

KEK informal seminar (June 6th, 2002)
seminar at Ochanomizu University (June 17th, 2002)

GAUGE THEORETICAL CONSTRUCTION OF NON-COMPACT CALABI-YAU MANIFOLDS

Annals of Physics **296** (2002) 347, hep-th/0110216

hep-th/0202064

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INTRODUCTION

We want to one-loop (or higher-loops) finite Field Theory

also

We want to obtain **Calabi-Yau metrics**

in order to give $\mathcal{N} = 1$ Supergravity in Four-dimensions
from String Theory (Candelas-Horowitz-Strominger-Witten, 1985)

↓

$\mathcal{N} = 2$ Supersymmetric Nonlinear Sigma Models in Two-dimensions

(= worldsheet in string theory)

with Ricci-flat Kähler (Calabi-Yau) target space

(= extra six-dimensions in superstring theory)

BUT...

It is difficult to obtain the explicit expression

of the metric of **compact** Calabi-Yau

Recently,

singular Calabi-Yau n -folds have been studied

in order to understand

the **nonperturbative effect** of

$\mathcal{N} = 1$ super-Yang–Mills theories in four-dimensions

[appearance of the new massless spectrum,
dynamical generations of the superpotential,
gauge/gravity dual, large- N duality, ...]

by

flop, conifold transition,

D-branes wrapped around the SUSY-cycles, etc.

Near the singular point of such Calabi-Yau,

this Calabi-Yau looks like **non-compact**

and

we can obtain this metric **very easily!**

Now we obtain the metrics of **non-compact Calabi-Yau n -folds**

\Rightarrow generalizations of the conifold

$\left(\begin{array}{l} \text{conifold : topologically } \mathbb{R}_+ \times \mathbf{S}^2 \times \mathbf{S}^3 \\ \text{appears in } \textit{conifold transition}, \textit{ Vafa's large-}N \textit{ dual, etc.} \end{array} \right)$

Our manifolds are constructed as

$$\mathbb{C}^1 \times G/H \simeq \mathbb{R}^1 \times G/H'$$

$$G/H = \text{Hermitian Symmetric Spaces (HSS)}, \quad H = H' \times U(1)$$

At large distance

the metrics tend to Ricci-flat “cones” over HSS.

But at the small distance

the “nodes” are **repaired by the HSS** and become smooth!

ORGANIZATION

Introduction

Review: Hermitian Symmetric Spaces

Canonical Line Bundle over Projective Space

Conifold

Gauge Theoretical Construction

Generalization

Summary

REVIEW: HERMITIAN SYMMETRIC SPACES

projective space: $\mathbb{C}\mathbf{P}^{N-1} = SU(N)/[SU(N-1) \times U(1)]$

$\vec{\phi} \in \mathbb{C}^N$: identification $\vec{\phi} \sim \lambda \vec{\phi}$, $\lambda \in \mathbb{C}^1$

$\vec{\phi}^T = (1, \varphi^i)$, ($i = 1, 2, \dots, N-1$): local coordinates

Kähler potential K :

$$K = c \log \vec{\phi}^\dagger \vec{\phi} = c \log \{1 + |\varphi^i|^2\}, \quad c = \text{constant}$$

quadric surface: $Q^{N-2} = SO(N)/[SO(N-2) \times U(1)]$

$\mathbb{C}\mathbf{P}^{N-1} + [\vec{\phi}^2 = 0]$:

$\vec{\phi}^T = (1, \varphi^i, -\frac{1}{2}(\varphi^i)^2)$, ($i = 1, 2, \dots, N-2$): local coordinates

$$K = c \log \vec{\phi}^\dagger \vec{\phi} = c \log \left\{1 + |\varphi^i|^2 + \frac{1}{4}(\varphi^i)^2(\varphi^{*j})^2\right\}$$

exceptional groups: $E_6/[SO(10) \times U(1)]$, $E_7/[E_6 \times U(1)]$

$\mathbb{C}\mathbf{P}^{26} + [\Gamma_{ijk}\phi^j\phi^k = 0]$, $\mathbb{C}\mathbf{P}^{55} + [d_{\alpha\beta\gamma\delta}\phi^\beta\phi^\gamma\phi^\delta = 0]$:

