Gauged Linear Sigma Models

for Noncompact Calabi-Yau Varieties

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Two-dimensional field theory is a powerful framework

We have studied 2-dim. SUSY nonlinear sigma models...

– 1999: Supersymmetric nonlinear sigma models on hermitian symmetric spaces
introducing an auxiliary gauge field
by Higashijima and Nitta

2000: $1/N$ expansion of SUSY NLSM on $Q^{N-2} = \dfrac{SO(N)}{SO(N-2) \times U(1)}$
two non-trivial vacua, asymptotically free
by Higashijima, Nitta, Tsuzuki and TK

2001 – 2002: Ricci-flat metrics on noncompact Kähler manifolds ($\equiv$ noncompact Calabi-Yau's)
by Higashijima, Nitta and TK

Metrics on Noncompact Calabi-Yau

(K.Higashijima, M.Nitta and TK, 2001, 2002)

line bundles	total dim. D	dual Coxeter C	"orbifolding" ℓ
$\mathbb{C} \ltimes \left(\mathbb{C}\mathrm{P}^{N-1} = \frac{SU(N)}{SU(N-1)\times U(1)}\right)$	$1+(N-1)$	N	N
$\mathbb{C} \ltimes \left(Q^{N-2} = \frac{SO(N)}{SO(N-2)\times U(1)}\right)$	$1+(N-2)$	$N-2$	$N-2$
$\mathbb{C} \ltimes E_6/[SO(10)\times U(1)]$	$1+16$	12	12
$\mathbb{C} \ltimes E_7/[E_6\times U(1)]$	$1+27$	18	18
$\mathbb{C} \ltimes \left(\mathrm{G}_{N,M} = \frac{U(N)}{U(N-M)\times U(M)}\right)$	$1+M(N-M)$	N	MN
$\mathbb{C} \ltimes SO(2N)/U(N)$	$1+\frac{1}{2}N(N-1)$	$N-1$	$N(N-1)$
$\mathbb{C} \ltimes Sp(N)/U(N)$	$1+\frac{1}{2}N(N+1)$	$N+1$	$N(N+1)$

$$\mathcal{K}'_{\mathrm{noncompact}}(\rho,\varphi) \;=\; \left(\mathrm{e}^{CX}+b\right)^{1/D}$$

$$X \;=\; \log|\rho^{1/\ell}|^2 + K_{\mathrm{compact}}(\varphi)\,, \qquad K_{\mathbb{C}\mathrm{P}^{N-1}}(\varphi) = r\log\left(1+\sum_{i=1}^{N-1}|\varphi_i|^2\right)$$

CFT descriptions (Virasoro- and current-algebras)?

global aspects of noncompact geometries?

and

mirror geometries?

Gauged linear sigma model E. Witten (1993), K. Hori and C. Vafa (2000)

$\mathcal{N} = (2,2)$ SUSY gauge theory with matters $(\mathrm{FI}: \ t \equiv r - i\theta)$

$$\mathscr{L} = \int \mathrm{d}^4\theta \Big\{ -\frac{1}{e^2}\overline{\Sigma}\Sigma + \sum_a \overline{\Phi}_a \, \mathrm{e}^{2Q_a V} \Phi_a \Big\} + \Big(\frac{1}{\sqrt{2}} \int \mathrm{d}^2\widetilde{\theta} \, (-\Sigma \, t) + c.c. \Big) + \Big(\int \mathrm{d}^2\theta \, W_{\mathrm{GLSM}}(\Phi_a) + c.c. \Big)$$

- Φ_a : charged chiral superfield, $\overline{D}_{\pm}\Phi_a = 0$
  Σ : twisted chiral superfield, $\overline{D}_+\Sigma = D_-\Sigma = 0$, $\Sigma = \frac{1}{\sqrt{2}}\overline{D}_+ D_- V$
- There exist at least two phases:

  $\mathrm{FI} \gg 0$: differential-geometric phase $\longrightarrow$ SUSY NLSM

  $\mathrm{FI} \ll 0$: algebro-gemetric phase $\longrightarrow$ LG, orbifold, SCFT

- Calabi-Yau/Landau-Ginzburg correspondence

  harmonic forms $\longleftrightarrow$ NS-NS chiral primary states

- "Mirror geometry" appears in the T-dual theory in terms of twisted chiral superfields Y_a

$$Y_a + \overline{Y}_a \equiv 2 \, \overline{\Phi}_a \, \mathrm{e}^{2Q_a V} \Phi_a$$

Effective theories

The potential energy density is given by

$$\mathcal{U}(\phi,\sigma) \;=\; \frac{e^2}{2}\mathcal{D}^2 + \sum_a |F_a|^2 + \mathcal{U}_\sigma(\phi,\sigma)$$

$$\mathcal{D} \;=\; \frac{1}{e^2}D \;=\; r - \sum_a Q_a|\phi_a|^2\,, \quad \overline{F}_a \;=\; -\frac{\partial}{\partial\phi_a}W_{\text{GLSM}}(\phi)\,, \quad \mathcal{U}_\sigma(\phi,\sigma) \;=\; 2|\sigma|^2\sum_a Q_a^2|\phi_a|^2$$

The supersymmetric vacuum manifold $\mathcal{M}$ is defined by

$$\boxed{\mathcal{M} \;=\; \Big\{(\phi_a,\sigma)\in\mathbb{C}^n \,\Big|\, \mathcal{D} = F_a = \mathcal{U}_\sigma = 0\Big\}\Big/U(1)}$$

In the IR limit $e \to \infty$, there appears the supersymmetric NLSM on $\mathcal{M}$ whose coupling is

$$r \;=\; \frac{1}{g^2}$$

Renormalization of the FI parameter is $$\boxed{r_0 \;=\; r_R + s\cdot\log\left(\frac{\Lambda_{\text{UV}}}{\mu}\right), \quad s \;=\; \sum_a Q_a}$$

