

日本物理学会第64回年次大会 at 立教学院池袋キャンパス (2009年3月30日)

Realization of AdS Vacua in Attractor Mechanism on Generalized Geometry

arXiv:0810.0937 [hep-th]

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いきなり結論

(拡張された)幾何でコンパクト化された超弦理論



4次元 $\mathcal{N} = 1$ 超重力の Superpotential (moduliの関数)の判別式が

正 : SUSY AdS 真空 負 : SUSY 平坦真空

拡張された幾何とは？

何故そんなものを考えるのか？

ご利益

全ての超重力理論を超弦理論から導出できる可能性

NS-NS, R-R fluxes を自然に取り込んだ真空解が得られる

Superpotential を自然に導出できる

余分なゼロ質量場を追い出す事ができる

高次元理論を通常^の幾何でコンパクト化しても
登場しない超重力理論がある(らしい)



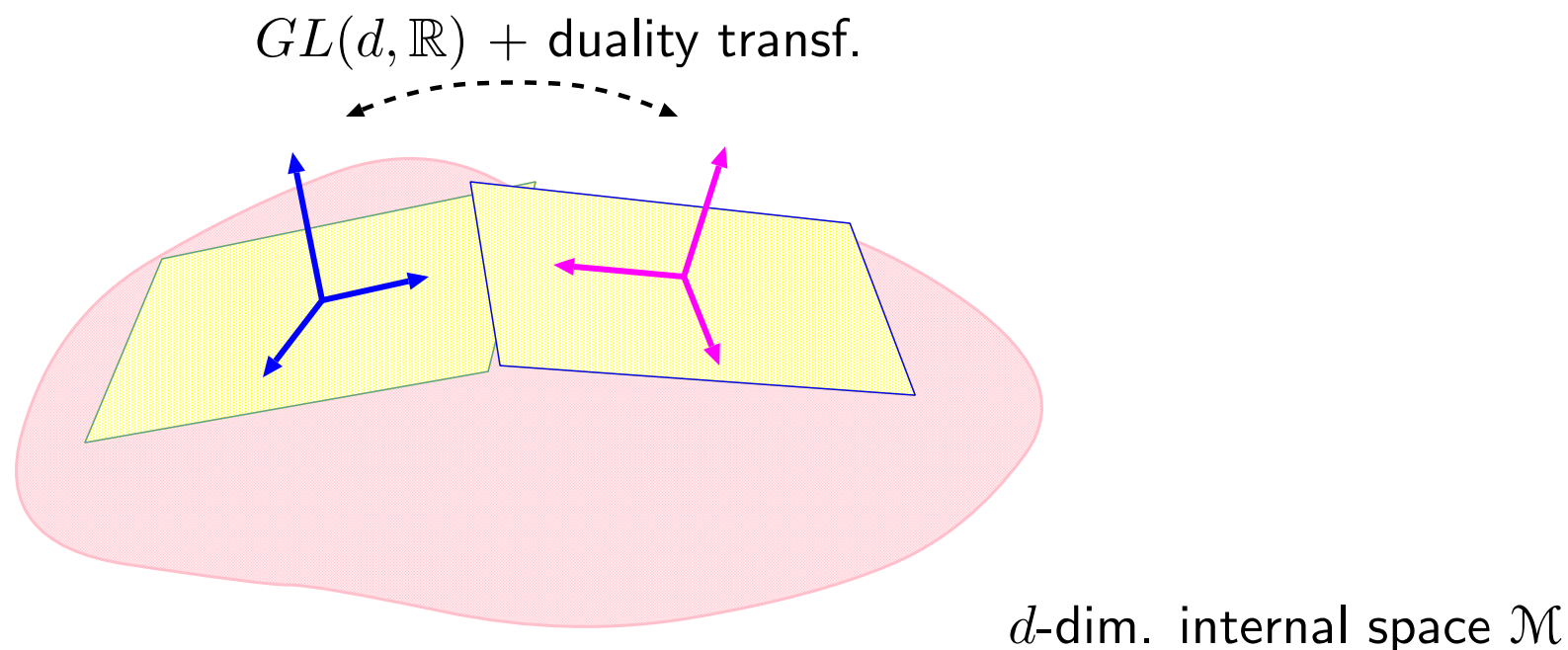
全ての超重力理論は超弦理論に由来すると考えるならば
超弦理論を通常でない幾何 (Nongeometric Background) で
コンパクト化したと考えてみよう

質問: Nongeometric Background (拡張された幾何)とは?

