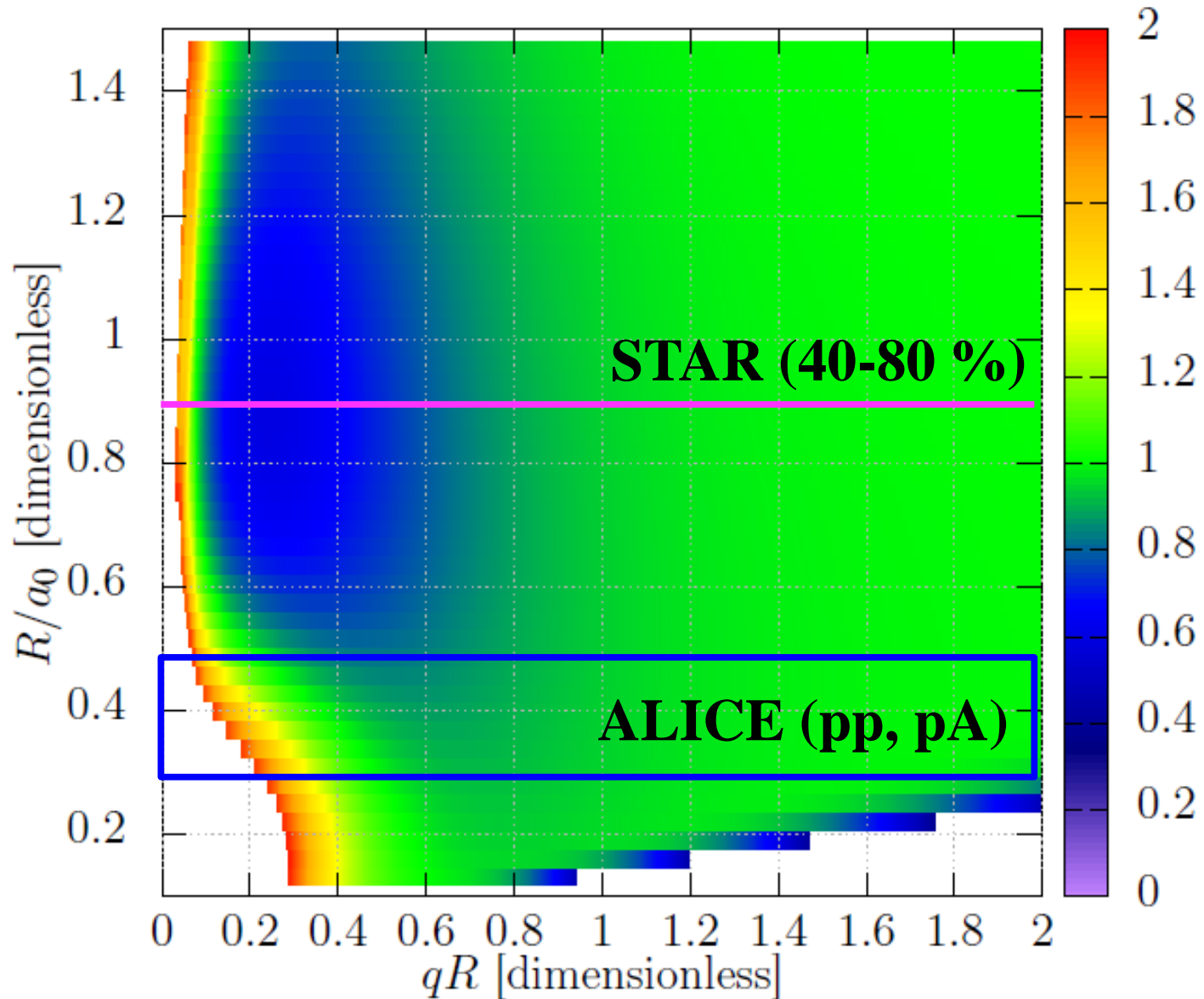
Meson exchange models,

More phenomenological models



Let us examine modern hh interactions !

