

# **Semi-doubled GLSM for Five-branes of Codimension Two**

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# **1 . 5-branes of codim-2**

F-string  
(1+1)-dim.



NS5-brane  
(1+5)-dim.



$B_{MN}$



Compactify on  $T^2$

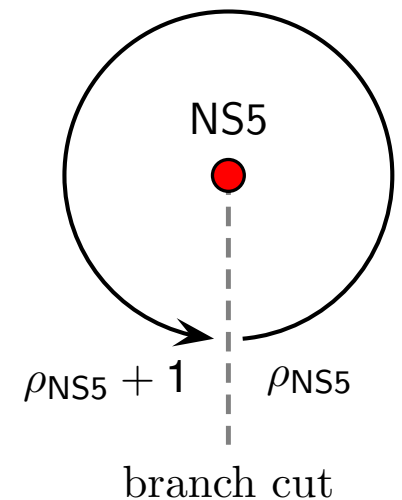


NS5-brane of codim-2 in 8D string theory = defect NS5-brane

$$ds^2 = \underbrace{dx_{012345}^2}_{\text{NS5-brane}} + H \left[ \underbrace{(dr)^2 + r^2(d\theta)^2}_{\text{codim-2}} + \underbrace{(d\varphi)^2 + (d\vartheta)^2}_{T^2} \right]$$

$$B_{\varphi\vartheta} = \frac{\theta}{2\pi}, \quad e^{2\phi} = H = \frac{1}{2\pi} \log \frac{\Lambda}{r}$$

$$\rho_{\text{NS5}} = B_{\varphi\vartheta} + iH = \frac{\theta}{2\pi} + \frac{i}{2\pi} \log \frac{\Lambda}{r}$$

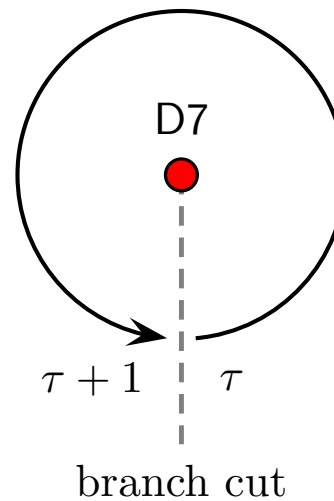


similar to D7-brane in type IIB string

D7-brane in 10D type IIB string is coupled to

$$\tau(z) \equiv C + ie^{-\phi} = \frac{\theta}{2\pi} + \frac{i}{2\pi} \log \frac{\Lambda}{\rho}$$

magnetic “charge” (**monodromy**):  $\tau \rightarrow \tau + 1$



feature of  $\tau \rightarrow$  F-theory

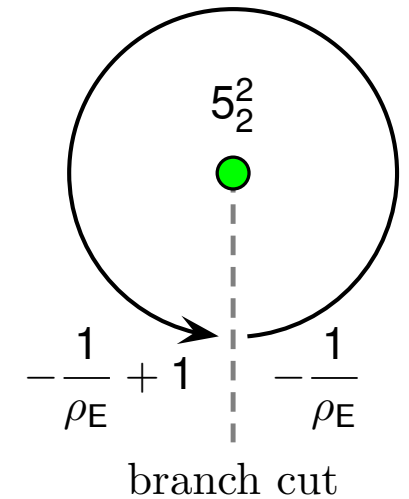
Vafa (1996)

defect NS5-brane  $\xrightarrow{\text{T-dual along } T^2}$  exotic  $5_2^2$ -brane

$$ds^2 = \underbrace{dx_{012345}^2}_{5_2^2\text{-brane}} + H \underbrace{[(dr)^2 + r^2(d\theta)^2]}_{\text{codim-2}} + \frac{H}{K} \underbrace{[(d\tilde{\varphi})^2 + (d\tilde{\vartheta})^2]}_{\text{dual } T^2}$$