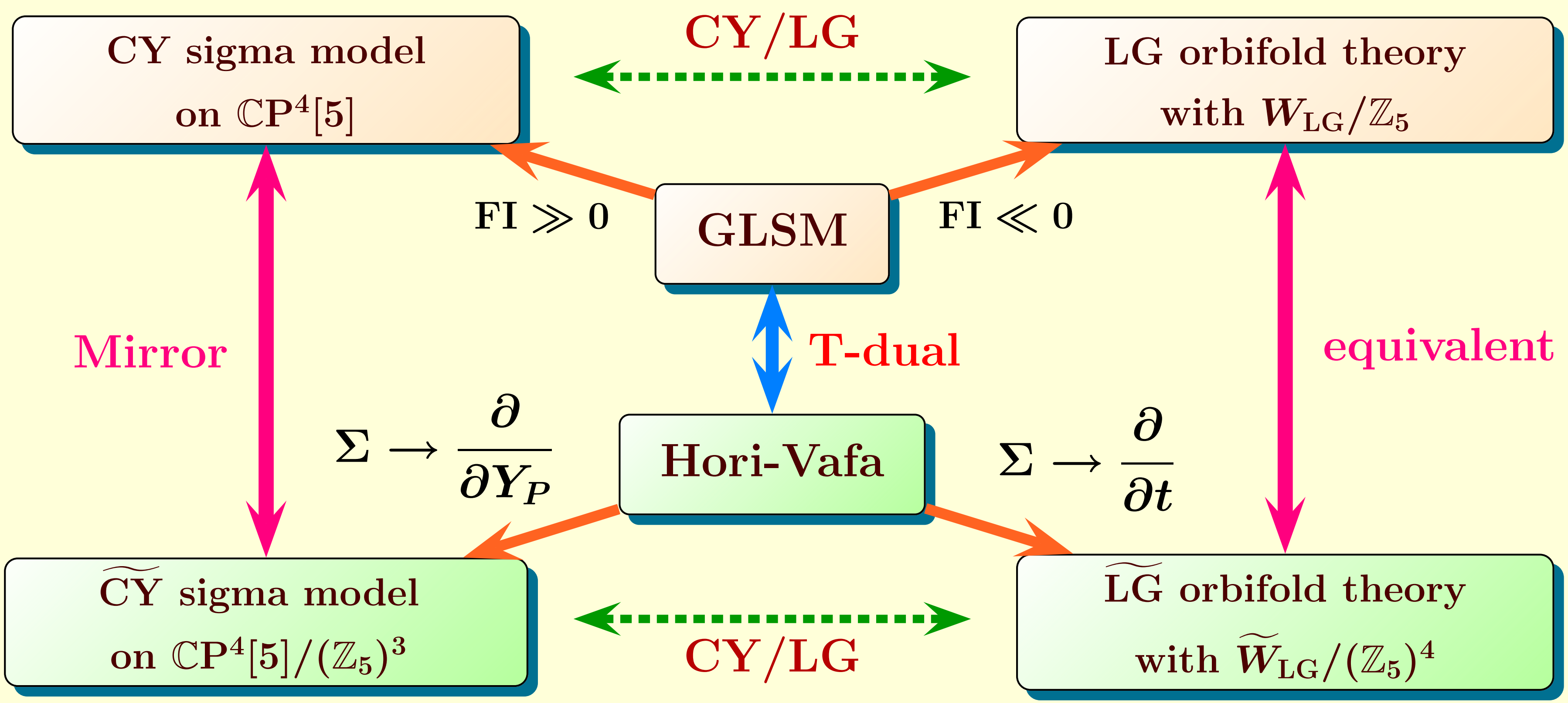
Thus we find that

$s > 0 \quad \longrightarrow \quad$ the theory is asymptotic free

$s = 0 \quad \longrightarrow \quad$ the theory is conformal

$s < 0 \quad \longrightarrow \quad$ the theory is infrared free

Example: quintic hypersurface and its mirror



Gauged Linear Sigma Model for $\mathcal{O}(-N+\ell)$ bundle on $\mathbb{C}\mathrm{P}^{N-1}[\ell]$

chiral superfield	S_1	$\cdots$	S_N	P_1	P_2
$U(1)$ charge	1	$\cdots$	1	$-\ell$	$-N+\ell$

$$W_{\mathrm{GLSM}} = P_1 \cdot G_\ell(S_i)$$

$G_\ell(S_i)$: homogeneous polynomial of degree ℓ

potential energy density:

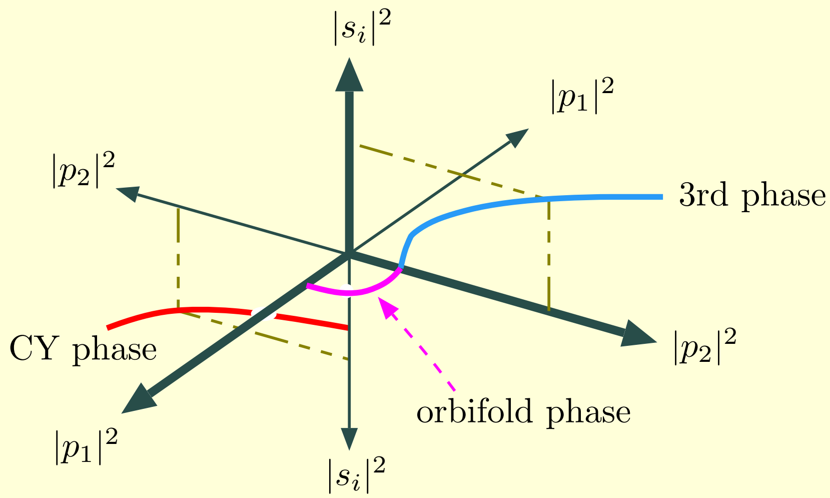
$$\mathcal{U} = \frac{e^2}{2}\mathcal{D}^2 + \big|G_\ell(s)\big|^2 + \sum_{i=1}^{N} \big|p_1 \partial_i G_\ell(s)\big|^2 + \mathcal{U}_\sigma$$

$$\mathcal{D} = r - \sum_{i=1}^{N} |s_i|^2 + \ell|p_1|^2 + (N-\ell)|p_2|^2$$

$$\mathcal{U}_\sigma = +2|\sigma|^2 \Big\{ \sum_{i=1}^{N} |s_i|^2 + \ell^2|p_1|^2 + (N-\ell)^2|p_2|^2 \Big\}$$

Let us analyze SUSY vacuum manifold $\mathcal{U} = 0$ and massless effective theories

Supersymmetric vacua



- **CY phase on $\mathcal{M}_{\rm CY}$**

  conformal sigma model on $\mathcal{M}_{\rm CY}$

- **orbifold phase on $\mathcal{M}^1_{r<0}$ (two "LG"s appear)**

  $\{$**CFT on** $\mathbb{C}^1 \otimes$ **LG with** $W_{\rm LG} = \langle p_1 \rangle G_\ell(S)\}/\mathbb{Z}_\ell$