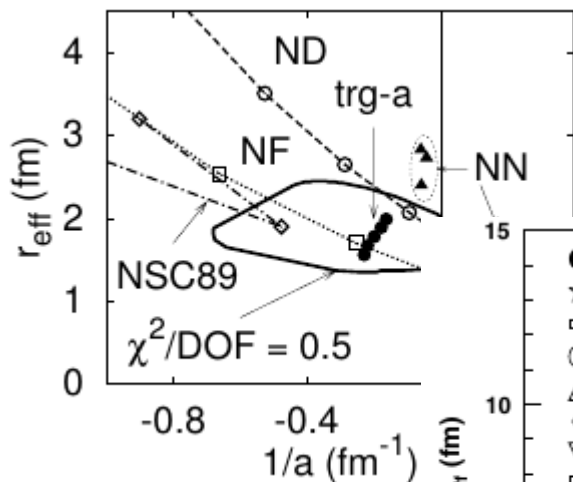
Correlation Function with Gaussian source



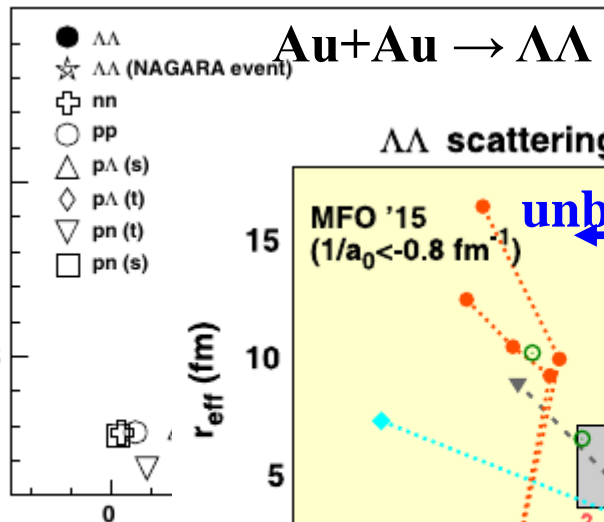
$N\Omega$ potential ($J=2$, HAL QCD, $a = 3.4$ fm) + Coulomb

$\Lambda\Lambda$ correlation and $\Lambda\Lambda$ interaction

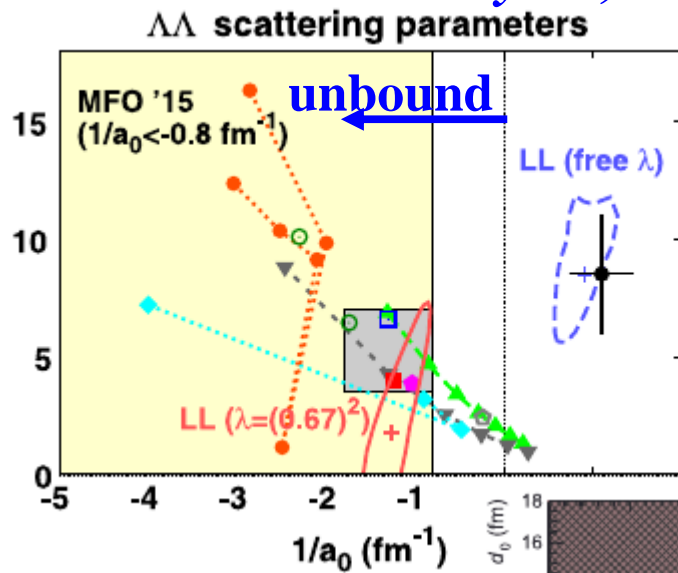
It is unlikely that $\Lambda\Lambda$ bound state exists.