構造群 = Diffeo群 ($GL(d, \mathbb{R})$) + 双対変換群 ($O(d, d; \mathbb{Z})$, S-双対変換, U-双対変換, etc.)



弦理論の双対性に起因すると考える



この分野における私の貢献

IIA 型超弦理論を

6D $SU(3) \times SU(3)$ Generalized geometry でコンパクト化



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判別式を見るだけで宇宙項の大きさを議論できる

$$\begin{aligned}
 V &= e^K \left(K^{\mathcal{M}\bar{\mathcal{N}}} D_{\mathcal{M}} \mathcal{W} \overline{D_{\bar{\mathcal{N}}} \mathcal{W}} - 3|\mathcal{W}|^2 \right) + \frac{1}{2} |D^a|^2 \\
 &\equiv V_{\mathcal{W}} + V_D
 \end{aligned}$$

Search of vacua $\partial_{\mathcal{P}} V|_* = 0$

$V_* > 0$: de Sitter space (non-SUSY)

$V_* = 0$: Minkowski space

$V_* < 0$: Anti-de Sitter space

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$$0 = \partial_{\mathcal{P}} V_{\mathcal{W}} = e^K \left\{ K^{\mathcal{M}\bar{\mathcal{N}}} D_{\mathcal{P}} D_{\mathcal{M}} \mathcal{W} \overline{D_{\bar{\mathcal{N}}} \mathcal{W}} + \partial_{\mathcal{P}} K^{\mathcal{M}\bar{\mathcal{N}}} D_{\mathcal{M}} \mathcal{W} \overline{D_{\bar{\mathcal{N}}} \mathcal{W}} - 2\overline{\mathcal{W}} D_{\mathcal{P}} \mathcal{W} \right\}$$

$$0 = \partial_{\mathcal{P}} V_D \quad \rightarrow \quad D^a = 0$$

Consider the SUSY condition $D_{\mathcal{P}} \mathcal{W} \equiv (\partial_{\mathcal{P}} + \partial_{\mathcal{P}} K) \mathcal{W} = 0$ in various cases.

1. Set a simple prepotential: $\mathcal{F} = D_{abc} \frac{X^a X^b X^c}{X^0}$
2. Consider the simplest model: single modulus t of Φ_+ (and U of Φ_-)

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The superpotential is reduced to

$$\mathcal{W} = \mathcal{W}^{\text{RR}} + U \mathcal{W}^{\text{Q}}$$

$$\mathcal{W}^{\text{RR}} = m_{\text{RR}}^0 t^3 - 3 m_{\text{RR}} t^2 + e_{\text{RR}} t + e_{\text{RR}0}$$

$$\mathcal{W}^{\text{Q}} = p_0^0 t^3 - 3 p_0 t^2 - e_0 t - e_{00}$$

Consider the SUSY condition:

$$D_t \mathcal{W} = 0 \quad \rightarrow \quad 0 = D_t \mathcal{W}^{\text{RR}} + U D_t \mathcal{W}^{\text{Q}}$$

$$D_U \mathcal{W} = 0 \quad \rightarrow \quad 0 = \frac{i}{\text{Im}U} \left(\mathcal{W}^{\text{RR}} + \text{Re}U \mathcal{W}^{\text{Q}} \right)$$

The discriminant of the superpotential \mathcal{W}^{RR} (and \mathcal{W}^{Q}) governs the SUSY solutions.