Γ_{ijk} : rank-3 symmetric tensor invariant under E_6

$d_{\alpha\beta\gamma\delta}$: rank-4 symmetric tensor invariant under E_7

$[E_6]$: $\vec{\phi}^T = (1, \varphi_\alpha, -\frac{1}{2\sqrt{2}}\varphi C\sigma_A^\dagger\varphi)$, ($\alpha = 1, 2, \dots, 16$; $A = 1, 2, \dots, 10$)

$$K = c \log \left\{1 + |\varphi_\alpha|^2 + \frac{1}{8}|\varphi C\sigma_A^\dagger\varphi|^2\right\}$$

$[E_7]$: $\vec{\phi}^T = (1, \varphi^i, \frac{1}{2}\Gamma_{ijk}\varphi^j\varphi^k, \frac{1}{6}\Gamma_{ijk}\varphi^i\varphi^j\varphi^k)$, ($i = 1, 2, \dots, 27$)

$$K = c \log \left\{1 + |\varphi^i|^2 + \frac{1}{4}|\Gamma_{ijk}\varphi^j\varphi^k|^2 + \frac{1}{36}|\Gamma_{ijk}\varphi^i\varphi^j\varphi^k|^2\right\}$$

Grassmannian: $G_{N,M} = U(N)/[U(N-M) \times U(M)]$

$\Phi : N \times M$ matrix, identification $\Phi \sim \Phi V$ [$V \in U(M)$]

$$\Phi = \begin{pmatrix} \mathbf{1}_M \\ \varphi_{Aa} \end{pmatrix}, (A = 1, 2, \dots, N - M; a = 1, 2, \dots, M)$$

Kähler potential K :

$$K = c \log \det \Phi^\dagger \Phi = c \log \det \{ \mathbf{1}_M + \varphi^\dagger \varphi \}$$

$Sp(N)/U(N)$

$G_{2N,N} + [\varphi^T - \varphi = 0]$:

$$\Phi = \begin{pmatrix} \mathbf{1}_N \\ \varphi_{ab} \end{pmatrix}, (1 \leq a \leq b \leq N)$$

$$K = c \log \det \Phi^\dagger \Phi = c \log \det \{ \mathbf{1}_N + \varphi^\dagger \varphi \}$$

$SO(2N)/U(N)$

$G_{2N,N} + [\varphi^T + \varphi = 0]$:

$$\Phi = \begin{pmatrix} \mathbf{1}_N \\ \varphi_{ab} \end{pmatrix}, (1 \leq a < b \leq N)$$

$$K = c \log \det \Phi^\dagger \Phi = c \log \det \{ \mathbf{1}_N + \varphi^\dagger \varphi \}$$

Hermitian symmetric spaces

	type	G/H	$\dim_{\mathbb{C}}(G/H)$	Kähler potential K
	AIII ₁	$\mathbb{C}\mathbf{P}^{N-1}$	$N - 1$	$c \log\{1 + \varphi^i ^2\}$
	AIII ₂	$G_{N,M}$	$M(N - M)$	$c \log \det\{\mathbf{1}_M + \varphi^\dagger \varphi\}$
∞	BDI	Q^{N-2}	$N - 2$	$c \log\{1 + \varphi^i ^2 + \frac{1}{4}(\varphi^i)^2(\varphi^{*j})^2\}$
	CI	$\frac{Sp(N)}{U(N)}$	$\frac{1}{2}N(N + 1)$	$c \log \det\{\mathbf{1}_N + \varphi^\dagger \varphi\}$
	DIII	$\frac{SO(2N)}{U(N)}$	$\frac{1}{2}N(N - 1)$	$c \log \det\{\mathbf{1}_N + \varphi^\dagger \varphi\}$
	EIII	$\frac{E_6}{SO(10) \times U(1)}$	16	$c \log\{1 + \varphi_\alpha ^2 + \frac{1}{8} \varphi C \sigma_A^\dagger \varphi ^2\}$
	EVII	$\frac{E_7}{E_6 \times U(1)}$	27	$c \log\{1 + \varphi^i ^2 + \frac{1}{4} \Gamma_{ijk} \varphi^j \varphi^k ^2 + \frac{1}{36} \Gamma_{ijk} \varphi^i \varphi^j \varphi^k ^2\}$