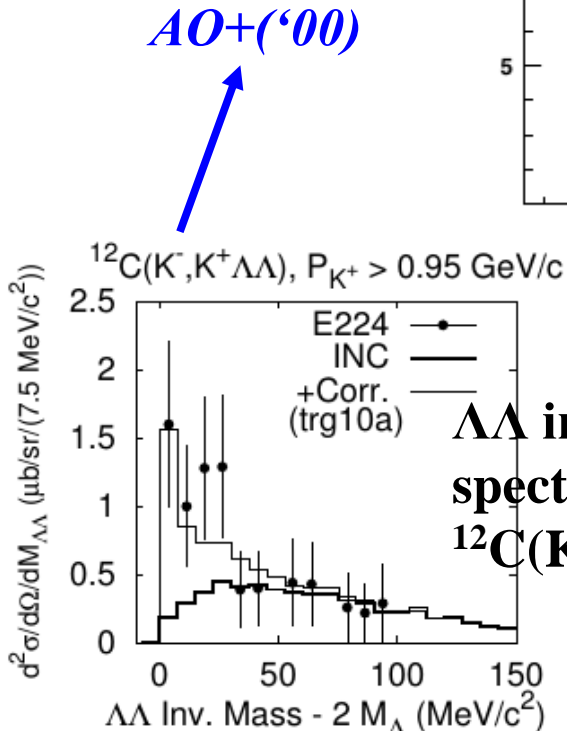
*L. Adamczyk+[STAR],
PRL114('15)022301*



*AO, K. Morita, K. Miyahara,
T. Hyodo, NPA954('16)294*



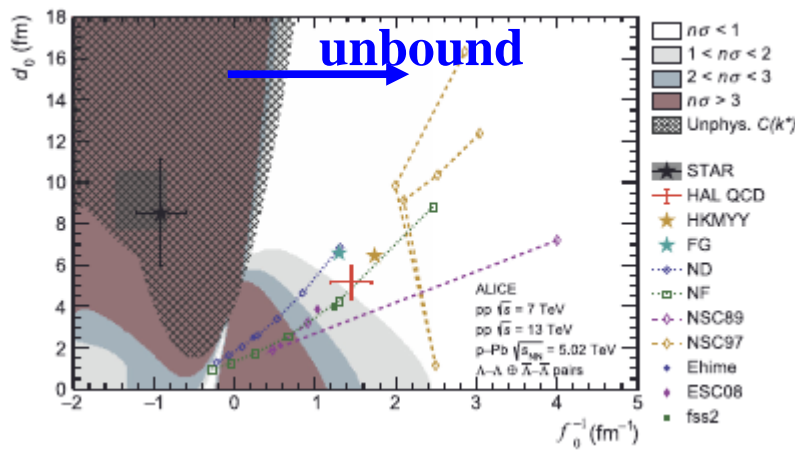
$\delta \sim -a_0 q$



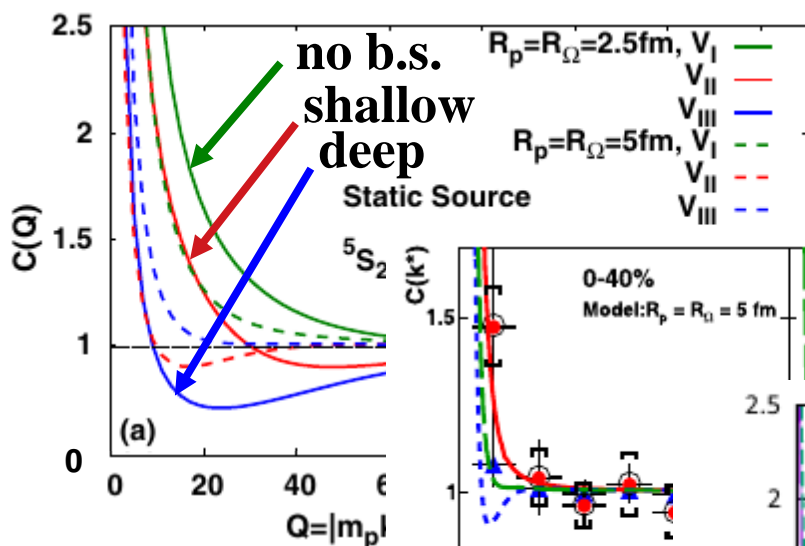
$\Lambda\Lambda$ inv. mass spectrum from $^{12}\text{C}(K^-, K^+\Lambda\Lambda)$

