

*Spectrum of Eleven-dimensional Supergravity
on a PP-wave Background*

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Introduction: Eleven-dimensional Supergravity

Lagrangian of eleven-dimensional supergravity (without torsion):

$$\begin{aligned}\mathcal{L} = & e\mathcal{R} - \frac{1}{2}e\bar{\Psi}_M\hat{\Gamma}^{MNP}D_N\Psi_P - \frac{1}{48}eF_{MNPQ}F^{MNPQ} \\ & + \frac{1}{192}e\bar{\Psi}_M\tilde{\Gamma}^{MNPQRS}\Psi_N F_{PQRS} + \frac{1}{(144)^2}\epsilon^{MNPQRSUVWXY}F_{MNPQ}F_{RSUV}C_{WXYZ}\end{aligned}$$

Classical Field Equations

$$\begin{aligned}0 &= \frac{1}{2}g_{MN}\mathcal{R} - \mathcal{R}_{MN} - \frac{1}{96}g_{MN}F_{PQRS}F^{PQRS} + \frac{1}{12}F_{MPQR}F_N{}^{PQR} \\ 0 &= \hat{\Gamma}^{MNP}D_N\Psi_P - \frac{1}{96}\tilde{\Gamma}^{MNPQRS}\Psi_N F_{PQRS} \\ 0 &= \nabla^Q\{eF_{QMNP}\} + \frac{18}{(144)^2}g_{MZ}g_{NK}g_{PL}\epsilon^{ZKLQRSUVWXY}F_{QRSU}F_{VWXY}\end{aligned}$$

Spectrum on $AdS_4 \times S^7$

$$\Phi_{\mu\nu\dots mn\dots}(x, y) = \sum_I \phi_{\mu\nu\dots}^I(x) \cdot Y_{mn\dots}^I(y)$$

$d = 11$	$d = 4 (AdS_4)$	spin	$S^7 (SO(8))$	number
$g_{MN}(x, y)$	$h_{\mu\nu}(x)$	2	1	1
	$V_\mu^{[ij]}(x)$	1	$K_m^{[ij]}(y)$	28
	$S^{[ijkl]}(x)$	0	$K_{(m}^{[ij} K_{n)}^{kl]} - \frac{1}{9} g_{mn} K_p^{[ij} K^{kl]p}$	35
$C_{MNP}(x, y)$	$P^{[ijkl]}(x)$	0	$\nabla_{[m} K_n^{[ij} \nabla_p K_q^{kl]}$	35
$\Psi_M(x, y)$	$\psi_\mu^I(x)$	3/2	$\eta^I(y)$	8
	$\chi^{[IJK]}(x)$	1/2	$\eta_m^{[IJK]} + \frac{1}{9} \hat{\Gamma}_m \eta^{[IJK]}$	56

$K_m^{[ij]}(y)$: Killing vector ($i = 1, \dots, 8$)

$\eta^I(y)$: Killing spinor ($I = 1, \dots, 8$)

Maximally Supersymmetric Spaces



Kowalski-Glikman Solution (PP-wave Background)

$$\begin{aligned}
 ds^2 &= -2dx^+dx^- + G_{++}(dx^+)^2 + \sum_{I=1}^9(dx^I)^2 & x^\pm &= \frac{1}{\sqrt{2}}(x^0 \pm x^{10}) \\
 G_{++} &= -\left[\left(\frac{\mu}{3}\right)^2 \sum_{\tilde{I}=1}^3(x^{\tilde{I}})^2 + \left(\frac{\mu}{6}\right)^2 \sum_{I'=4}^9(x^{I'})^2 \right] & \mu &= F_{+123} \neq 0
 \end{aligned}$$