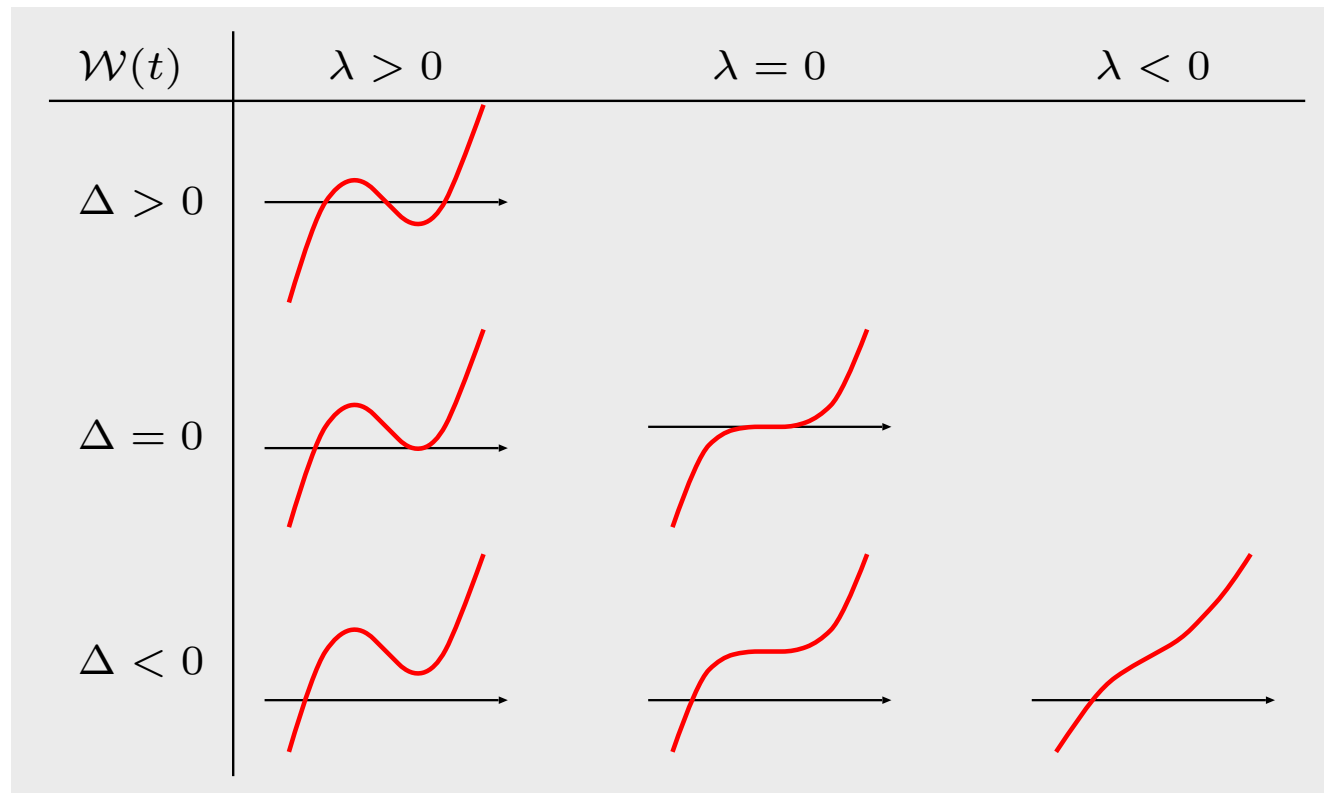
► Discriminant of cubic equation

Consider a cubic function and its derivative:
$$\begin{cases} \mathcal{W}(t) = at^3 + bt^2 + ct + d \\ \partial_t \mathcal{W}(t) = 3at^2 + 2bt + c \end{cases}$$

Discriminants $\Delta(\mathcal{W})$ and $\Delta(\partial_t \mathcal{W})$ are

$$\Delta(\mathcal{W}) \equiv \Delta = -4b^3d + b^2c^2 - 4ac^3 + 18abcd - 27a^2d^2$$

$$\Delta(\partial_t \mathcal{W}) \equiv \lambda = 4(b^2 - 3ac)$$



$\Delta^{\text{RR}} > 0$ case: always $\lambda^{\text{RR}} > 0$, and exists a zero point: $D_t \mathcal{W}^{\text{RR}} = 0$

$$\begin{aligned}
 D_t \mathcal{W}^{\text{RR}}|_* &= 0 \\
 t_*^{\text{RR}} &= \frac{6(3m_{\text{RR}}^0 e_{\text{RR}0} + m_{\text{RR}} e_{\text{RR}})}{\lambda^{\text{RR}}} - 2i \frac{\sqrt{3\Delta^{\text{RR}}}}{\lambda^{\text{RR}}} \\
 \mathcal{W}_*^{\text{RR}} &= -\frac{24\Delta^{\text{RR}}}{(\lambda^{\text{RR}})^3} \left(36(m_{\text{RR}})^3 + 36(m_{\text{RR}}^0)^2 e_{\text{RR}0} - 3m_{\text{RR}} \lambda^{\text{RR}} - 4i m_{\text{RR}}^0 \sqrt{3\Delta^{\text{RR}}} \right)
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 \end{aligned}$$