*S. Acharya+[ALICE],
PLB797('19)134822*

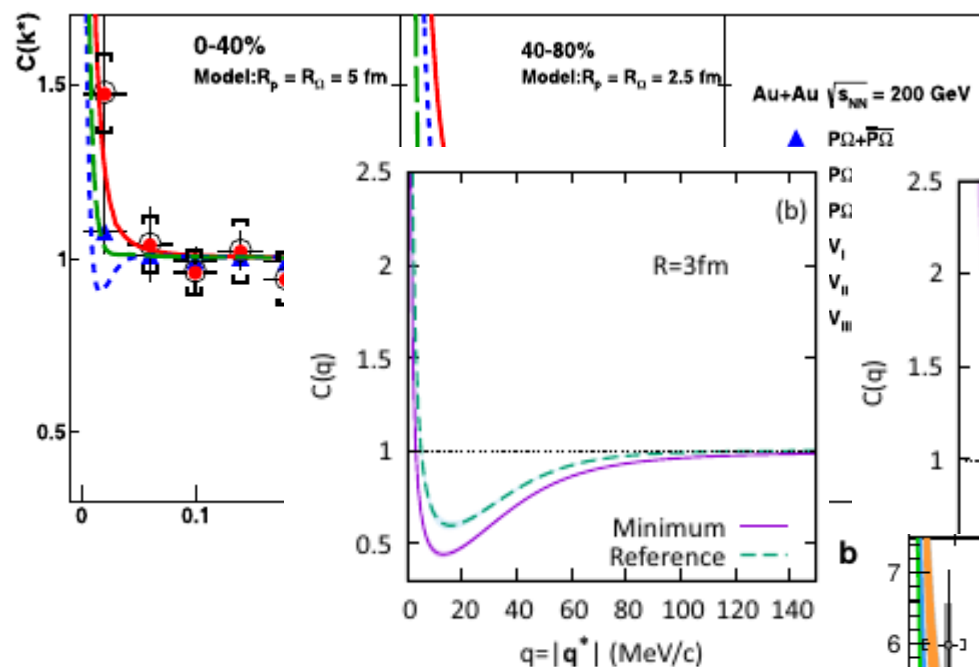
$\delta \sim +a_0 q$



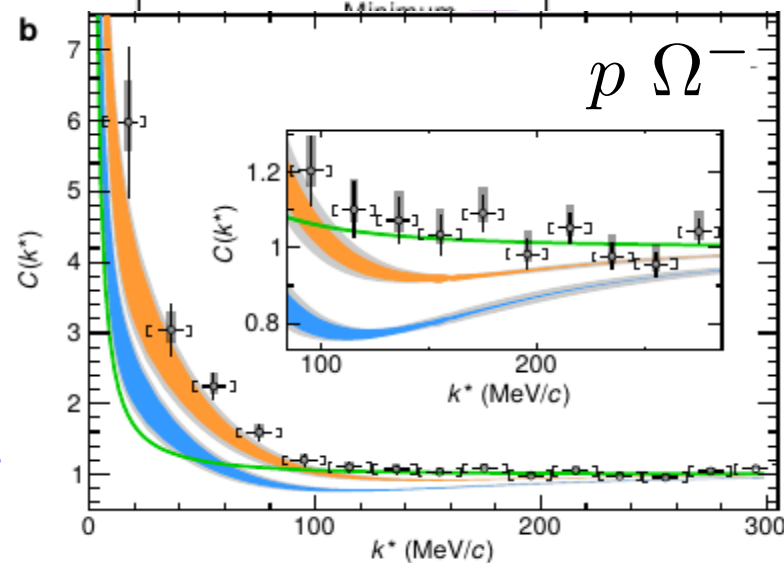
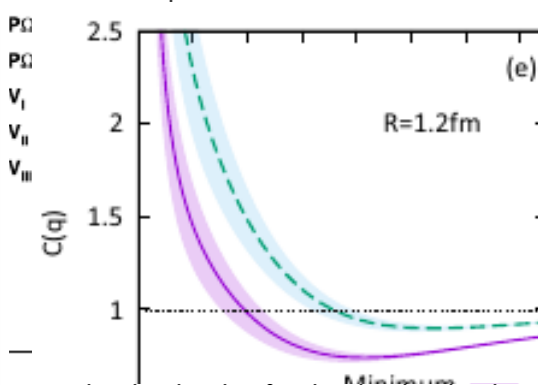
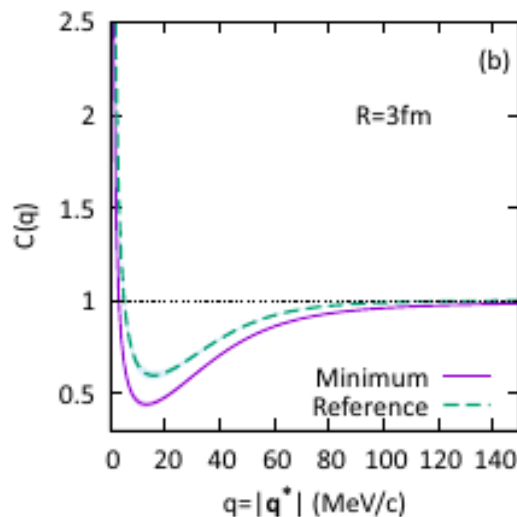
$p\Omega^-$ correlation



*K. Morita, AO, F. Etminan,
T. Hatsuda, PRC94('16)031901(R)
(w/ Lattice potential with heavier quark mass)*