CANONICAL LINE BUNDLE OVER PROJECTIVE SPACE

the simplest construction:

$$\vec{\phi}^T \equiv \sigma(1, \varphi^i), \quad \sigma \in \mathbb{C}^1$$

Assumption: Kähler potential \mathcal{K} (non-compact manifold)

$$\mathcal{K} = \mathcal{K}(X), \quad X \equiv \log \vec{\phi}^\dagger \vec{\phi} = \log |\sigma|^2 + \Psi$$

X : non-compact direction σ \times compact Kähler potential Ψ ($= K|_{c=1}$)

$$\mathbb{C} \times \frac{SU(N)}{SU(N-1) \times U(1)}$$

Ricci-flatness condition

holomorphic coordinates: $\phi^\mu = \{\sigma, \varphi^i\}$

metric: $g_{\mu\nu^*} = \partial_\mu \partial_{\nu^*} \mathcal{K}(X)$

$$\text{Ricci tensor : } \mathcal{R}_{\mu\nu^*} = -\partial_\mu \partial_{\nu^*} \log \det g_{\kappa\lambda^*}$$

Ricci-flatness condition:

$$\mathcal{R}_{\mu\nu^*} = 0 \longrightarrow \begin{cases} \det g_{\mu\nu^*} = (\text{constant}) \times |F|^2 \\ F = \text{holomorphic function} \end{cases}$$

consider in $\sigma \neq 0$ region:

$$g_{\sigma\sigma^*} = \frac{d^2\mathcal{K}}{dX^2} \cdot \frac{\partial X}{\partial\sigma} \frac{\partial X}{\partial\sigma^*}, \quad g_{\sigma j^*} = \frac{d^2\mathcal{K}}{dX^2} \cdot \frac{\partial X}{\partial\sigma} \frac{\partial X}{\partial\varphi^{*j}},$$

$$g_{ij^*} = \frac{d^2\mathcal{K}}{dX^2} \cdot \frac{\partial X}{\partial\varphi^i} \frac{\partial X}{\partial\varphi^{*j}} + \frac{d\mathcal{K}}{dX} \cdot \frac{\partial^2 X}{\partial\varphi^i \partial\varphi^{*j}}$$

determinant:

$$\det g_{\mu\nu^*} = \frac{1}{|\sigma|^2} \frac{d^2\mathcal{K}}{dX^2} \cdot \det \left\{ \frac{d\mathcal{K}}{dX} \cdot \frac{\partial^2 X}{\partial\varphi^i \partial\varphi^{*j}} \right\}$$

$$= \frac{1}{|\sigma|^2} \frac{d^2\mathcal{K}}{dX^2} \left(\frac{d\mathcal{K}}{dX} \right)^{N-1} \cdot \det \tilde{g}_{ij^*} \quad \left(\partial_i \partial_{j^*} X = \partial_i \partial_{j^*} \Psi \equiv \tilde{g}_{ij^*} \right)$$

Ricci-flatness condition is a [Partial Differential Equation](#)

not soluble in general...

BUT fortunately,

\mathbf{CP}^{N-1} is [Einstein-Kähler](#)

The condition reduces to an [Ordinary Differential Equation \(ODE\)](#)!

$$-\partial_i \partial_{j^*} \log \det \tilde{g}_{kl^*} = \tilde{\mathcal{R}}_{ij^*} = \mathcal{C} \tilde{g}_{ij^*} = \mathcal{C} \partial_i \partial_{j^*} \Psi$$

$$\rightarrow \det \tilde{g}_{ij^*} \sim \exp(-\mathcal{C}\Psi) \sim \exp(-\mathcal{C}X)$$

(up to holomorphic functions)