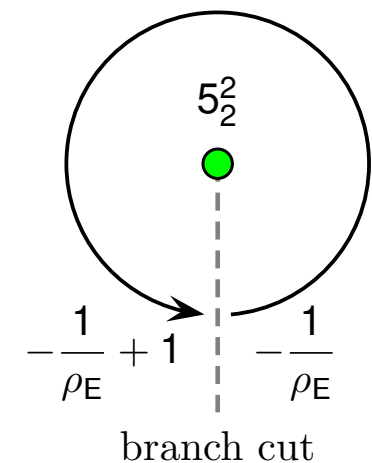
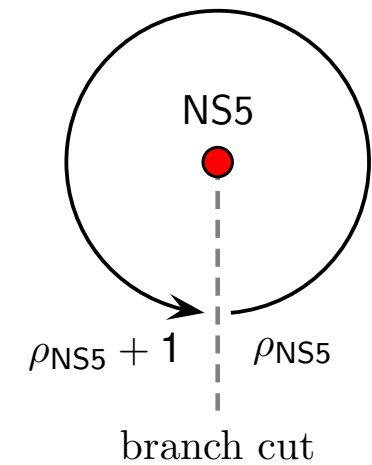
$$B_{\varphi\vartheta} = -\frac{\theta}{2\pi K}, \quad e^{2\phi} = \frac{H}{K}, \quad K = H^2 + \left(\frac{\theta}{2\pi}\right)^2$$

$$\rho_E = -\left[\frac{\theta}{2\pi} + \frac{i}{2\pi} \log \frac{\Lambda}{r}\right]^{-1} = -\frac{1}{\rho_{\text{NS5}}}$$



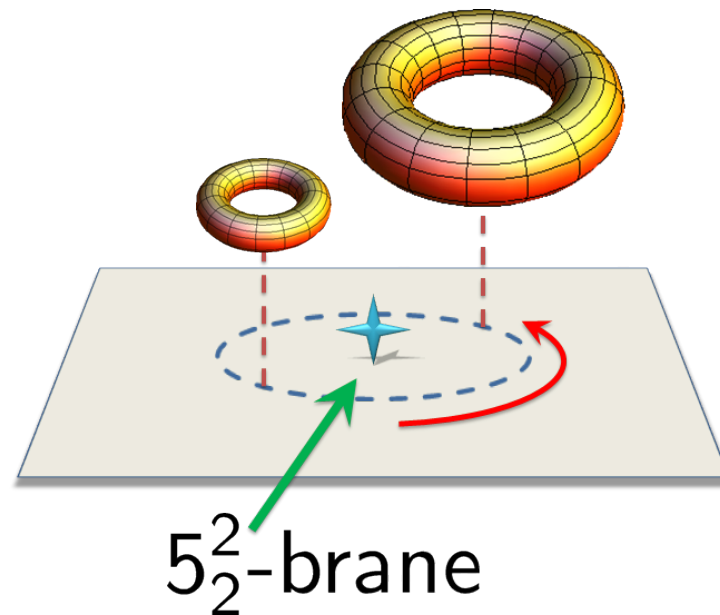
codim-2 branes	D7	defect NS5
$D$	10B	8 (+ $T^2$ )
coupled to	$\tau = C + ie^{-\phi}$	$\rho_{\text{NS5}} = \text{Kähler str. of } T^2$
monodromy	$SL(2, \mathbb{Z})_S$	$SL(2, \mathbb{Z})_K$
pair	NS7 ( $[0, 1]7$ or $7_3$ )	$5_2^2$

$SL(2, \mathbb{Z})_K \subset SO(2, 2; \mathbb{Z}) : \text{T-duality along } T^2$



We investigated 5-branes of codim-2 in string worldsheet sigma model,  
and understood the **classical** background geometry.

We also investigated the **worldsheet instanton corrections along  $S^1 \subset T^2$**   
in terms of ANO vortex corrections **in the UV gauge theory (GLSM)**.



How can we obtain the **full corrections along  $T^2$**  ?