*J. Adam+[STAR],
PLB790('19)490.*

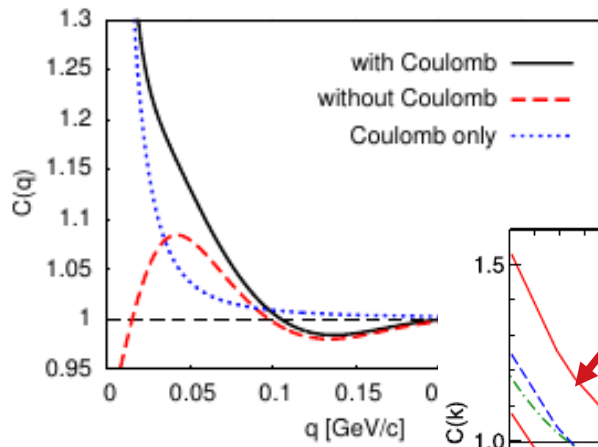


*K. Morita, S. Gongyo, T. Hatsuda,
T. Hyodo, Y. Kamiya, AO,
PRC 101('20)015201. (w/ Lattice
potential at physical quark mass)*

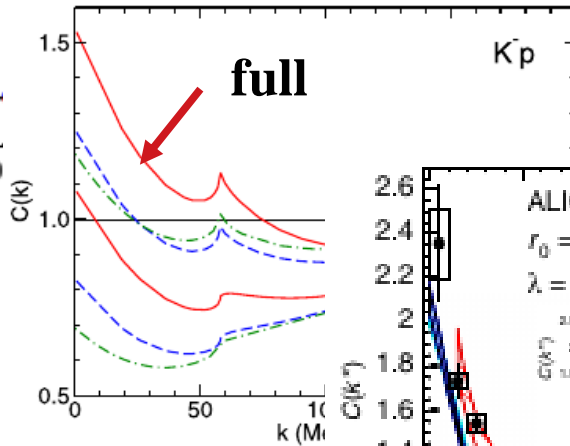
*S. Acharya+[ALICE],
2005.11495 [nucl-ex]
(pp 13 TeV)*

Bound state ?