  $\{$**"LG" with** $W_{\rm LG} = P_1 \cdot G_\ell(S)\}/\mathbb{Z}_{N-\ell}$

- **3rd phase on $\mathcal{M}^2_{\rm r<0}$ NEW!**

  conformal sigma model on $\mathcal{M}^2_{r<0}$

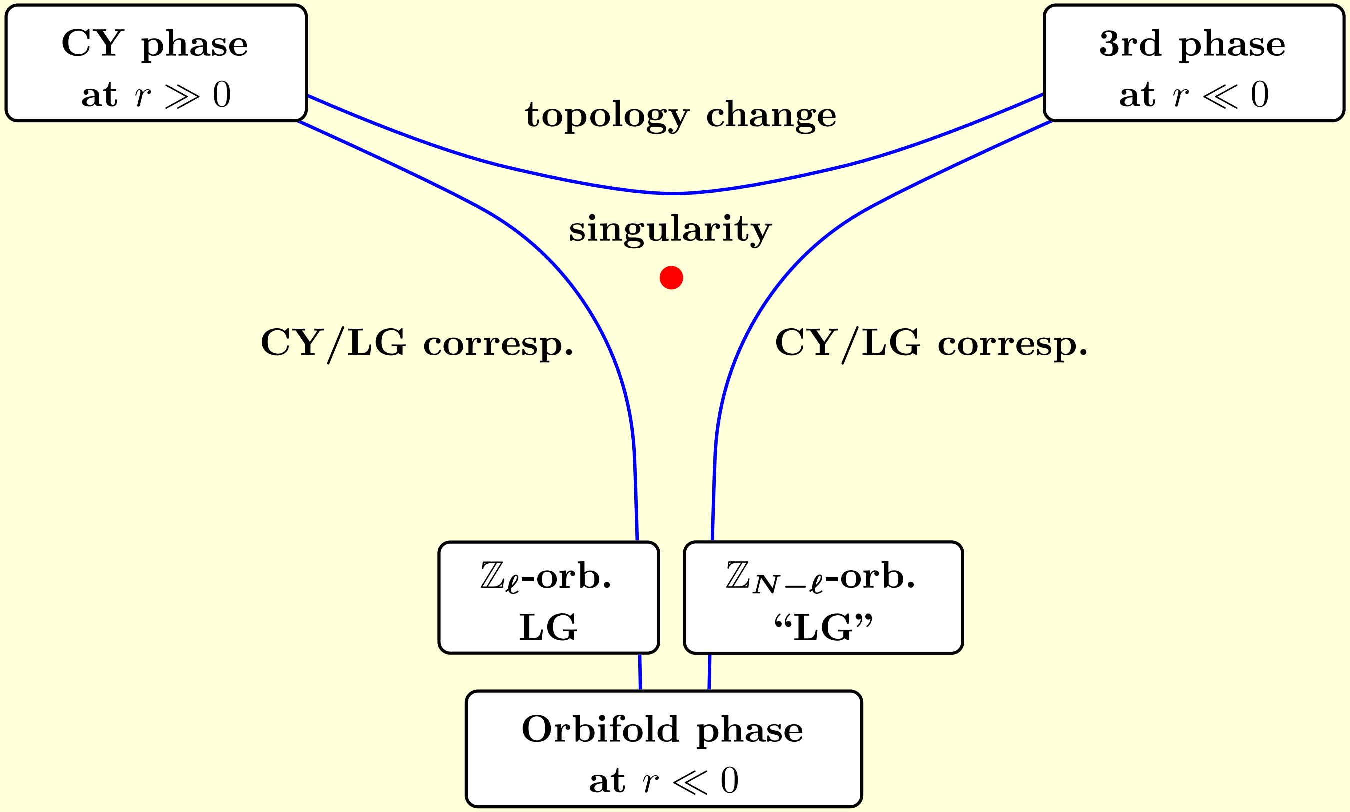
$$\mathcal{M}_{\rm CY} = \Big\{ (s_i; p_2) \in \mathbb{C}^{N+2} \Big| \mathcal{D} = G_\ell = 0,\ r > 0 \Big\} \Big/ U(1) \equiv \mathcal{O}(-N+\ell) \text{ bundle on } \mathbb{CP}^{N-1}[\ell]$$

$$\mathcal{M}^1_{r<0} = \Big\{ (p_1, p_2) \in \mathbb{C}^2 \Big| \mathcal{D} = 0, r < 0 \Big\} \Big/ U(1) \equiv \mathrm{W}\mathbb{CP}^1_{\ell, N-\ell}$$

$$\mathcal{M}^2_{r<0} = \Big\{ (s_i; p_2) \in \mathbb{C}^{N+2} \Big| \mathcal{D} = G_\ell = 0,\ r < 0 \Big\} \Big/ U(1)$$

CY/LG correspondence and topology change

The four theories are related to each other via CY/LG correspondence and topology change:



as a Conjecture

We also notice that we have obtained various massless effective theories by decomposing (not by integrating out) all massive modes. Thus they are just approximate descriptions.

T-dual description of the GLSM is also powerful to investigate low energy theories.

Analyzing them

we will re-investigate the massless effective theories in the original GLSM.

T-dual Theory

$$\mathscr{L} = \int d^4\theta \Big\{ -\frac{1}{e^2}\overline{\Sigma}\Sigma - \sum_a \Big(\frac{1}{2}(Y_a + \overline{Y}_a)\log(Y_a + \overline{Y}_a)\Big)\Big\} + \Big(\frac{1}{\sqrt{2}}\int d^2\widetilde{\theta}\,\widetilde{W} + c.c.\Big)$$

$$\widetilde{W} = \Sigma\Big(\sum_{i=1}^{N} Y_i - \ell Y_{P_1} - (N-\ell)Y_{P_2} - t\Big) + \sum_{i=1}^{N} e^{-Y_i} + e^{-Y_{P_1}} + e^{-Y_{P_2}}$$

Period integral : $\widehat{\Pi} \equiv \int d\Sigma \prod_{i=1}^{N} dY_i\, dY_{P_1}\, dY_{P_2}\,(\ell\Sigma)\,\exp\big(-\widetilde{W}\big)$

chiral superfield	S_1	S_2	$\cdots$	S_N	P_1	P_2
$U(1)$ charge	1	1	$\cdots$	1	$-\ell$	$-N+\ell$
twisted chiral	Y_1	Y_2	$\cdots$	Y_N	Y_{P_1}	Y_{P_2}

$$2\,\overline{\Phi}_a\, e^{2Q_aV}\,\Phi_a = Y_a + \overline{Y}_a$$

$U(1)$ phase rotation symmetry on Φ_a ➡ shift symmetry on Y_a: $Y_a \equiv Y_a + 2\pi i$

In the IR limit $e \to \infty$, the gauge field Σ is no longer dynamical and should be integrated out.