$\Delta^{\text{RR}} < 0$ case: only $\lambda^{\text{RR}} < 0$ is physically allowed, and exists a zero point: $\mathcal{W}^{\text{RR}} = 0$

$$\begin{aligned}
 \mathcal{W}_*^{\text{RR}} &= m_{\text{RR}}^0 (t_* - e)(t_* - \alpha)(t_* - \bar{\alpha}) = 0, \quad t_* = \alpha^{\text{RR}} = \alpha_1 + i\alpha_2 \\
 \alpha_1 &= \frac{\lambda^{\text{RR}} + F^{2/3} + 12m_{\text{RR}} F^{1/3}}{12m_{\text{RR}}^0 F^{1/3}} \\
 (\alpha_2)^2 &= \frac{1}{m_{\text{RR}}^0} \left(e_{\text{RR}} - 6m_{\text{RR}} \alpha_1 + 3m_{\text{RR}}^0 (\alpha_1)^2 \right) \\
 e &= -\frac{1}{m_{\text{RR}}^0} \left(-3m_{\text{RR}} + 2m_{\text{RR}}^0 \alpha_1 \right) \\
 F &= 108(m_{\text{RR}}^0)^2 e_{\text{RR}0} + 12m_{\text{RR}}^0 \sqrt{-3\Delta^{\text{RR}}} + 108(m_{\text{RR}})^3 - 9m_{\text{RR}} \lambda^{\text{RR}} \\
 D_t \mathcal{W}^{\text{RR}}|_* &= 2i m_{\text{RR}}^0 (e - \alpha^{\text{RR}}) \alpha_2
 \end{aligned}$$

... Analysis of \mathcal{W}^{Q} is also discussed.

Three types of solutions to satisfy $0 = D_t \mathcal{W}^{\text{RR}} + U D_t \mathcal{W}^{\text{Q}}$ and $0 = \mathcal{W}^{\text{RR}} + \text{Re}U \mathcal{W}^{\text{Q}}$:

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- SUSY AdS vacuum: moduli are (mostly) stabilized

$$\Delta^{\text{RR}} > 0, \quad \Delta^{\text{Q}} > 0; \quad D_t \mathcal{W}^{\text{RR}}|_* = 0 = D_t \mathcal{W}^{\text{Q}}|_*$$

$$t_*^{\text{RR}} = t_*^{\text{Q}}, \quad \text{Re} U_* = -\frac{\mathcal{W}_*^{\text{RR}}}{\mathcal{W}_*^{\text{Q}}}$$

$$V_* = -3e^K |\mathcal{W}_*|^2 = -\frac{4}{[\text{Re}(\mathcal{C}\mathcal{G}_0)]^2} \sqrt{\frac{\Delta^{\text{Q}}}{3}} \ll \mathcal{O}(1)$$

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- SUSY Minkowski vacuum: moduli are stabilized

$$\Delta^{\text{RR}} < 0, \quad \Delta^{\text{Q}} < 0; \quad \mathcal{W}_*^{\text{RR}} = 0 = \mathcal{W}_*^{\text{Q}}$$

$$\alpha^{\text{RR}} = \alpha^{\text{Q}}, \quad U_* = -\frac{D_t \mathcal{W}^{\text{RR}}|_*}{D_t \mathcal{W}^{\text{Q}}|_*} \neq 0$$

$$V_* = 0$$

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- SUSY Minkowski vacuum: moduli are stabilized

$$\begin{aligned} \Delta^{\text{RR}} < 0, \quad \Delta^{\text{Q}} < 0; \quad \mathcal{W}_*^{\text{RR}} = 0 = \mathcal{W}_*^{\text{Q}} \\ \hline \alpha^{\text{RR}} = \alpha^{\text{Q}}, \quad U_* = -\frac{D_t \mathcal{W}^{\text{RR}}|_*}{D_t \mathcal{W}^{\text{Q}}|_*} \neq 0 \\ V_* = 0 \end{aligned}$$

- SUSY AdS vacua, but moduli t and U are not fixed: non-stabilized point

$$U = -\frac{D_t \mathcal{W}^{\text{RR}}(t)}{D_t \mathcal{W}^{\text{Q}}(t)}, \quad \text{Re}U = -\frac{\mathcal{W}^{\text{RR}}(t)}{\mathcal{W}^{\text{Q}}(t)}$$