## Worksheet instantons on

defect NS5 : point-like instantons, i.e., small instantons arrayed along  $T^2$

→ deforms background geometry

→ **KK-modes** along  $T^2$

Witten (1996); Tong (2002)

$5_2^2$ -brane : disk instantons

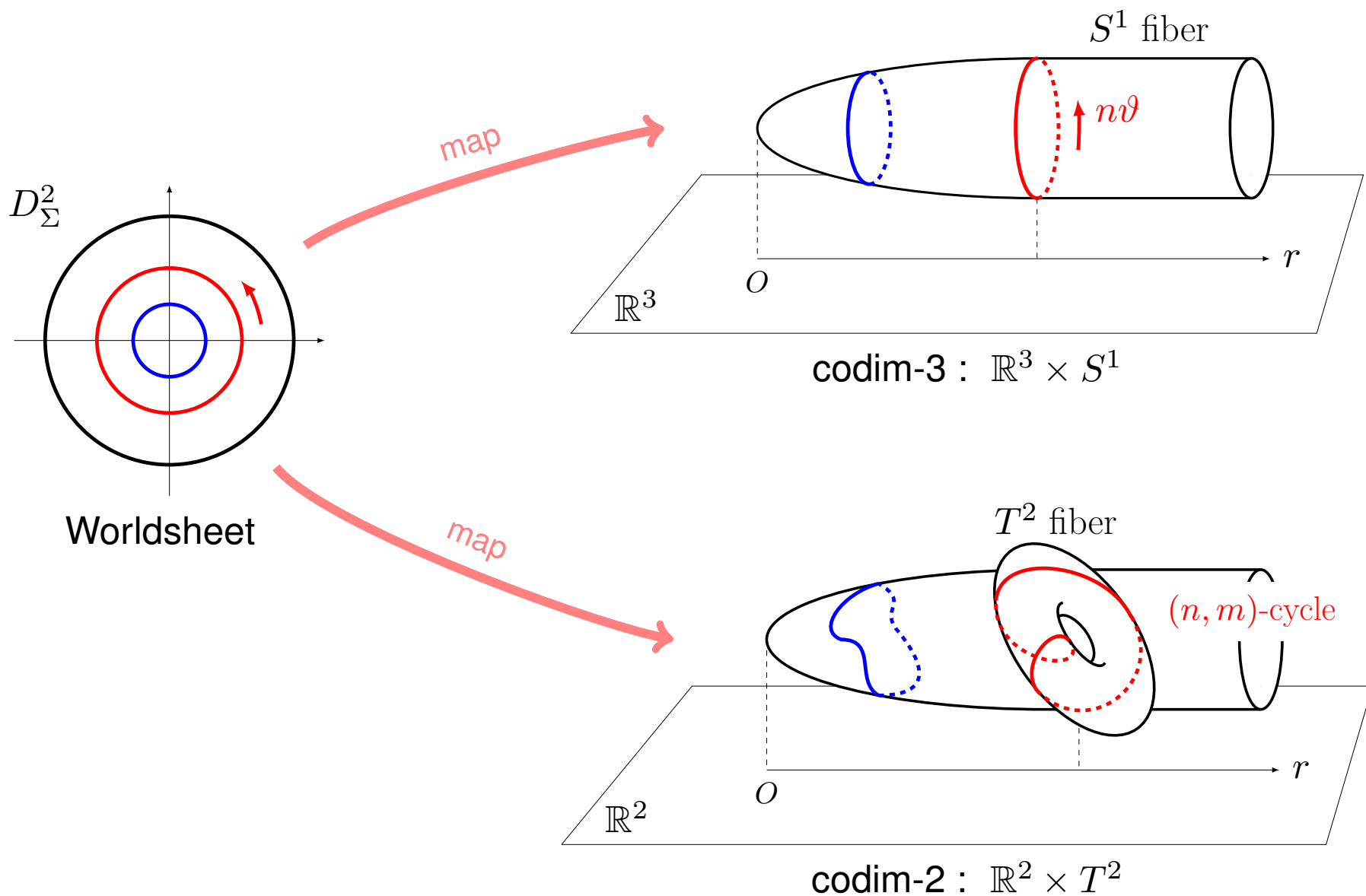
→ does not appear as deformation of background geometry