in order to obtain LG theory or geometry, we replace Σ to $\Sigma \to \frac{\partial}{\partial t}$ or $\Sigma \to \frac{\partial}{\partial Y_P}$

Twisted Landau-Ginzburg theory: There exist consistent solutions

— Solution one: $\mathbb{Z}_\ell$-type orbifold symmetry —

Solve Y_{P_1} by using the constraint derived from integrating out Σ:

$$Y_{P_1} = \frac{1}{\ell}\Big\{t - \sum_{i=1}^{N} Y_i + (N-\ell)Y_{P_2}\Big\}$$

Field re-definition preserving canonical measure in $\widehat{\Pi}$:

$$X_i \equiv e^{-\frac{1}{\ell}Y_i}, \quad X_{P_2} \equiv e^{\frac{N-\ell}{\ell}Y_{P_2}}, \quad X_i \to \omega_i X_i, \quad X_{P_2} \to \omega_{P_2} X_{P_2}, \quad (\mathbb{Z}_\ell)^N \text{ symmetry}$$

Thus we obtain the twisted LG superpotential:

$$\boxed{\Big\{\widetilde{W}_\ell = X_1^\ell + \cdots + X_N^\ell + X_{P_2}^{-\frac{\ell}{N-\ell}} + e^{t/\ell} X_1 \cdots X_N X_{P_2}\Big\}\Big/(\mathbb{Z}_\ell)^N}$$

The negative power term describes $\mathcal{N}=2$ Kazama-Suzuki model on $SL(2,\mathbb{R})_k/U(1)$:

$$\frac{\ell}{N-\ell} = k = \frac{2}{Q^2}$$

Thus we argue that

this effective theory is the LG minimal model coupled to the KS model with $(\mathbb{Z}_\ell)^N$ orbifold symmetry

Twisted mirror geometry: There also exist consistent solutions

— Solution one: $\mathbb{Z}_\ell$-type orbifold symmetry —

We replace $\ell\Sigma$ to $\dfrac{\partial}{\partial Y_{P_1}}$ and obtain

$$\widehat{\Pi} = \int \prod_{i=1}^{N} \mathrm{d}Y_i \, (\mathrm{e}^{-Y_{P_1}} \mathrm{d}Y_{P_1}) \, \mathrm{d}Y_{P_2} \, \delta\Big(\sum_i Y_i - \ell Y_{P_1} - (N-\ell) Y_{P_2} - t \Big) \exp\Big(- \sum_i \mathrm{e}^{-Y_i} - \mathrm{e}^{-Y_{P_1}} - \mathrm{e}^{-Y_{P_2}} \Big)$$

Re-defining the variables in order to obtain the canonical measure, we obtain

$$\widetilde{\mathcal{M}}_\ell = \Big\{ \{ \mathcal{F}(Z_i) = 0 \} / \mathbb{C}^* \, , \, \mathcal{G}(Z_b; u, v) = 0 \Big\} \Big/ (\mathbb{Z}_\ell)^{N-2}$$

$$\mathcal{F}(Z_i) = Z_1^\ell + \cdots + Z_\ell^\ell + \psi Z_1 \cdots Z_\ell \, , \quad \psi = \mathrm{e}^{t/\ell} Z_{\ell+1} \cdots Z_N$$

$$\mathcal{G}(Z_b; u, v) = Z_{\ell+1}^\ell + \cdots + Z_N^\ell + 1 - uv$$

$Z_a \mapsto \lambda \, \omega_a \, Z_a$ for $a = 1, \cdots, \ell$ (homogeneous coordinates of $\mathbb{C}\mathrm{P}^{\ell-1}[\ell]$)

$Z_b \mapsto \omega_b \, Z_b$ for $b = \ell+1, \cdots, N$ (homogeneous coordinates of $\mathbb{C}^{N-\ell}$)

$$\omega_a^\ell = \omega_b^\ell = \omega_1 \cdots \omega_N = 1 \, , \qquad \lambda : \mathbb{C}^*\text{-value}$$

Return to the original GLSM

Recall the following two arguments:

- $\mathcal{N}=2$ SCFT on $SL(2,\mathbb{R})_k/U(1)$ is equivalent to $\mathcal{N}=2$ Liouville theory via T-duality
- If a CFT $\mathcal{C}$ has an abelian discrete symmetry group Γ, the orbifold CFT $\mathcal{C}'=\mathcal{C}/\Gamma$ has a symmetry group Γ' which is isomorphic to Γ. Furthermore a new orbifold CFT $\mathcal{C}'/\Gamma'$ is identical to the original CFT $\mathcal{C}$.

Thus we insist that

"{CFT on $\mathbb{C}^1 \otimes$ LG with $W_{\text{LG}}=\langle p_1\rangle G_\ell(S)\}/\mathbb{Z}_\ell$" in the original GLSM

is described by

$\{\mathcal{N}=2$ Liouville theory coupled to the LG minimal model with $W_{\text{LG}}\}/\mathbb{Z}_\ell$