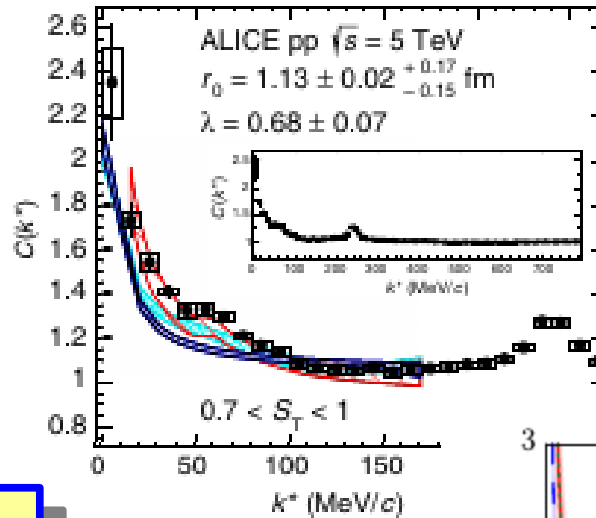
pK^- correlation



S. Cho+ [ExHIC], PPNP95('17)279.
(Insufficient coupled-channel effects)



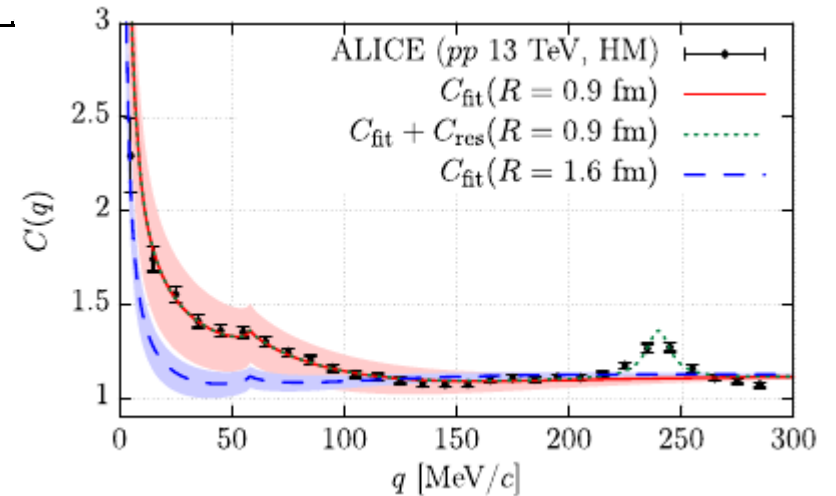
J. Haidenbauer, NPA981('19)1.
(w/ CC effects, w/o Coulomb)



S. Acharya+[ALICE],
PRL124('20)092301

- ◆ $K^+p \oplus K^+p$
- Coulomb
- Coulomb+Strong (Kyoto Model)
- Coulomb+Strong (Jülich Model)

Source size dep. shows interesting feature.