Hamiltonian

We will encounter Klein-Gordon type equations of motion and have to evaluate its energy spectrum. Klein-Gordon type equation of motion for a field $\phi(x)$:

$$(\square - \alpha \mu i \partial_-) \phi(x^+, x^-, x^I) = 0$$

α : arbitrary constant x^+ : evolution parameter

We express the Hamiltonian $H = i\partial_+$:

$$H = \frac{1}{3}\mu \sum_{\tilde{I}} \bar{a}^{\tilde{I}} a^{\tilde{I}} + \frac{1}{6}\mu \sum_{I'} \bar{a}^{I'} a^{I'} + \frac{1}{2}\mu (2 - \alpha)$$

Last term = zero-point energy E_0 of the system (eigenvalue of H):

$$E_0 = \frac{1}{2}\mu \mathcal{E}_0(\phi) \quad \mathcal{E}_0(\phi) = 2 - \alpha$$

Bosonic/Fermionic Spectrum

Fluctuation fields are expanded around classical fields as follows:

$$g_{MN} \rightarrow g_{MN} + h_{MN}$$

g_{MN} : pp-wave background

$$\Psi_M \rightarrow 0 + \psi_M$$

$$C_{MNP} \rightarrow C_{MNP} + \mathcal{C}_{MNP}$$

$$4\partial_{[+}C_{123]} = F_{+123} = \mu$$

Constraints:

$$\left. \begin{aligned} h_{-M} &= h^{+N} = 0 \\ \psi_- &= 0 \\ \mathcal{C}_{-NP} &= 0 \end{aligned} \right\}$$

light-cone gauge-fixing

Bosonic Fields

$$\begin{aligned}
 0 &= \square \mathcal{C}_{\tilde{I}J'K'} & \tilde{I} &: SO(3) & I' &: SO(6) \\
 0 &= (\square - \mu i \partial_-) H_{\tilde{I}J'} & H_{\tilde{I}J'} &= h_{\tilde{I}J'} + i \mathcal{C}_{\tilde{I}J'} & \mathcal{C}_{\tilde{I}J'} &= \frac{1}{2} \varepsilon_{\tilde{I}\tilde{K}\tilde{L}} \mathcal{C}_{\tilde{K}\tilde{L}J'} \\
 0 &= \square h_{\tilde{I}\tilde{J}}^\perp & h_{\tilde{I}\tilde{J}}^\perp &= h_{\tilde{I}\tilde{J}} - \frac{1}{3} \delta_{\tilde{I}\tilde{J}} h_{\tilde{K}\tilde{K}} \\
 0 &= (\square - 2\mu i \partial_-) h & h &= h_{\tilde{I}\tilde{I}} + i \mathcal{C} & \mathcal{C} &= 2 \mathcal{C}_{123} \\
 0 &= (\square - \mu i \partial_-) \mathcal{C}_{I'J'K'}^\ominus & \mathcal{C}_{I'J'K'}^\ominus &= -\frac{i}{3!} \varepsilon^{I'J'K'L'M'N'} \mathcal{C}_{L'M'N'}^\ominus
 \end{aligned}$$

Now we have derived the spectrum:

energy	bosonic fields			degrees of freedom
4	\bar{h}			1
3	$\bar{H}_{\tilde{I}J'}$	$\mathcal{C}_{I'J'K'}^\oplus$		18 + 10
2	$\mathcal{C}_{\tilde{I}J'K'}$	$h_{\tilde{I}\tilde{J}}^\perp$	$h_{I'J'}$	45 + 5 + 20
1	$H_{\tilde{I}J'}$	$\mathcal{C}_{I'J'K'}^\ominus$		18 + 10
0	h			1

Results

energy	bosons			fermions		D.O.F.
4	\bar{h}					1
7/2				$\psi_L^{\oplus\parallel}$		8
3	$\bar{H}_{\tilde{I}J'}$	$\mathcal{C}_{I'J'K'}^{\oplus}$				28
5/2				$\psi_{\tilde{I}R}^{\oplus\perp}$	$\psi_{I'L}^{\oplus\perp}$	56
2	$\mathcal{C}_{\tilde{I}J'K'}$	$h_{\tilde{I}J}^{\perp}$	$h_{I'J'}^{\perp}$			70
3/2				$\psi_{\tilde{I}L}^{\oplus\perp}$	$\psi_{I'R}^{\oplus\perp}$	56
1	$H_{\tilde{I}J'}$	$\mathcal{C}_{I'J'K'}^{\ominus}$				28
1/2				$\psi_R^{\oplus\parallel}$		8
0	h					1

Energy spectrum and degrees of freedom of physical modes in 11-dimensional supergravity on the pp-wave background.