as an exact effective theory

Summary

- We found three non-trivial phases and four effective theories in the GLSM

  two CY sigma models

  two orbifolded LG theories coulped to 1-dim. SCFT

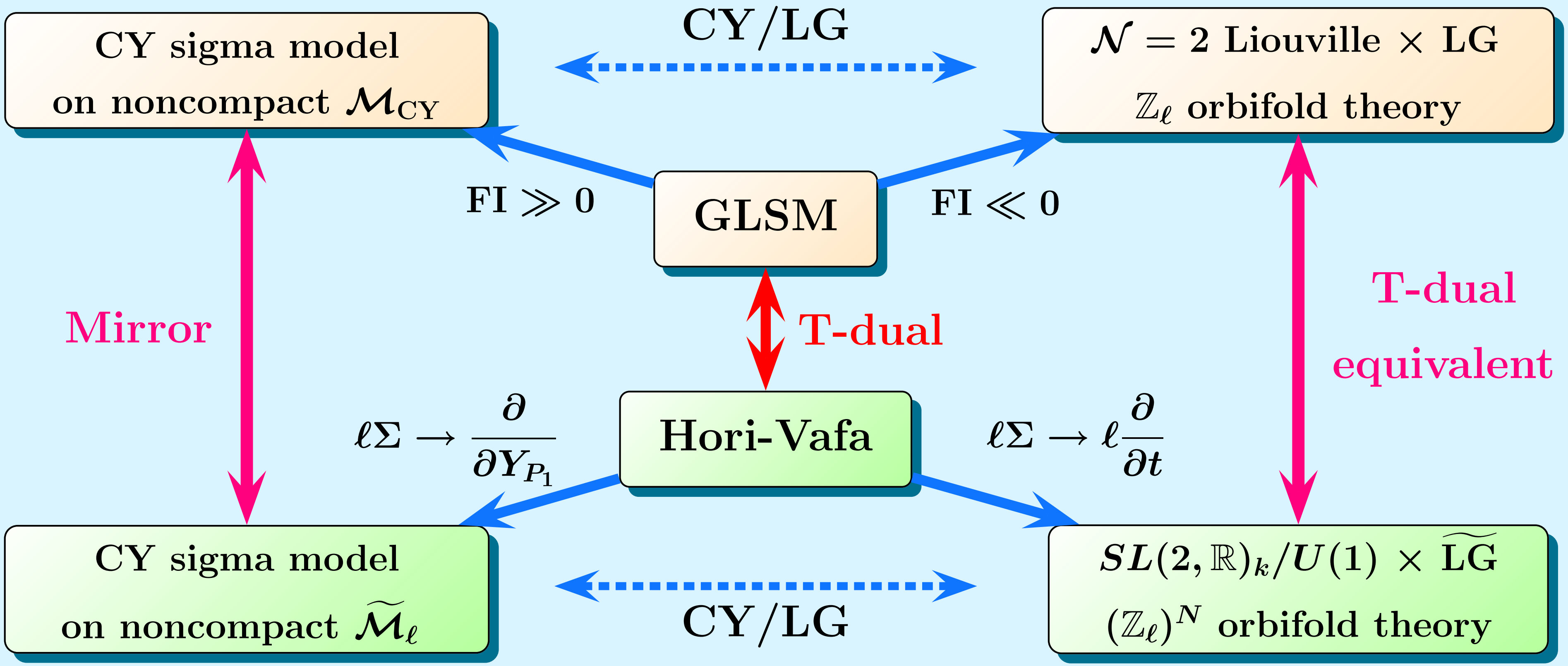
- We constructed four exact effective theories in the T-dual theory

  two NLSMs on mirror CY geometries

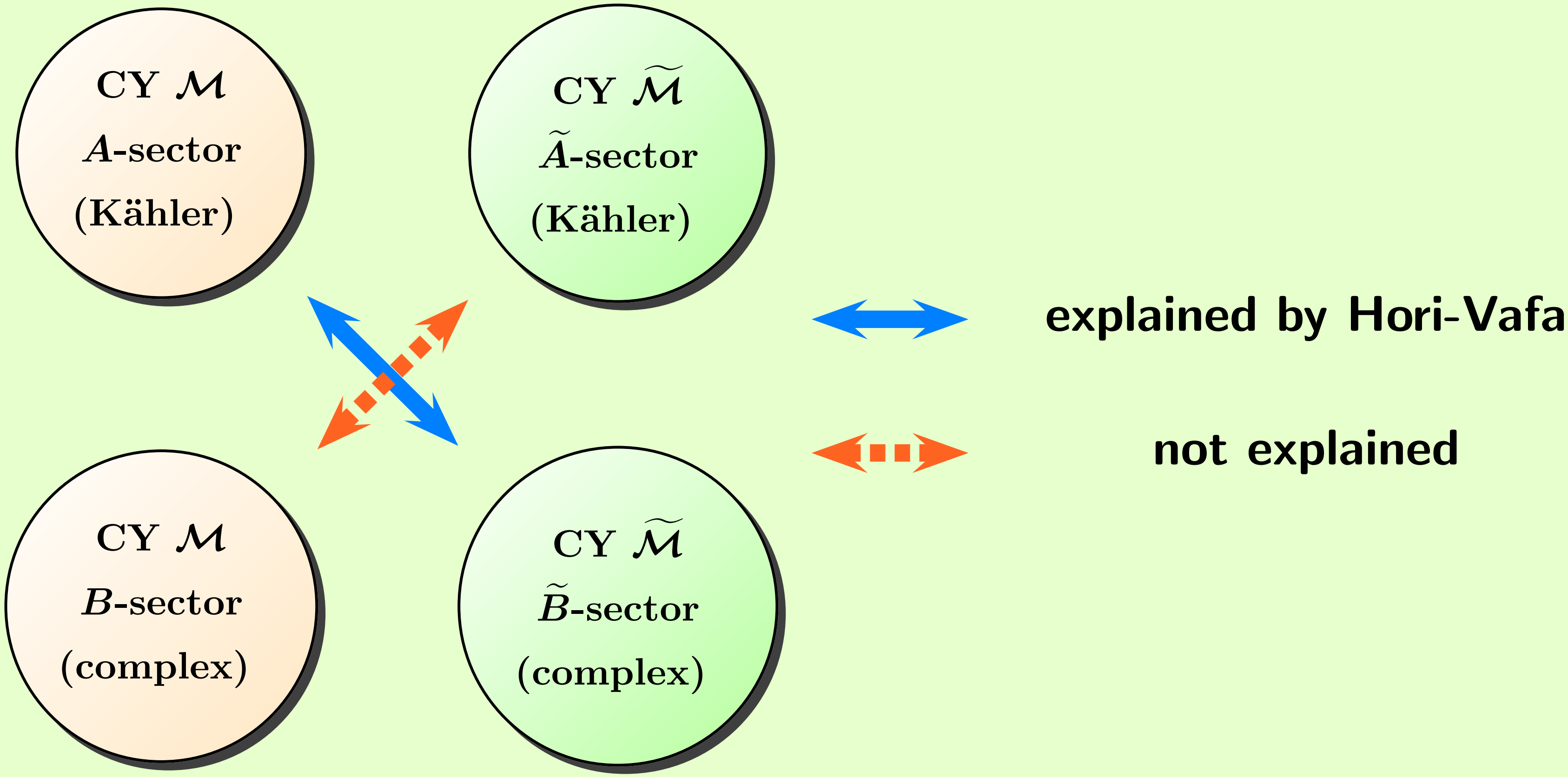
  two orbifolded LG theories including a term with negative power $-k$

  This term represents a gauged WZW model on $SL(2,\mathbb{R})_k/U(1)$ at level k

- We argue that the LG theories in the orignal GLSM can be interpreted as $\mathcal{N}=2$ Liouville theories coupled to LG minimal models

