Femtoscopy and scattering of Λ - α with Λ potential from chiral EFT

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- Introduction hyperon puzzle and Λ potential at high densities
- **\land** Λ - α potential from chiral effective field theory
- Λ-α correlation function
- **Summary**

Work in progress



Hyperon puzzle

- Hyperonic matter EOS with empirical Hyperon-Nucleon (YN) interactions cannot support 2 solar mass neutron stars
- Proposed solutions
 - YN interaction details ? YNN 3-body int. ? Many-body theories ? Gradual transition to quark matter ? Modified gravity ?
 - A plausible answer = repulsive Λ potential at high densities caused by YNN 3-body int.

Chiral EFT: Gerstung, Kaiser, Weise (2001.10563), Kohno (1802.05388)



A. Ohnishi @ 3rd J-PARC HEF-ex WS, Mar. 14, 2023 2

Verification of repulsive U_{Λ} at high densities

- Directed flow of Λ in heavy-ion collisions Y. Nara, A. Jinno, K. Murase, AO (2208.01297)
 - U_Λ from chiral EFT w/ 3-body force (Chi3) reasonably explains the directed flow slope.
 - Other softer potentials at high densities also explains the data (other effects are also important).

Skyrme Hartree-Fock calculation using Chiral EFT U_A w/ 3BF

A. Jinno, K. Murase, Y. Nara, AO, in prep. [Talk by Jinno, Thu.]³⁰

Explains the Λ binding energies of hypernuclei comparably well to empirical SHF potentials.
D. E. Lanskoy, Y. Yamamoto, PRC55, 2330 (1997);
N. Guleria, S. K. Dhiman, R. Shyam, NPA886, 71 (2012).



Ζ

Can we distinguish the density dep. of U_A ? \rightarrow How about A-nucleus femtoscopy or scattering ?

Λ -a femtoscopy or scattering in HEFex









Λ potential (U_{Λ})

Skyrme Hartree-Fock equation

$$\begin{bmatrix} -\boldsymbol{\nabla} \cdot \left(\frac{\hbar^2}{2m_B^*(\boldsymbol{r})}\right) \boldsymbol{\nabla} + U_B^{\text{local}}(\boldsymbol{r}) - i\boldsymbol{W}_B(\boldsymbol{r}) \cdot (\boldsymbol{\nabla} \times \boldsymbol{\sigma}) \end{bmatrix} \phi_i^B(\boldsymbol{r}) = \varepsilon_i \ \phi_i^B(\boldsymbol{r}) ,$$
$$U_{\Lambda}^{\text{local}} = a_1^{\Lambda} \rho_N + a_2^{\Lambda} \tau_N - a_3^{\Lambda} \triangle \rho_N + a_4^{\Lambda} \rho_N^{4/3} + a_5^{\Lambda} \rho_N^{5/3} ,$$
$$\frac{\hbar^2}{2m_{\Lambda}^*} = \frac{\hbar^2}{2m_{\Lambda}} + a_2^{\Lambda} \rho_N , \ \rho_N = \sum_{i \in N} |\phi_i|^2 , \ \tau_N = \sum_{i \in N} |\boldsymbol{\nabla} \phi_i|^2$$

Λ-α potential

$$\begin{bmatrix} -\frac{\hbar^2 \nabla^2}{2\mu} + U(\boldsymbol{r}) \end{bmatrix} \phi(\boldsymbol{r}) = \varepsilon_{\Lambda} \ \phi(\boldsymbol{r}) ,$$
$$U = a_1^{\Lambda} \rho_N + a_2^{\Lambda} \tau_N - a_3^{\Lambda} \triangle \rho_N + a_4^{\Lambda} \rho_N^{4/3} + a_5^{\Lambda} \rho_N^{5/3} - a_2^{\Lambda} \nabla \cdot \rho_N \nabla$$

Nucleon densities are measured from the center of mass of α
μ = reduced mass of Λ-α



Wave function of α particle

A simple gaussian wave function of α

$$\psi = \prod_{i=1}^{A} \phi(\boldsymbol{r}_i), \ \phi(\boldsymbol{r}) = \left(\frac{2\nu}{\pi}\right)^{3/4} \exp[-\nu \boldsymbol{r}^2].$$