Discussions: Developments around 11-dimensional Theory

11-dimensional Supergravity was born (1978)

Freund-Rubin ansatz, spontaneous compactification ($AdS_{4(7)} \times S^7(4)$) (\sim 1982)

but, *Supergravity is Dead!* (by A.Salam, 1987)

Supermembrane in flat 11-dimensions (1988)

unstable, continuous spectrum (1989)

⋮

M-theory, BFSS conjecture (1996)

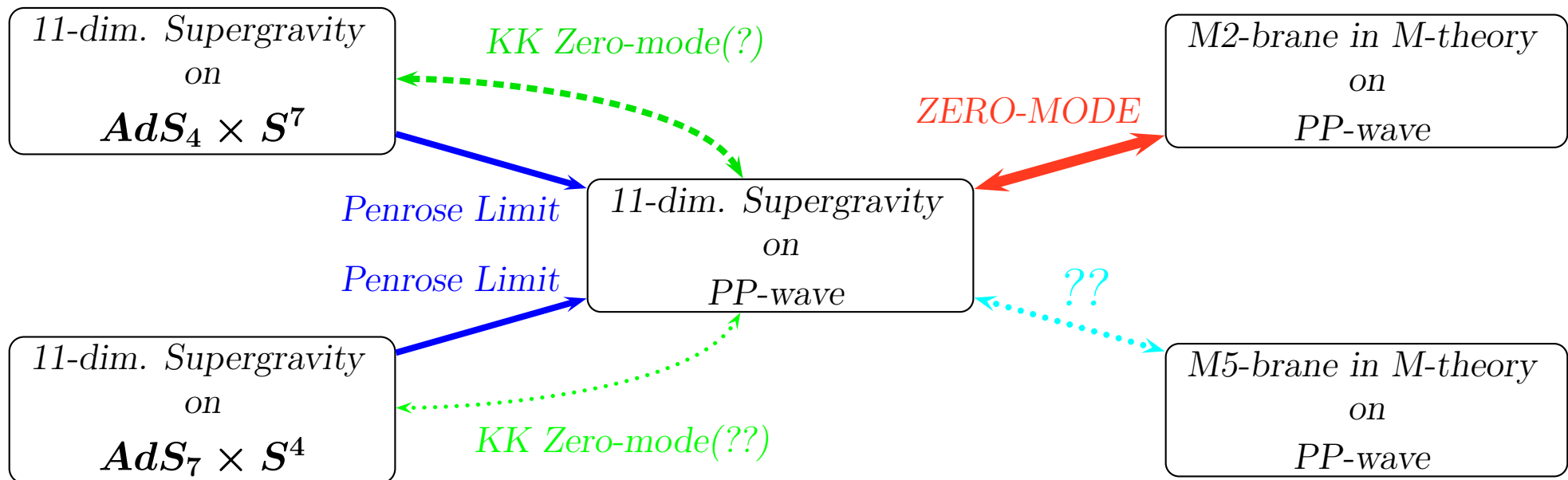
matrix model (supermembrane) revisited

supermembrane (matrix model) on PP-wave background (2002)

Discretized Spectrum, Mass Gap, etc. (2002)

⇓

Supergravity revisited!!



Work in Progress

- Comparison with KK zero-modes of $AdS_p \times S^q$
- Propagators and energy-momentum tensors of h_{MN} , \mathcal{C}_{MNP} and ψ_M
- Dimensional reduction to type IIA supergravity (only 24 supercharges)