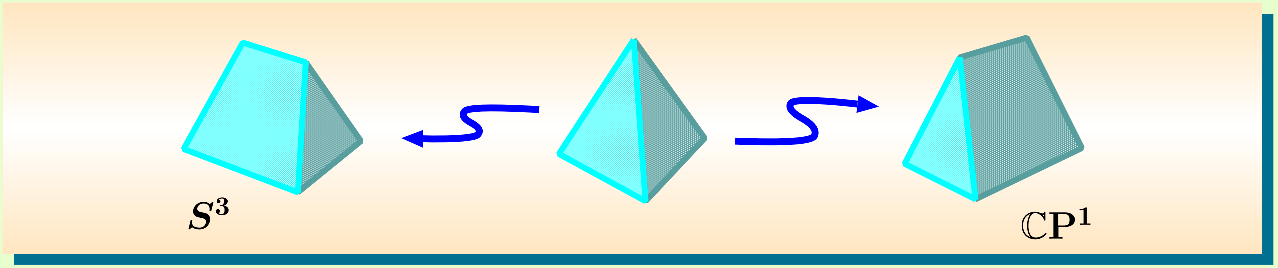

CY sigma model
on noncompact $\mathcal{M}_{\mathrm{CY}}$
CY/LG
$\mathcal{N}=2$ Liouville $\times$ LG
$\mathbb{Z}_\ell$ orbifold theory
FI $\gg 0$
GLSM
FI $\ll 0$
Mirror
T-dual
T-dual
equivalent
$\ell\Sigma \to \frac{\partial}{\partial Y_{P_1}}$
Hori-Vafa
$\ell\Sigma \to \ell\frac{\partial}{\partial t}$
CY sigma model
on noncompact $\widetilde{\mathcal{M}}_\ell$
CY/LG
$SL(2,\mathbb{R})_k/U(1) \times \widetilde{\mathrm{LG}}$
$(\mathbb{Z}_\ell)^N$ orbifold theory

Discussions

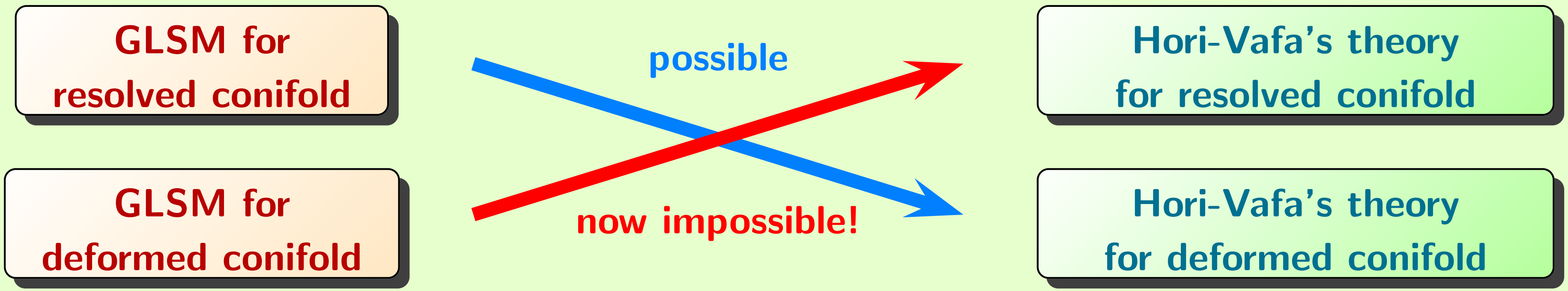


- Hori-Vafa's T-dual theory is only valid when we consider the GLSM without a superpotential or with a superpotential given simply by a homogeneous polynomial such as $W_{\rm GLSM} = P \cdot G_\ell(S)$. Even though the polynomial $G_\ell(S)$ has an additional symmetry, the period integral $\widehat{\Pi}$ cannot recognize the existence of this additional symmetry. Thus the T-dual theory does not map all structures of the CY $\mathcal{M}$ to the mirror geometry completely.

Example: resolved/deformed conifold



- deformed conifold: deformation of complex moduli
  resolved conifold: deformation of Kähler moduli



Derivation of the T-dual Lagrangian

Here we briefly review the T-duality of a generic GLSM without any superpotentials. We start from

$$\mathscr{L}' \ = \ \int \mathrm{d}^4\theta \Big\{ -\frac{1}{e^2}\overline{\Sigma}\Sigma + \sum_a \Big(\mathrm{e}^{2Q_a V + B_a} - \frac{1}{2}(Y_a + \overline{Y}_a) B_a \Big) \Big\} + \Big(\frac{1}{\sqrt{2}} \int \mathrm{d}^2\widetilde{\theta}\, (-\Sigma\, t) + (c.c.) \Big) \, , \tag{1}$$

where Y_a and B_a are twisted chiral superfields and a real superfields B_a.

Integrating out twisted chiral superfields Y_a, we obtain $\overline{D}_+ D_- B_a = D_+ \overline{D}_- B_a = 0$, whose solutions are written in terms of chiral superfields Ψ_a and $\overline{\Psi}_a$ such as $B_a = \Psi_a + \overline{\Psi}_a$. When we substitute them into (1), $\mathscr{L}_{\mathrm{GLSM}}$ appears:

$$\mathscr{L}'\Big|_{B_a = \Psi_a + \overline{\Psi}_a} \ = \ \int \mathrm{d}^4\theta \Big\{ -\frac{1}{e^2}\overline{\Sigma}\Sigma + \sum_a \overline{\Phi}_a\, \mathrm{e}^{2Q_a V} \Phi_a \Big\} + \Big(\frac{1}{\sqrt{2}} \int \mathrm{d}^2\widetilde{\theta}\, (-\Sigma\, t) + (c.c.) \Big) \ \equiv \ \mathscr{L}_{\mathrm{GLSM}} \, , \tag{2}$$

where we re-wrote $\Phi_a = \mathrm{e}^{\Psi_a}$.