→ **Winding modes** along  $T^2$

Gregory, Harvey, Moore (1997); Tong (2002); Harvey, Jensen (2005); Okuyama (2005); TK, Sasaki (2013)

analyzed in our DFT computation (TK, Sasaki, Shiozawa arXiv:1803.11087)

# Worldsheet instantons as disk instantons



# 2. Formulation

2D  $\mathcal{N} = (2, 2)$  superfield formalism

$$e^{\int \vartheta F}$$

theta-angle in 2D path integral

$$e^{\int \vartheta F}$$

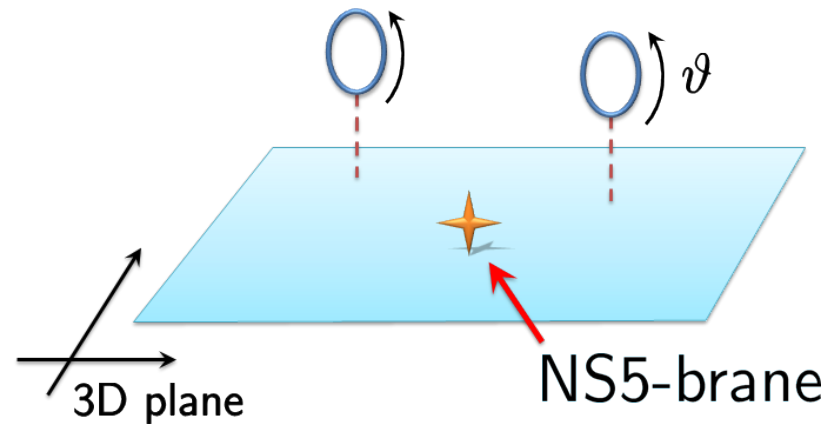
is provided by twisted superpotential

$$e^{\int \Theta \Sigma}$$

$$\Theta, \Sigma = \bar{D}_+ D_- V : \text{twisted chiral}$$

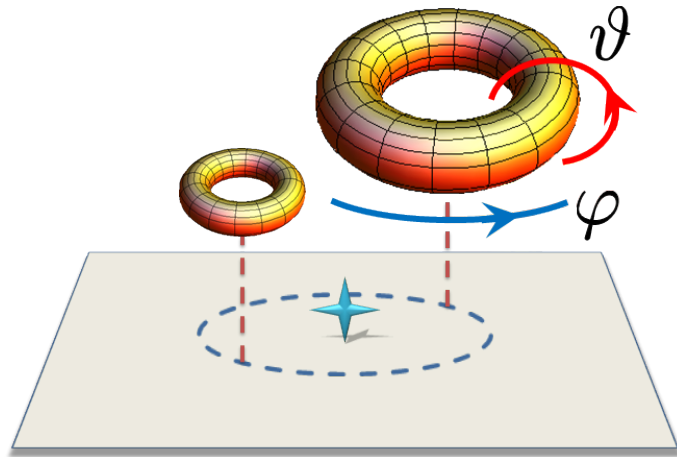
$\vartheta$ , as imaginary part of  $\Theta$ , receives corrections by Chern class.

This is applied to worldsheet instanton corrections of 5-branes of **codim-3**.



Tong (2002); Harvey, Jensen (2005); Okuyama (2005)

Apply this to gauge theory for 5-brane of codim-2 with  $S^1_{\vartheta} \times S^1_{\varphi} = T^2$ .



$$\rightarrow e^{\int \vartheta F + \int \varphi \hat{F}}$$

$$\rightarrow e^{\int \Theta \Sigma + i \int \Theta \hat{\Sigma}}$$

Worldsheet instantons can be mapped to corrections to  $T^2$ .

Available duality transformation w/ irreducible superfields ?