[Ricci-flatness condition](#):

$$e^{-NX} \frac{d}{dX} \left(\frac{d\mathcal{K}}{dX} \right)^N \equiv (\text{constant})$$

$$\frac{d\mathcal{K}}{dX} = (\lambda e^{NX} + b)^{\frac{1}{N}}$$

λ : positive real parameter

b : integration constant, **very important parameter**

$b \neq 0$ case:

metric: ($ds^2 = g_{\sigma\sigma^*} d\sigma d\sigma^* + \dots$)

$$g_{\sigma\sigma^*} = \lambda (\lambda e^{NX} + b)^{\frac{1-N}{N}} e^{N\Psi} |\sigma|^{2N-2}$$

The metric is ill-defined (degenerate) at $\sigma = 0$

But the curvature is still **finite** in the $\sigma \rightarrow 0$ limit

The coordinates $\phi^\mu = \{\sigma, \varphi^i\}$ has a coordinate singularity at $\sigma = 0$

↓

$$\text{transformation : } \rho \equiv \frac{\sigma^N}{N}$$

metric after transformation:

$$g_{\rho\rho^*} = \lambda (\lambda e^{NX} + b)^{\frac{1-N}{N}} e^{N\Psi}$$

$$g_{\rho j^*} = \lambda N (\lambda e^{NX} + b)^{\frac{1-N}{N}} e^{N\Psi} \rho^* \partial_{j^*} \Psi$$

$$g_{ij^*} = \lambda N^2 (\lambda e^{NX} + b)^{\frac{1-N}{N}} e^{N\Psi} |\rho|^2 \partial_i \Psi \partial_{j^*} \Psi + (\lambda e^{NX} + b)^{\frac{1}{N}} \partial_i \partial_{j^*} \Psi$$

metric of $\rho = 0$ ($d\rho = 0$) sub-manifold

$$g_{ij^*} \Big|_{\rho=0} = b^{\frac{1}{N}} \partial_i \partial_{j^*} \Psi \Rightarrow \mathbb{C}\mathbf{P}^{N-1} \text{ metric!}$$

The new coordinate system $\{\rho, \varphi^i\}$ is **well-defined** at $\rho = 0$.

$b = 0$ case:

Kähler potential is very simple:

$$\mathcal{K}|_{b=0} = \lambda^{\frac{1}{N}} e^X = \lambda^{\frac{1}{N}} |\sigma|^2 (1 + |\varphi^i|^2)$$

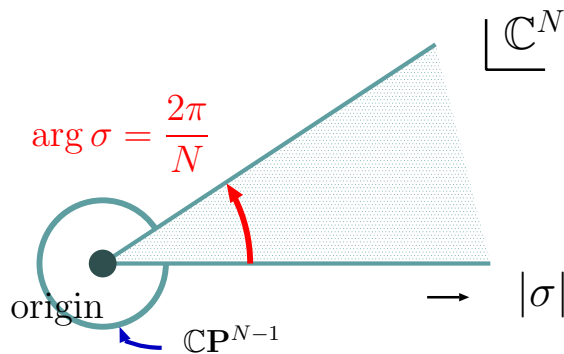
If we re-write the coordinates as $\phi^1 = \sigma$, $\phi^i = \sigma \varphi^{i-1}$: $\{\vec{\phi}^T = \sigma(1, \varphi^i)\}$

$$\mathcal{K} = \lambda^{\frac{1}{N}} \vec{\phi}^\dagger \vec{\phi} \iff \text{flat metric?}$$

But we need to transform $\rho = \sigma^N / N$
in order to avoid the coordinate singularity

↓

Orbifold $\mathbb{C}^N / \mathbb{Z}_N$



Result:

- This manifold has a parameter b .
- If $b \neq 0$, this is $\mathbb{C}^1 \times \mathbb{C}\mathbb{P}^{N-1}$ and has no singularity at the origin.
- Orbifold singularity $\mathbb{C}^N / \mathbb{Z}_N$ appears in $b = 0$ limit.