On the other hand, when we first integrate out B_a in $\mathscr{L}'$, we obtain $B_a \ = \ -2Q_a V + \log \Big(\dfrac{Y_a + \overline{Y}_a}{2} \Big)$. Let us insert these solutions into (1). By using a deformation $\int \mathrm{d}^4\theta\, Q_a V Y_a = -\dfrac{Q_a}{2} \int \mathrm{d}^2\widetilde{\theta}\, \overline{D}_+ D_- V Y_a = -\dfrac{Q_a}{\sqrt{2}} \int \mathrm{d}^2\widetilde{\theta}\, \Sigma Y_a$, we find that a Lagrangian of twisted chiral superfields appears:

$$\begin{aligned} \mathscr{L}_{\mathrm{T}} \ &= \ \int d^4\theta \Big\{ -\frac{1}{e^2}\overline{\Sigma}\Sigma - \sum_a \Big(\frac{1}{2}(Y_a + \overline{Y}_a) \log(Y_a + \overline{Y}_a) \Big) \Big\} + \Big(\frac{1}{\sqrt{2}} \int \mathrm{d}^2\widetilde{\theta}\, \widetilde{W} + (c.c.) \Big) \, , \\ \widetilde{W} \ &= \ \Sigma \Big(\sum_a Q_a Y_a - t \Big) + \mu \sum_a \mathrm{e}^{-Y_a} \, . \end{aligned}$$

Notice that the twisted superpotential $\widetilde{W}$ is corrected by instanton effects where the instantons are the vortices of the gauge theory. In attempt to analyze a model satisfying $\sum_a Q_a = 0$, the scale parameter μ is omitted by field re-definitions.

linear dilaton CFT and Liouville theory

$$\mathbb{R}^{9,1} = \underbrace{\mathbb{R}^{d-1,1}}_{\text{free SCFT}} \times \underbrace{X^{2n}}_{\text{singular CY}} \sim \mathbb{R}^{d-1,1} \times \underbrace{\mathbb{R}_\phi \times S^1}_{\text{linear dilaton SCFT}} \times \underbrace{\mathcal{M}/U(1)}_{\mathcal{N}=2\text{ Landau-Ginzburg}}$$

linear dilaton: $\Phi = -\frac{Q}{2}\phi$

Landau-Ginzburg: $W_{\text{LG}} = F(Z_a),\ \ F(\lambda^{r_a} Z_a) = \lambda F(Z_a)$

$$c_{\text{total}} = c_d + c_{\text{dilaton}} + c_{\text{LG}} \quad \longrightarrow \quad 15 = \frac{3}{2}d + \left(\frac{3}{2} + 3Q^2\right) + 3\sum_{a=1}^{n+1}(1 - 2r_a)$$

$\mathcal{N} = 2$ "LG" on $\mathbb{R}_\phi \times S^1 \times \mathcal{M}/U(1)$: $W = -\mu Z_0^{-k} + F(Z_a)$

$$k = \frac{1}{r_\Omega} = \frac{2}{Q^2}, \qquad r_\Omega \equiv \textstyle\sum_a r_a - 1$$

linear dilaton SCFT on $\mathbb{R}_\phi \times S^1$ $\equiv$ "LG" with 「$W = -\mu Z_0^{-k}$」

$\equiv$ Kazama-Suzuki model on $SL(2,\mathbb{R})_k/U(1)$

$\overset{\text{T-dual}}{\equiv}$ Liouville theory of charge Q

Strictly, we consider the Euclidean black hole: $SL(2,\mathbb{R})_k/U(1) \longrightarrow [SL(2,\mathbb{C})_k/SU(2)]/U(1)$

Derivation of the twisted geometry

Let us study how to obtain the geometry with $\mathbb{Z}_\ell$-type orbifold symmetry. Replacing $\ell\Sigma$ in $\widehat{\Pi}$ to

$$\ell\Sigma \;\to\; \frac{\partial}{\partial Y_{P_1}} \, ,$$

we can perform the integration of Σ and obtain

$$\widehat{\Pi} \;=\; \int \prod_{i=1}^{N} \mathrm{d}Y_i \, (\mathrm{e}^{-Y_{P_1}} \mathrm{d}Y_{P_1}) \, \mathrm{d}Y_{P_2} \, \delta\Big(\sum_i Y_i - \ell Y_{P_1} - (N-\ell) Y_{P_2} - t \Big) \exp\Big(- \sum_i \mathrm{e}^{-Y_i} - \mathrm{e}^{-Y_{P_1}} - \mathrm{e}^{-Y_{P_2}} \Big) \, . \qquad (3)$$

We perform the re-definitions of the variables Y_i, Y_{P_1} and Y_{P_2}:

$$\mathrm{e}^{-Y_{P_1}} \;=\; \widetilde{P}_1 \, , \quad \mathrm{e}^{-Y_a} \;=\; \widetilde{P}_1 \, U_a \quad \text{for } a = 1, \dots, \ell \, , \qquad \mathrm{e}^{-Y_{P_2}} \;=\; \widetilde{P}_2 \, , \quad \mathrm{e}^{-Y_b} \;=\; \widetilde{P}_2 \, U_b \quad \text{for } b = \ell+1, \dots, N \, .$$

Substituting these re-defined variables into (3), we continue the calculation:

$$\begin{aligned}
\widehat{\Pi} \;&=\; \int \prod_{i=1}^{N} \Big(\frac{\mathrm{d}U_i}{U_i} \Big) \, \mathrm{d}\widetilde{P}_1 \Big(\frac{\mathrm{d}\widetilde{P}_2}{\widetilde{P}_2} \Big) \, \delta\Big(\log\Big(\prod_i U_i \Big) + t \Big) \exp\Big\{ - \widetilde{P}_1 \Big(\sum_{a=1}^{\ell} U_a + 1 \Big) - \widetilde{P}_2 \Big(\sum_{b=\ell+1}^{N} U_b + 1 \Big) \Big\} \\
&=\; \int \prod_i \Big(\frac{\mathrm{d}U_i}{U_i} \Big) \, \mathrm{d}\widetilde{P}_2 \, \mathrm{d}u \, \mathrm{d}v \, \delta\Big(\log\Big(\prod_i U_i \Big) + t \Big) \, \delta\Big(\sum_a U_a + 1 \Big) \exp\Big\{ - \widetilde{P}_2 \Big(\sum_b U_b + 1 - uv \Big) \Big\} \\
&=\; \int \prod_i \Big(\frac{\mathrm{d}U_i}{U_i} \Big) \, \mathrm{d}u \, \mathrm{d}v \, \delta\Big(\log\Big(\prod_i U_i \Big) + t \Big) \, \delta\Big(\sum_a U_a + 1 \Big) \, \delta\Big(\sum_b U_b + 1 - uv \Big) \, , \qquad (4)
\end{aligned}$$

where we introduced new variables u and v taking values in $\mathbb{C}$ and used a following equation

$$\frac{1}{\widetilde{P}_2} \;=\; \int \mathrm{d}u \, \mathrm{d}v \, \exp\big(\widetilde{P}_2 \, uv \big) \, .$$

It is obvious that (4) still includes a non-canonical integral measure. Thus we perform further re-definitions such as

$$U_a = \mathrm{e}^{-t/\ell} \frac{Z_a^\ell}{Z_1 \cdots Z_N}, \qquad U_b = Z_b^\ell .$$