Roček, Verlinde (1991); Hori, Vafa (2000)

$$-\frac{1}{g^2} \int d^4\theta |\Theta|^2 - \int d^2\tilde{\theta} \Theta (\Sigma + i\hat{\Sigma}) + \text{h.c.}$$

dualize imaginary part  $\downarrow$  twisted chiral  $\Theta \rightarrow$  chiral  $\Gamma$

$$+\frac{g^2}{2} \int d^4\theta (\Gamma + \bar{\Gamma} + 2(V + i\hat{V}))^2$$

dualize real part  $\downarrow$  chiral  $\Gamma \rightarrow$  twisted chiral

??

use **reducible** superfields

Grisaru, Massar, Sevrin, Troost (1998); TK (2015)

Duality transformation by reducible superfield :

$$\mathcal{L} = \int d^4\theta \left( -\frac{1}{g^2} |R|^2 - 2(R + \bar{R})V - 2i(R - \bar{R})\hat{V} - RL - \bar{R}\bar{L} \right)$$

- integrate out  $L$  :  $\mathcal{L} = -\frac{1}{g^2} \int d^4\theta |\Theta|^2 - \int d^2\tilde{\theta} \Theta (\Sigma + i\hat{\Sigma}) + \text{h.c.}$

$\Updownarrow$  dual ( $1/g^2 \leftrightarrow g^2$ )

- integrate out  $R$  :  $\mathcal{L} = +g^2 \int d^4\theta |L + 2(V + i\hat{V})|^2$

$L$  : twisted linear superfield  $0 = \bar{D}_+ D_- L$

$$L = X + \bar{W} + Y$$

$X, W$  : chiral,  $Y$  : twisted chiral

Grisaru, Massar, Sevrin, Troost (1998); TK (2015)



# 3. **Worksheet instantons**

Intrinsic part :

$$S = \int d^2\sigma \left\{ \frac{1}{2e^2} F_{12}^2 + |D_m q|^2 + \frac{e^2}{2} (|q|^2 - t)^2 + i\vartheta F_{12} \right. \\ \left. + \frac{1}{2\hat{e}^2} \hat{F}_{12}^2 + |D_m p|^2 + \frac{\hat{e}^2}{2} (|p|^2 - \hat{t})^2 + i\varphi \hat{F}_{12} \right\} + \dots$$

$$\geq |\vec{t}|^2 \sqrt{|\vec{Q}|^2} - i\vec{\vartheta} \cdot \vec{Q}$$

$$\vec{\vartheta} = (\vartheta, \varphi), \quad \vec{t} = (t, \hat{t}), \quad \vec{Q} = \left( - \int d^2\sigma F_{12}, - \int d^2\sigma \hat{F}_{12} \right)$$

Chern numbers

BPS eqs are given as Abrikosov-Nielsen-Olesen (ANO) vortex eq.

$$F_{12} \mp e^2 (|q|^2 - t) = 0, \quad (D_1 \pm iD_2)q = 0$$

- Semi-doubled NLSM (classical)

$$\mathcal{L} = -\frac{H}{2} \{ (\partial_m x)^2 + (\partial_m y)^2 + (\partial_m \varphi)^2 + (\partial_m \vartheta)^2 \} + \Omega \epsilon^{mn} (\partial_m \varphi) (\partial_n \vartheta) \\ + \epsilon^{mn} (\partial_m \varphi) (\partial_n \tilde{\gamma}) + \epsilon^{mn} (\partial_m \vartheta) (\partial_m \gamma)$$

provides string sigma model for defect NS5 and  $5_2^2$  (skip)

$$\text{with } H = \log \frac{\Lambda}{r}, \quad \Omega = \arctan \left( \frac{y}{x} \right), \quad r = \sqrt{x^2 + y^2}$$

- Worldsheet instantons originate from ANO vortices

$$H \rightarrow H + \sum_{(n,m) \neq (0,0)} e^{in\vartheta} e^{im\varphi} e^{-r\sqrt{n^2+m^2}}$$

Witten (1993); Morrison, Plesser (1994); Schroers (1996)

## Worldsheet instantons on

defect NS5 : point-like instantons, i.e., small instantons arrayed along  $T^2$   
→ deforms background geometry  
→ **KK-modes** along  $T^2$