↓

canonical line bundle over $\mathbb{C}\mathbb{P}^{N-1}$

CONIFOLD

the similar set-up to the line bundle over $\mathbb{C}\mathbf{P}^{N-1}$:

$$\vec{\phi}^T = \sigma \left(1, \varphi^i, -\frac{1}{2}(\varphi^i)^2 \right), \quad \sigma \in \mathbb{C}^1$$

Assumption:

$$\mathcal{K} = \mathcal{K}(X), \quad X = \log \vec{\phi}^\dagger \vec{\phi} = \log |\sigma|^2 + \Psi$$

Ψ : Kähler potential of Q^{N-2} ($c = 1$)

the solution of Ricci-flatness condition: $\frac{d\mathcal{K}}{dX} = (\lambda e^{(N-2)X} + b)^{\frac{1}{N-1}}$

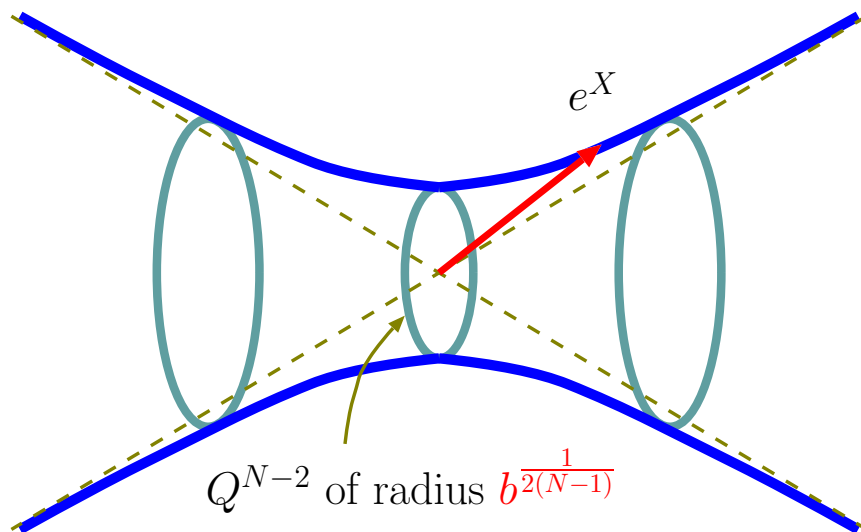
metric with $b \neq 0$: ($ds^2 = g_{\sigma\sigma^*} d\sigma d\sigma^* + \dots$)

$$g_{\sigma\sigma^*} = \frac{N-2}{N-1} \lambda (\lambda e^{(N-2)X} + b)^{-\frac{N-2}{N-1}} e^{(N-2)\Psi} |\sigma|^{2N-6}$$

coordinate transformation: $\rho \sim \sigma^{N-2}$

$\rho = 0$ ($d\rho = 0$) sub-manifold:

$$g_{ij^*} \Big|_{\rho=0} = b^{\frac{1}{N-1}} \partial_i \partial_{j^*} \Psi \iff Q^{N-2} \text{ metric!}$$



$N = 3$ solution : [Eguchi-Hanson gravitational instanton](#)

$$\mathcal{K}(X) = 2\sqrt{\lambda e^X + b} + \sqrt{b} \log \left(\frac{\sqrt{\lambda e^X + b} - \sqrt{b}}{\sqrt{\lambda e^X + b} + \sqrt{b}} \right)$$

re-definition: $\varrho^4 \equiv 4(\lambda e^X + b)$, $a^4 \equiv 4b$

$$\rightarrow \mathcal{K} = \varrho^2 + \frac{a^2}{2} \log \left(\frac{\varrho^2 - a^2}{\varrho^2 + a^2} \right)$$

Kähler potential of [Eguchi-Hanson](#) (Gibbons-Pope, 1979)

the node is repaired by $Q^1 = SO(3)/U(1) \simeq \mathbf{S}^2$

$N = 4$ solution : [\$Q^2\$ Resolved Conifold](#)

the node is repaired by

$$Q^2 = SO(4)/[SO(2) \times U(1)] \simeq \mathbf{S}^2 \times \mathbf{S}^2 \text{ (same radius)}$$

neither [deformation](#) (\mathbf{S}^3) nor [small resolution](#) (\mathbf{S}^2)

exceptional groups:

line bundle over $\mathbb{C}\mathbf{P}^{26} + [\Gamma_{ijk}\phi^j\phi^k = 0]$

\implies line bundle over $E_6/[SO(10) \times U(1)]$

line bundle over $\mathbb{C}\mathbf{P}^{55} + [d_{\alpha\beta\gamma\delta}\phi^\beta\phi^\gamma\phi^\delta = 0]$