Naive estimate of densities

$$\rho(\boldsymbol{r}) = A\left(\frac{2\nu}{\pi}\right)^{3/2} \exp[-2\nu\boldsymbol{r}^2] \ (A=4) \,,$$

$$\tau(\boldsymbol{r}) = \sum_{i} |\boldsymbol{\nabla}\phi_{i}|^{2} = 4\nu^{2}\boldsymbol{r}^{2}\rho(\boldsymbol{r}), \ \Delta\rho(\boldsymbol{r}) = \frac{1}{r}\frac{\partial^{2}}{\partial^{2}r}r\rho(r) = -4\nu(3-4\nu r^{2})\rho(\boldsymbol{r})$$

Densities measured from the center-of-mass

$$\rho(\mathbf{r}) = \int d\mathbf{r}_1 \cdots d\mathbf{r}_A |\psi(\mathbf{r}_1, \cdots, \mathbf{r}_A)|^2 \sum_i \delta(\mathbf{r}_i - \mathbf{r}_G - \mathbf{r})$$

$$=A\left(\frac{2\nu_{c}}{\pi}\right)^{3/2}\exp[-2\nu_{c}\boldsymbol{r}^{2}] \left[\nu_{c}=\nu/(1-1/A)\right],$$

$$\Delta \rho(\boldsymbol{r}) = -4\nu_c (3 - 4\nu_c r^2)\rho(\boldsymbol{r}), \ \tau(\boldsymbol{r}) = \rho(\boldsymbol{r}) \left(4\nu^2 \boldsymbol{r}^2 + \frac{3\nu}{A}\right)$$

Width parameter from rmsr=1.67824(83) fm J.J. Krauth et al., Nature 589, 527 (2021); $\nu = \frac{9}{16\langle r^2 \rangle_A} = 0.20 \text{ fm}^{-2}$



Densities should include nucleon spread.

Λ potential parameters

Parameters in Empirical SHF for Λ and in Chiral EFT

	a1	a2	a3	a4	a5	BE(5HeΛ)	BE(pot.)
LY4	-500.89	16.00	20.000	548.411	0	3.63184	2.70
HPL2	-302.76	23.72	29.853	514.263	0	2.48189	_
Chi2	-352.21	39.35	52.158	-356.95	1000.81	3.61234	2.46
Chi3	-388.28	47.28	36.558	-405.67	1256.75	4.78802	2.83

⁵ He binding energy 3.12 ± 0.02 MeV, Juric et al., NPB52 (1973) 1;

LY4 from D. E. Lanskoy, Y. Yamamoto, PRC55, 2330 (1997); HPL2 from N. Guleria, S. K. Dhiman, R. Shyam, NPA886, 71 (2012).

Chi2/Chi3 fit results of Chiral EFT w/o and w/ 3BF Gerstung, Kaiser, Weise EPJA56('20),175(2001.1056 Kohno, PRC97('18), 035206 (1802.05388).

- We ignore LS potential for Λ
- a3 parameter in Chi2 and Chi3 is tuned to fit ¹³ C binding energy.

Jinno, Murase, Nara, AO, in prep. Jinno's talk on Thu.





Λ -a potential

- Significant difference in the potential shape (local part)
 - LY4 shows Woods-Saxon type Λ - α potential (volume type)
 - Chi3 shows a dip at r ~ 1.2 fm and central repulsion (surface type)
 - Potential shape reflects the density dependence of U_{Λ}



 Λ - α potential (local part)

Λ - α scattering phase shift



 $(a_0, r_{\text{eff}}) = (4.13, 1.72) \text{fm} (\text{Chi}3), (3.94, 1.76) \text{fm} (\text{LY}4)$



Λ - α correlation function



Femtoscopic study of hadron-hadron interaction

- How can we study interactions between short-lived particles $? \rightarrow$ Femtoscopy !
- Correlation function (CF)
 - Koonin-Pratt formula