Note that the period integral (4) is invariant under the following transformations acting on the new variables Z_i:

$$Z_a \mapsto \lambda \omega_a Z_a , \quad Z_b \mapsto \omega_b Z_b , \qquad \omega_a^\ell = \omega_b^\ell = \omega_1 \cdots \omega_N = 1 ,$$

where λ is an arbitrary number taking in $\mathbb{C}^*$. The ω_i come from the shift symmetry of the original variables $Y_i \equiv Y_i + 2\pi i$. Combining these transformations we find that $\widehat{\Pi}$ has $\mathbb{C}^* \times (\mathbb{Z}_\ell)^{N-2}$ symmetries. Substituting Z_i into (4), we obtain

$$\widehat{\Pi} = \int \frac{1}{\mathrm{vol.}(\mathbb{C}^*)} \prod_{i=1}^N \mathrm{d}Z_i \, \mathrm{d}u \, \mathrm{d}v \, \delta\Big(\sum_{a=1}^\ell Z_a^\ell + \mathrm{e}^{t/\ell} Z_1 \cdots Z_N \Big) \, \delta\Big(\sum_{b=\ell+1}^N Z_b^\ell + 1 - uv \Big) ,$$

which indicates that the resulting mirror geometry is described by

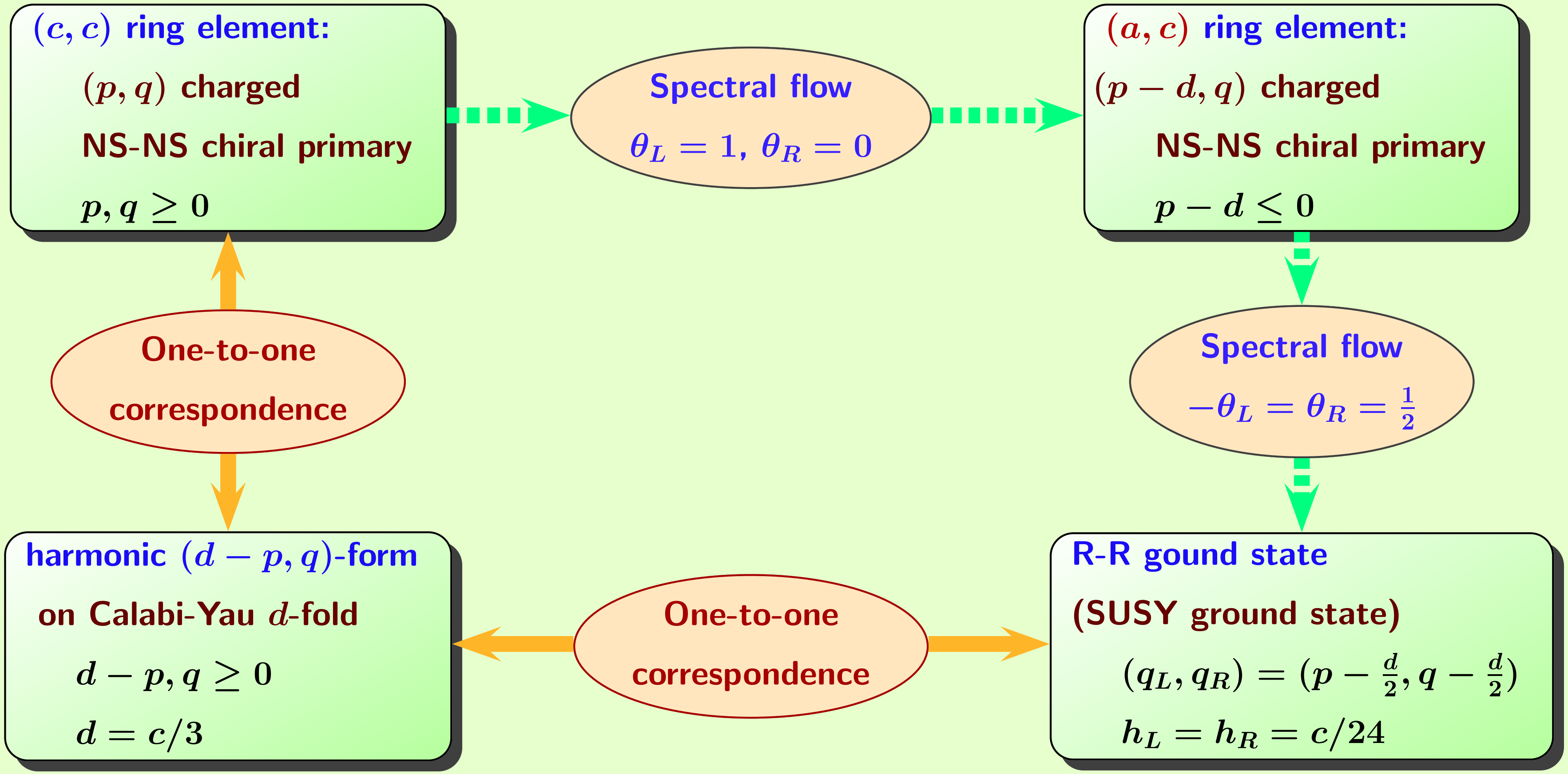
$$\widetilde{\mathcal{M}}_\ell = \Big\{ (Z_i; u, v) \in \mathbb{C}^{N+2} \Big| \big\{ \mathcal{F}(Z_i) = 0 \big\} / \mathbb{C}^* , \ \mathcal{G}(Z_b; u, v) = 0 \Big\} \Big/ (\mathbb{Z}_\ell)^{N-2} ,$$

$$\mathcal{F}(Z_i) = \sum_{a=1}^\ell Z_a^\ell + \psi Z_1 \cdots Z_\ell , \qquad \mathcal{G}(Z_b; u, v) = \sum_{b=\ell+1}^N Z_b^\ell + 1 - uv , \qquad \psi = \mathrm{e}^{t/\ell} Z_{\ell+1} \cdots Z_N .$$

This is an $(N-1)$-dimensional complex manifold.

The equation $\mathcal{F}(Z_i) = 0$ denotes that the complex variables Z_a consist of the degree ℓ hypersurface in the projective space: $\mathbb{C}\mathrm{P}^{\ell-1}[\ell]$. This subspace itself is a compact CY manifold, which is parametrized by a parameter ψ which is subject to the equation $\mathcal{G}(Z_b; u, v) = 0$. Moreover we can also interpret that the total space is a noncompact CY manifold whose compact directions are described by Z_i, while the variables u and v run in the noncompact directions under the equations.

$\mathcal{N} = (2,2)$ SCFT and (compact) Calabi-Yau geometry



relation among the NS-NS chiral primary states, R-R ground states and harmonic forms