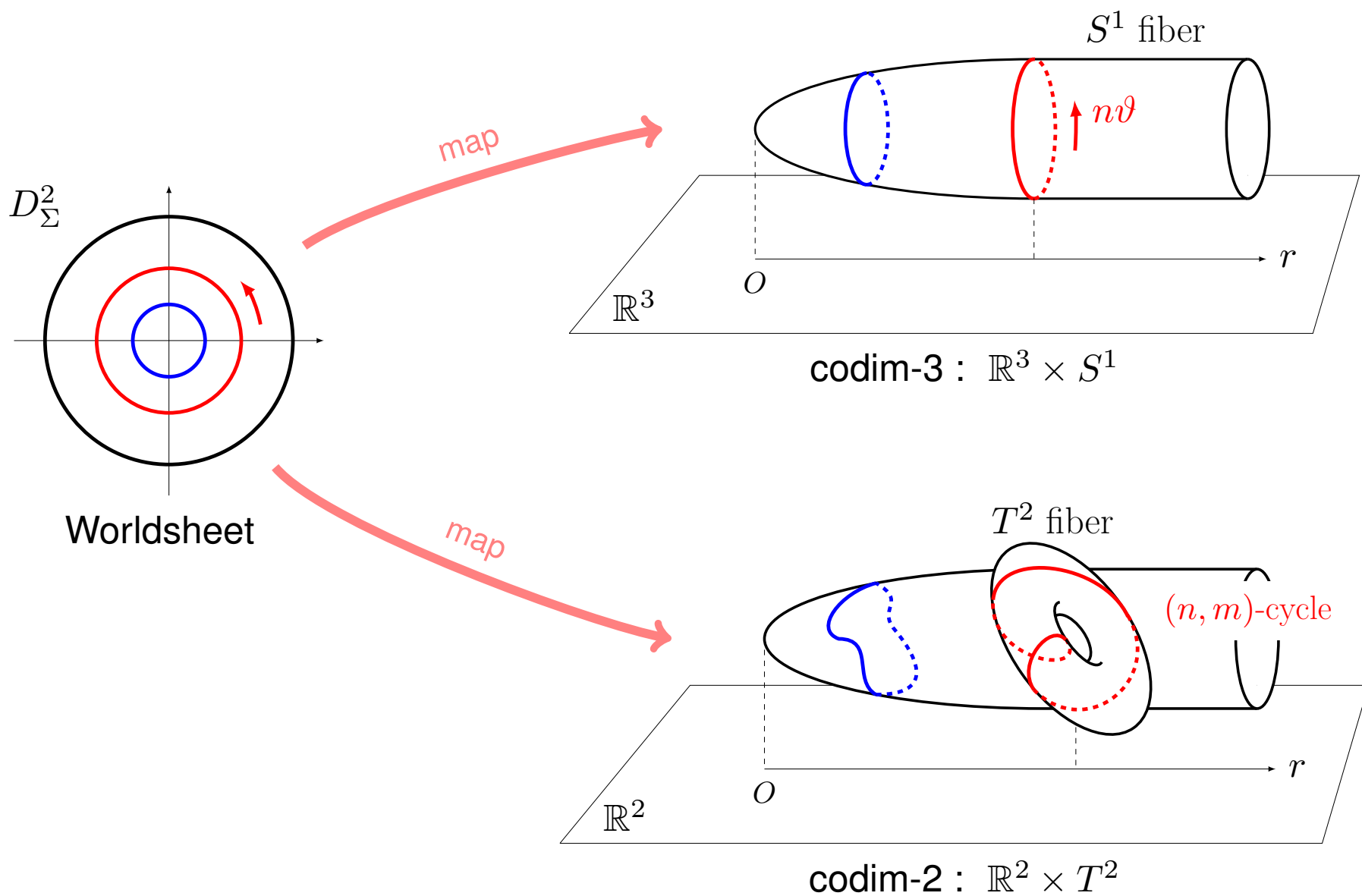
Witten (1996); Tong (2002)

$5_2^2$ -brane : disk instantons  
→ does not appear as deformation of background geometry  
→ **Winding modes** along  $T^2$

Gregory, Harvey, Moore (1997); Tong (2002); Harvey, Jensen (2005); Okuyama (2005); TK, Sasaki (2013)

confirmed our DFT computation (TK, Sasaki, Shiozawa arXiv:1803.11087)