Koonin('77), Pratt+('86), Lednicky+('82)



source fn. relative w.f. $C(\boldsymbol{p}_{1}, \boldsymbol{p}_{2}) = \frac{N_{12}(\boldsymbol{p}_{1}, \boldsymbol{p}_{2})}{N_{1}(\boldsymbol{p}_{1})N_{2}(\boldsymbol{p}_{2})} \simeq \int d\boldsymbol{r} S_{12}(\boldsymbol{r}) |\varphi_{\boldsymbol{q}}(\boldsymbol{r})|^{2}$

- Source size from quantum stat. + CF (Femtoscopy) Hanbury Brown & Twiss ('56); Goldhaber, Goldhaber, Lee, Pais ('60)
- Hadron-hadron interaction from source size + CF
 - CF of non-identical pair from Gaussian source R. Lednicky, V. L. Lyuboshits ('82); K. Morita, T. Furumoto, AO ('15)

$$C(\boldsymbol{q}) = 1 + \int d\boldsymbol{r} S(r) \left\{ |\varphi_0(r)|^2 - |j_0(qr)|^2 \right\} \quad (\varphi_0 = \text{s-wave w.f.})$$

CF shows how much $|\varphi|^2$ is enhanced $\rightarrow V_{\mu\nu}$ effects !



Measured Correlation Functions (examples)





Theoretical femtoscopic study of hh int. (examples)





Haidenbauer(1808.05049), *Morita*+(1408.6682)





Y.Kamiya, K.Sasaki, et al., (2108.09644)

Ξ-p R = 1.182 fm

Coulomb Coulomb

100

150

p-Ξ ⊕ p-Ξ* (ALICE p-Pb collisions at 13 TeV)

 $\pm p$

200

3.0

2.5

20

1.5

1.0

50

C (K)



(covariant χEFT, S=-2) *Mrówczyński*, *Słón (1904.08320, K⁻d)*, Haidenbauer (2005.05012, Ad), Etminan, Firoozabadi (1908.11484, Ωd), L.-S. Geng K.Ogata+ (Ξ⁻ d,2103.00100)

Z.-W. Liu, K.-W. Li, (2201.04997)





Λ -a correlation function

Comparison of Λ-α correlation functions from empirical (LY4) and chiral (Chi3) potentials for gaussian source with R=1 and 3 fm.





Λ -a correlation function

- For a large source (R=3 fm), LY4 and Chi3 give almost the same C(q).

 - Lednicky-Lyuboshitz formula based on asymptotic w.f. works well.
- For a small source (R=1 fm),

Chi3 gives a slightly smaller C(q) at low momenta.

- Central repulsion pushes out the w.f. to the outer region (r>1 fm).
- In all cases, the bound state $({}^{5}_{\Lambda}He)$ causes a dip in C(q).
 - The bound state eats the yield at low momenta.





Summary

- The density dependence of Λ , $U_{\Lambda}(\rho)$, is important to understand and/or to solve the hyperon puzzle.
- Repulsive U_A at high densities from chiral EFT with 3-body force effects seems to be consistent with

the flow data of Λ from heavy-ion collisions

and the hypernuclear binding energies. Gerstung, Kaiser, Weise EPJA56('20),175(2001.10563); Kohno, PRC97('18), 035206 (1802.05388). Nara, Jinno, Murase, AO (PRC106 ('22), 044902 (2208.01297); Jinno, Murase, Nara, AO, in prep. However, U_{Λ} with different density dependence can also explain the data.

- A-α correlation function may be helpful to distinguish the density dependence, if C(q) is obtained at the precision of a few %.
 - Large source → scattering parameters (STAR BES II?)
 - Small source \rightarrow difference in the potential shape (J-PARC?)



Summary (cont.)

- There are many problems...
 - Feasible in experiment ?
 - α production from a small source (R=1 fm \rightarrow pp or pA)?
 - Does pA collision at $\sqrt{s_{NN}} < 5$ GeV lead to "chaotic" source ?
 - Is a few % precision achievable in experiment ?
 - Other effects may modify C(q) E.g. m*=m approximation leads to different shape of C(q).
 - Other "smoking gun" observable?