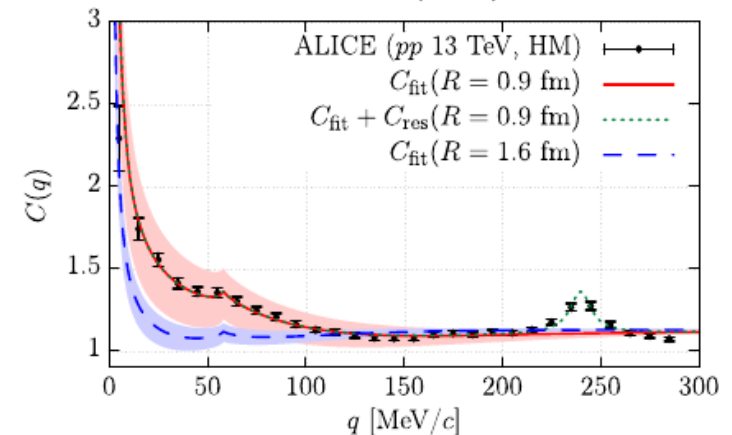
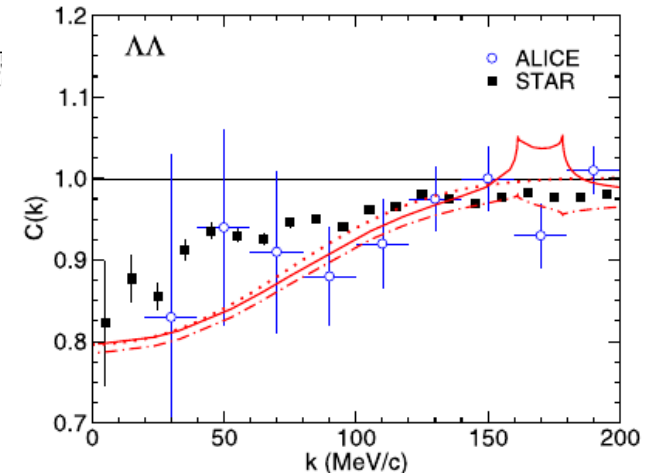
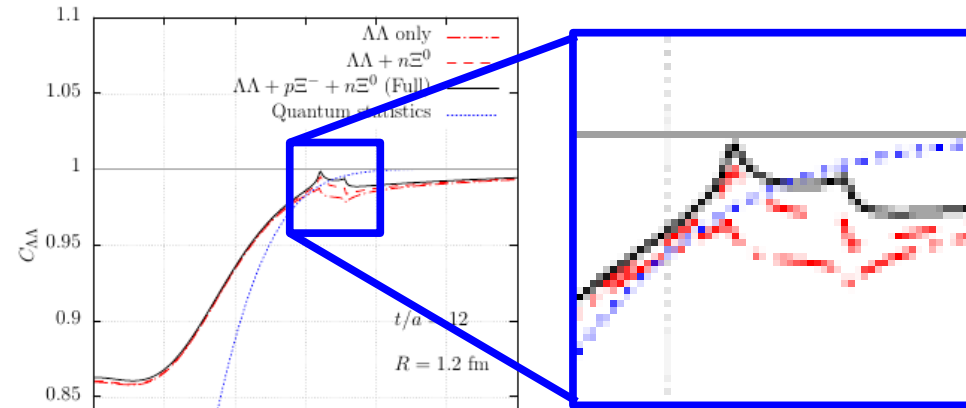


Y. Kamiya, T. Hyodo, K. Morita, AO,
W. Weise, PRL124('20)132501.
(Chiral SU(3) dynamics)

Other bound states ?

■ $\Lambda\Lambda$ - $N\Xi$

- $C(q)$ in AA(RHIC) and pp(LHC) are similar (No b.s. below $\Lambda\Lambda$).
- LQCD predicts a virtual pole near $N\Xi$ threshold, which can be detected as the cusp in $C(q)$. NLO(600) potential predicts the same. (The fate of H particle)
K. Sasaki+[HAL QCD], NPA998('20)121737;
Y. Kamiya+, in prep.; Haidenbauer('19).



■ KN

- $\Lambda(1405)$ is believed to be the bound state of KN, and “dip” is expected at $R \sim a$.
- However, Coulomb and coupled-channel effects modify the dip-like behavior.

Kamiya+ ('20).