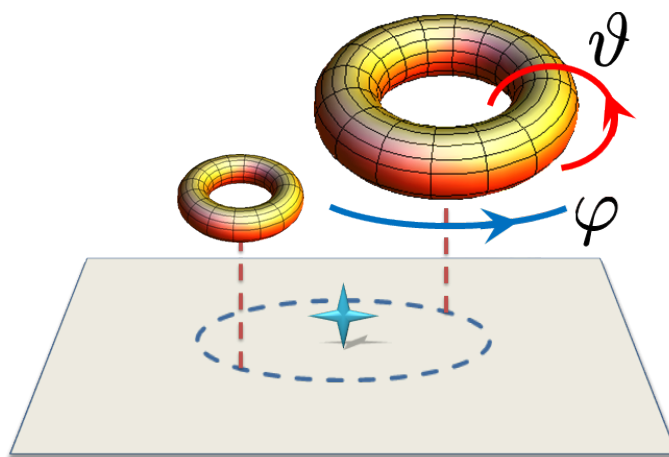
# Worldsheet instantons as disk instantons



# 4. **Summary**

- We captured string worldsheet instanton corrections via ANO vortices in GLSM.
- These instanton corrections are understood as corrections by KK-modes and/or winding modes **along  $T^2$** .

To realize above, we developed duality transformations by **reducible** superfield.



**Thanks**

# Appendix



# **A.** **Duality transformations**

Duality transformation :

$$\mathcal{L} = -\frac{1}{2g^2}(\partial_m\vartheta)^2 + \epsilon^{mn}(\partial_m\vartheta)(\partial_n\gamma)$$

- integrate out  $\gamma$  :  $\mathcal{L} = -\frac{1}{2g^2}(\partial_m\vartheta)^2$
- integrate out  $\vartheta$  :  $\mathcal{L} = -\frac{g^2}{2}(\partial_m\gamma)^2$

It is well known if **only  $\vartheta$  in  $\Theta$**  is dualized.

Roček, Verlinde (1991); Hori, Vafa (2000)

However, how to dualize **both  $\vartheta$  and  $\varphi$**  in  $\Theta$  ?

$$e^{\int \Theta \Sigma + i \int \Theta \hat{\Sigma}}$$

technically disturbs the standard T-duality (Hori-Vafa) transformation :

~~$$\Theta + \bar{\Theta} \Leftrightarrow -(\Gamma + \bar{\Gamma}) - 2V$$

and/or  $\Theta - \bar{\Theta} \Leftrightarrow -(\Gamma - \bar{\Gamma}) + 2i\hat{V}$

by parity transform along  $x^-$ -direction~~

Duality transformation by irreducible superfields :

Roček, Verlinde (1991); Hori, Vafa (2000)

$$\mathcal{L} = \int d^4\theta \left( -\frac{1}{2g^2} B^2 - 2BV - 2(\Gamma + \bar{\Gamma})B \right)$$

- integrate out  $\Gamma$  :  $\mathcal{L} = -\frac{1}{g^2} \int d^4\theta |\Theta|^2 - \int d^2\tilde{\theta} \Theta \Sigma + \text{h.c.}$
- integrate out  $B$  :  $\mathcal{L} = +\frac{g^2}{2} \int d^4\theta (\Gamma + \bar{\Gamma} + 2V)^2$

How to generalize this method to  $\Theta(\Sigma + i\hat{\Sigma})$  ?

Role of  $X, W, Y$  :

$$\text{Im}(X + \overline{W}) = \text{Im } \Theta \quad \xleftrightarrow{\text{dual}} \quad \text{Im}(X - \overline{W})$$

$$\text{Re}(X + \overline{W}) = \text{Re } \Theta \quad \xleftrightarrow{\text{dual}} \quad \text{Re}(X - \overline{W})$$

$Y$  : link LHS with RHS

Different from the duality transformation by irreducible superfields,  
we can dualize real/imaginary part of  $\Theta$  without any disturbance !