\implies line bundle over $E_7/[E_6 \times U(1)]$

the solution of Ricci-flatness condition:

$$\frac{d\mathcal{K}}{dX} = (\lambda e^{cX} + b)^{\frac{1}{D}}, \quad \rho = \sigma^n/n$$

line bundle	\mathcal{C}	D	n
$\mathbb{C} \times \frac{E_6}{SO(10) \times U(1)}$	12	17	12
$\mathbb{C} \times \frac{E_7}{E_6 \times U(1)}$	18	28	18

GAUGE THEORETICAL CONSTRUCTION

Tool: **Supersymmetric Nonlinear Sigma Model as Gauge Theory**

(introduce the gauge field as an auxiliary field)

Grassmannian $G_{N,M}$:

$\Phi : N \times M$ matrix-valued chiral superfield

$U(N)_L \times U(M)_R$ transformation:

$$\Phi \rightarrow \Phi' = g_L \Phi g_R, \quad (g_L, g_R) \in (U(N)_L, U(M)_R).$$

$U(M)_R$ is gauged (introduce a $U(M)$ gauge field):

$$\Phi \rightarrow \Phi' = \Phi e^{-i\Lambda}, \quad e^V \rightarrow e^{V'} = e^{i\Lambda} e^V e^{-i\Lambda^\dagger}.$$

$U(N)_L \times U(M)_R$ invariant Lagrangian (linear sigma model):

$$\mathcal{L} = \int d^4\theta \left\{ \text{tr}(\Phi^\dagger \Phi e^V) - c \text{tr} V \right\}.$$

V : vector superfield ($U(M)^\mathbb{C}$ gauge group)

integrate out $V \implies$ target space becomes Grassmannian $G_{N,M}$

gauge-fixing: $\Phi = \begin{pmatrix} \mathbf{1}_M \\ \varphi_{Aa} \end{pmatrix},$

$\varphi_{Aa} : (N - M) \times M$ matrix-valued chiral superfield

$$\mathcal{K} = c \log \det (\mathbf{1}_M + \varphi^\dagger \varphi)$$

$G_{2N,N} + \text{F-term constraints}$:

$G_{2N,N}$ Lagrangian + Superpotential term: $W = \text{tr}(\Phi_0 \Phi^T J' \Phi)$

Φ_0 : $N \times N$ matrix-valued auxiliary superfield

$$J' = \begin{pmatrix} 0 & \mathbf{1}_N \\ \epsilon \mathbf{1}_N & 0 \end{pmatrix}, \quad \epsilon = \pm 1$$

$$\mathcal{L} = \int d^4\theta \mathcal{K}(\Phi, \Phi^\dagger, V) + \left(\int d^2\theta W(\Phi_0, \Phi) + \text{c.c.} \right)$$

integrate out $V, \Phi_0 \rightarrow \text{NLSM}$

$$\mathcal{K} = c \log \det \{ \mathbf{1}_N + \varphi^\dagger \varphi \}$$

$$\frac{Sp(N)}{U(N)} \quad (\epsilon = -1, \varphi^T - \varphi = 0), \quad \frac{SO(2N)}{U(N)} \quad (\epsilon = +1, \varphi^T + \varphi = 0)$$

line bundle over $G_{N,M}$:

$U(M)$ gauge symmetric Lagrangian:

$$\mathcal{K}_0(\Phi, \Phi^\dagger, V) = f(\text{tr}(\Phi^\dagger \Phi e^V)) - c \text{tr} V$$

f : arbitrary function of $\text{tr}(\Phi^\dagger \Phi e^V)$, $V = V^a T_a$, $T_a \in U(M)$

c : FI constant \Rightarrow vector superfield C

equation of motion for V and C :

$$\partial \mathcal{L} / \partial V = f' \cdot \Phi^\dagger \Phi e^V - C \cdot \mathbf{1}_M = 0$$

$$\partial \mathcal{L} / \partial C = -\text{tr} V = 0$$

under $\text{tr} V = 0$, perform the trace and determinant of the 1st eq.:

$$f' \cdot \text{tr}(\Phi^\dagger \Phi e^V) = M \cdot C, \quad (f')^M \cdot \det \Phi^\dagger \Phi = C^M$$

$$\text{delete } C: \text{tr}(\Phi^\dagger \Phi e^V) = M [\det \Phi^\dagger \Phi]^{\frac{1}{M}}$$

\Downarrow

under eq. of motion $\partial \mathcal{L} / \partial C = 0$ ($U(M) \rightarrow SU(M)$ gauge group):

$$\mathcal{K}_0 = f(M[\det \Phi^\dagger \Phi]^{\frac{1}{M}}) \equiv \mathcal{K}(\underbrace{\log \det \Phi^\dagger \Phi}_{\parallel X})$$

introduce $C =$ ungauged the $U(1)$ symmetry

the simplest construction of the line bundle

$$\Phi = \sigma \begin{pmatrix} \mathbf{1}_M \\ \varphi_{Aa} \end{pmatrix}, \sigma \in \mathbb{C}^1 \implies X = M^2 \log |\sigma|^2 + \Psi$$

Later discussion is same as the canonical line bundle over $\mathbb{C}\mathbf{P}^{N-1}$.

Solutions of Ricci-flatness condition:

$$\frac{d\mathcal{K}}{dX} = (\lambda e^{cX} + b)^{\frac{1}{D}}, \quad \rho = \sigma^n/n$$

line bundle	\mathcal{C}	D	n
$\mathbb{C} \times G_{N,M}$	N	$1 + M(N - M)$	MN
$\mathbb{C} \times \frac{Sp(N)}{U(N)}$	$N + 1$	$1 + \frac{1}{2}N(N + 1)$	$N(N + 1)$
$\mathbb{C} \times \frac{SO(2N)}{U(N)}$	$N - 1$	$1 + \frac{1}{2}N(N - 1)$	$N(N - 1)$

Isomorphism and duality:

Eguchi-Hanson space (complex two-dimensions)

$$\mathbb{C}^1 \times \mathbb{C}\mathbf{P}^1 \simeq \mathbb{C}^1 \times \frac{SO(4)}{U(2)} \simeq \mathbb{C}^1 \times \frac{Sp(1)}{U(1)} \simeq \mathbb{C}^1 \times Q^1$$

Line bundle over $\mathbb{C}\mathbf{P}^3$ (complex four-dimensions)

$$\mathbb{C}^1 \times \mathbb{C}\mathbf{P}^3 \simeq \mathbb{C}^1 \times \frac{SO(6)}{U(3)}$$

Another complex four-dimensional manifolds

$$\mathbb{C}^1 \times \frac{Sp(2)}{U(2)} \simeq \mathbb{C}^1 \times Q^3$$

Line bundle over the Klein quadric (complex five-dimensions)

$$\mathbb{C}^1 \times G_{4,2} \simeq \mathbb{C}^1 \times Q^4$$

Duality between Grassmannians

$$\mathbb{C}^1 \times G_{N,M} \simeq \mathbb{C}^1 \times G_{N,N-M}$$

Ricci-flatness solutions:

$$\frac{d\mathcal{K}}{dX} = (\lambda e^{cX} + b)^{\frac{1}{D}}, \quad \rho = \sigma^n/n.$$