Lednický-Lyuboshits (LL) model

■ Lednický-Lyuboshits analytic model

- **Asymp. w.f. + Eff. range corr. + $\psi^{(-)} = [\psi^{(+)}]^*$**

$$\psi_0(r) \rightarrow \psi_{\text{asy}}(r) = \frac{e^{-i\delta}}{qr} \sin(qr + \delta) = \mathcal{S}^{-1} \left[\frac{\sin qr}{qr} + f(q) \frac{e^{iqr}}{r} \right]$$

$$\begin{aligned} \Delta C_{\text{LL}}(q) &= \int dr S_{12}(r) (|\psi_{\text{asy}}(r)|^2 - |j_0(qr)|^2) \\ &= \frac{|f(q)|^2}{2R^2} F_3\left(\frac{r_{\text{eff}}}{R}\right) + \frac{2\text{Re}f(q)}{\sqrt{\pi}R} F_1(x) - \frac{\text{Im}f(q)}{R} F_2(x) \end{aligned}$$

($x = 2qR$, $R = \text{Gaussian size}$, F_1, F_2, F_3 : Known functions)

■ Phase shifts

$$q \cot \delta = -\frac{1}{a_0} + \frac{1}{2} r_{\text{eff}} q^2 + \mathcal{O}(q^4) \rightarrow \delta \simeq -a_0 q + \mathcal{O}(q^3)$$

$$\sin(qr + \delta) \simeq \sin(q(r - a_0) + \dots) \quad \text{Node at } r \sim a_0 \text{ for small } q$$

$C(q)$ in the low momentum limit

- Correlation function at small q (and $r = 0$) $\rightarrow F = 1, F = 0, F = 1$

$$\Delta C_{LL}(q) \rightarrow \frac{|f(0)|^2}{2R^2} + \frac{2\text{Re}f(0)}{\sqrt{\pi}R} \quad (q \rightarrow 0)$$

$$f(q) = (q \cot \delta - iq)^{-1} \simeq \left(-\frac{1}{a_0} + \frac{1}{2}r_{\text{eff}}q^2 - iq \right)^{-1} \rightarrow -a_0$$

$$C_{LL}(q \rightarrow 0) = 1 + \frac{a_0^2}{2R^2} - \frac{2a_0}{\sqrt{\pi}R} = 1 - \frac{2}{\pi} + \frac{1}{2} \left(\frac{a_0}{R} - \frac{2}{\sqrt{\pi}} \right)^2$$

$$1 - 2/\pi \simeq 0.36, \quad \sqrt{\pi}/2 \simeq 0.89$$

$C(q \rightarrow 0)$ takes a minimum of 0.36 at $R/a = 0.89$ in the LL model.

Correlation Function with Coupled-Channels Effects

J. Haidenbauer, NPA 981('19)1; R. Lednicky, V. V. Lyuboshits, V. L. Lyuboshits, Phys. At. Nucl. 61('98)2950.

- **Single channel, w/o Coulomb (non-identical pair)**

$$C(\mathbf{q}) = \underline{1} + \int d\mathbf{r} S(\mathbf{r}) \left[\underline{|\chi^{(-)}(r, q)|^2} - \underline{|j_0(qr)|^2} \right]$$

- **Single channel, w/ Coulomb**

$$C(\mathbf{q}) = \int d\mathbf{r} S(\mathbf{r}) \left[\underline{|\varphi^{C,\text{full}}(\mathbf{q}, \mathbf{r})|^2} + \underline{|\chi^{C,(-)}(r, q)|^2} - \underline{|j_0^C(qr)|^2} \right]$$

Full free

s-wave w.f.

s-wave

Coulomb w.f.

with Coul.

Coul. w.f.

- **Coupled channel, w/ Coulomb**

$$C_i(\mathbf{q}) = \int d\mathbf{r} S_i(\mathbf{r}) \left[\underline{|\varphi^{C,\text{full}}(\mathbf{q}, \mathbf{r})|^2} + \underline{|\chi_i^{C,(-)}(r, q)|^2} - \underline{|j_0^C(qr)|^2} \right] \\ + \sum_{j \neq i} \omega_j \int d\mathbf{r} S_j(\mathbf{r}) \underline{|\chi_j^{C,(-)}(r, q)|^2}$$

s-wave w.f.

in j-th channel

Outgoing B.C. in the i-th channel, ω = Source weight ($\omega = 1$)