$L = X + \overline{W} + Y$  carries both original and dual fields  $\rightarrow$  “doubled” GLSM

**B.** **Another GLSM**  
**for 5-branes of codim-2**

“Semi-doubled” GLSM for 5-branes of codim-2  $\mathcal{L} = \mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{CHM}} + \mathcal{L}_{\text{NHM}}$  :

$$\mathcal{L}_{\text{gauge}} = \int d^4\theta \left\{ \frac{1}{e^2} (-|\Sigma|^2 + |\Phi|^2) + \frac{1}{\widehat{e}^2} (-|\widehat{\Sigma}|^2 + |\widehat{\Phi}|^2) \right\}$$

$$\begin{aligned} \mathcal{L}_{\text{CHM}} = & \int d^4\theta \left\{ |Q|^2 e^{+2V} + |\widetilde{Q}|^2 e^{-2V} \right\} - \int d^2\theta \widetilde{Q} \Phi Q + \text{h.c.} \\ & + \int d^4\theta \left\{ |P|^2 e^{+2\widehat{V}} + |\widetilde{P}|^2 e^{-2\widehat{V}} \right\} - \int d^2\theta \widetilde{P} \widehat{\Phi} P + \text{h.c.} \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{\text{NHM}} = & \int d^4\theta \left\{ \frac{1}{g^2} |\Psi|^2 + g^2 |L + 2(V + i\widehat{V})|^2 \right\} \\ & + \left\{ \int d^2\theta (s - \Psi) \Phi + \int d^2\widetilde{\theta} t \Sigma + \text{h.c.} \right\} + \epsilon^{mn} \partial_m (\vartheta A_n) \\ & + \left\{ \int d^2\theta (\widehat{s} - \Psi) \widehat{\Phi} + \int d^2\widetilde{\theta} \widehat{t} \widehat{\Sigma} + \text{h.c.} \right\} + \epsilon^{mn} \partial_m (\varphi \widehat{A}_n) \end{aligned}$$

skip the detail...

## Semi-doubled NLSM

$$\begin{aligned} \mathcal{L} = & -\frac{H}{2} \{ (\partial_m x)^2 + (\partial_m y)^2 + (\partial_m \varphi)^2 + (\partial_m \vartheta)^2 \} + \Omega \epsilon^{mn} (\partial_m \varphi) (\partial_n \vartheta) \\ & + \epsilon^{mn} (\partial_m \varphi) (\partial_n \tilde{\gamma}) + \epsilon^{mn} (\partial_m \vartheta) (\partial_n \gamma) \end{aligned}$$

provides string sigma model for defect NS5, KK-vortex, and  $5_2^2$ .

$$\text{dNS5: } \mathcal{L} = -\frac{H}{2} \{ (\partial_m x)^2 + (\partial_m y)^2 + (\partial_m \varphi)^2 + (\partial_m \vartheta)^2 \} + \Omega \epsilon^{mn} (\partial_m \varphi) (\partial_n \vartheta)$$

$$\text{KK-v: } \mathcal{L} = -\frac{H}{2} \{ (\partial_m x)^2 + (\partial_m y)^2 + (\partial_m \varphi)^2 \} - \frac{1}{2H} (\partial_m \gamma - \Omega \partial_m \varphi)^2$$

$$5_2^2: \mathcal{L} = -\frac{H}{2} \{ (\partial_m x)^2 + (\partial_m y)^2 \} - \frac{H}{K} \{ (\partial_m \tilde{\gamma})^2 + (\partial_m \gamma)^2 \} - \frac{\Omega}{K} \epsilon^{mn} (\partial_m \tilde{\gamma}) (\partial_n \gamma)$$

$$H = \log \frac{\Lambda}{r}, \quad \Omega = \arctan \left( \frac{y}{x} \right), \quad r = \sqrt{x^2 + y^2}, \quad K = H^2 + \Omega^2$$

$$\text{instanton corrections: } H \rightarrow H + \sum_{(n,m) \neq (0,0)} e^{in\vartheta} e^{im\varphi} K_0(r\sqrt{n^2 + m^2})$$

modified Bessel of 2nd kind