Solutions:

type	$\mathbb{C} \times G/H$	\mathcal{C}	D	n
AIII ₁	$\mathbb{C} \times \mathbb{C}\mathbf{P}^{N-1}$	N	$1 + (N - 1)$	N
AIII ₂	$\mathbb{C} \times G_{N,M}$	N	$1 + M(N - M)$	MN
BDI	$\mathbb{C} \times Q^{N-2}$	$N - 2$	$1 + (N - 2)$	$N - 2$
CI	$\mathbb{C} \times Sp(N)/U(N)$	$N + 1$	$1 + \frac{1}{2}N(N + 1)$	$N(N + 1)$
DIII	$\mathbb{C} \times SO(2N)/U(N)$	$N - 1$	$1 + \frac{1}{2}N(N - 1)$	$N(N - 1)$
EIII	$\mathbb{C} \times E_6/[SO(10) \times U(1)]$	12	$1 + 16$	12
EVII	$\mathbb{C} \times E_7/[E_6 \times U(1)]$	18	$1 + 27$	18

$$D = \dim_{\mathbb{C}}(\mathbb{C} \times G/H), \quad \mathcal{C} = \frac{1}{2}C_2(G)$$

$$Q^1 \simeq \mathbb{C}\mathbf{P}^1 \simeq SO(4)/U(2) \simeq Sp(1)/U(1)$$

$$\mathbb{C}\mathbf{P}^3 \simeq SO(6)/U(3)$$

$$Q^2 \simeq \mathbb{C}\mathbf{P}^1 \times \mathbb{C}\mathbf{P}^1$$

$$Q^4 \simeq G_{4,2}$$

$$Q^3 \simeq Sp(2)/U(2)$$

$$G_{N,M} \simeq G_{N,N-M}$$

GENERALIZATION

If the base manifolds are **Einstein-Kähler**,

We can construct the other line bundles:

$$\mathbb{C}^1 \times (\text{non-symmetric spaces}) : \quad \mathbb{C}^1 \times \frac{SU(\ell + m + n)}{S[U(\ell) \times U(m) \times U(n)]}$$

$$\mathbb{C}^1 \times (\text{direct product of HSS}) : \quad \mathbb{C}^1 \times \mathbb{C}\mathbf{P}^{N-1} \times \mathbb{C}\mathbf{P}^{M-1}$$

omit the explanation (mathematically very complicated!),

please see [hep-th/0202064](#).

SUMMARY

Supersymmetric Nonlinear Sigma Model as Gauge Theory

↓

$U(1)$ ungauged \rightarrow non-compact Kähler manifold

G/H is Einstein-Kähler \rightarrow “Ricci-flatness condition = ODE”

coordinate transformation $\rho \sim \sigma^n \rightarrow$ “delete coordinate singularity”

integration constant $b \neq 0 \rightarrow$ “delete **conical singularity**”

(the node is repaired by HSS)

We obtained the canonical line bundle over HSS

↓

“**Non-compact Calabi-Yau Manifolds**”

Eguchi-Hanson space in real four-dimensions

New deformation of the conifold: $\mathbb{C}^1 \times Q^2 \simeq \mathbb{C}^1 \times \mathbf{S}^2 \times \mathbf{S}^2$

New metric: $\mathbb{C}^1 \times G_{N,M}$, etc.

General Solutions of Kähler Potential:

$$\frac{d\mathcal{K}}{dX} = (\lambda e^{cX} + b)^{\frac{1}{D}}$$

$$\mathcal{K}(X) = \frac{D}{c} \left[(\lambda e^{cX} + b)^{\frac{1}{D}} + b^{\frac{1}{D}} \cdot I(b^{-\frac{1}{D}} (\lambda e^{cX} + b)^{\frac{1}{D}}; D) \right]$$

$$\begin{aligned} I(y; D) &\equiv \int^y \frac{dt}{t^D - 1} = \frac{1}{D} \left[\log(y - 1) - \frac{1 + (-1)^D}{2} \log(y + 1) \right] \\ &\quad + \frac{1}{D} \sum_{r=1}^{\lfloor \frac{D-1}{2} \rfloor} \cos \frac{2r\pi}{D} \log \left(y^2 - 2y \cos \frac{2r\pi}{D} + 1 \right) \\ &\quad + \frac{2}{D} \sum_{r=1}^{\lfloor \frac{D-1}{2} \rfloor} \sin \frac{2r\pi}{D} \arctan \left[\frac{\cos(2r\pi/D) - y}{\sin(2r\pi/D)} \right] \end{aligned}$$