

Efimov physics of hadrons



Tetsuo Hyodo

Yukawa Institute for Theoretical Physics, Kyoto Univ.

2018, Jan 29th 1

Contents



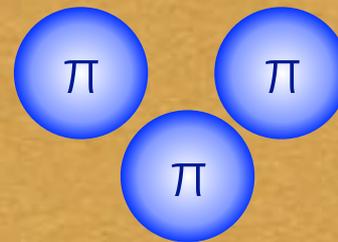
Introduction

- Universal physics in few-body systems
- Tuning hadron interactions by quark mass



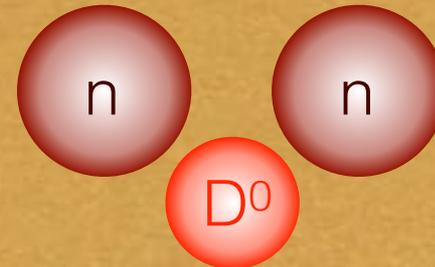
Three-pion systems

[T. Hyodo, T. Hatsuda, Y. Nishida, Phys. Rev. C89, 032201\(R\) \(2014\)](#)



Two neutrons with a flavored meson (K^-/D^0)

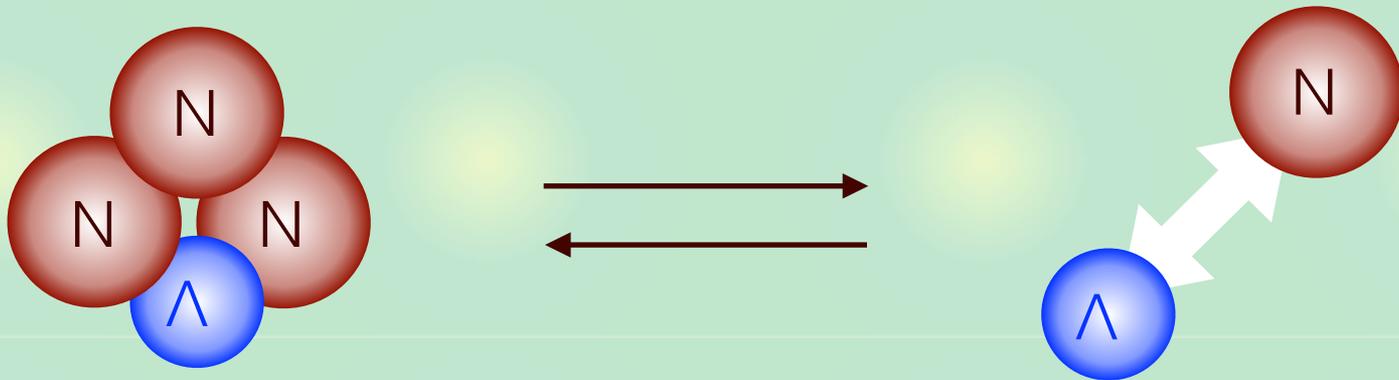
[U. Raha, Y. Kamiya, S.-I. Ando, T. Hyodo, arXiv:1708.03369 \[nucl-th\]](#)



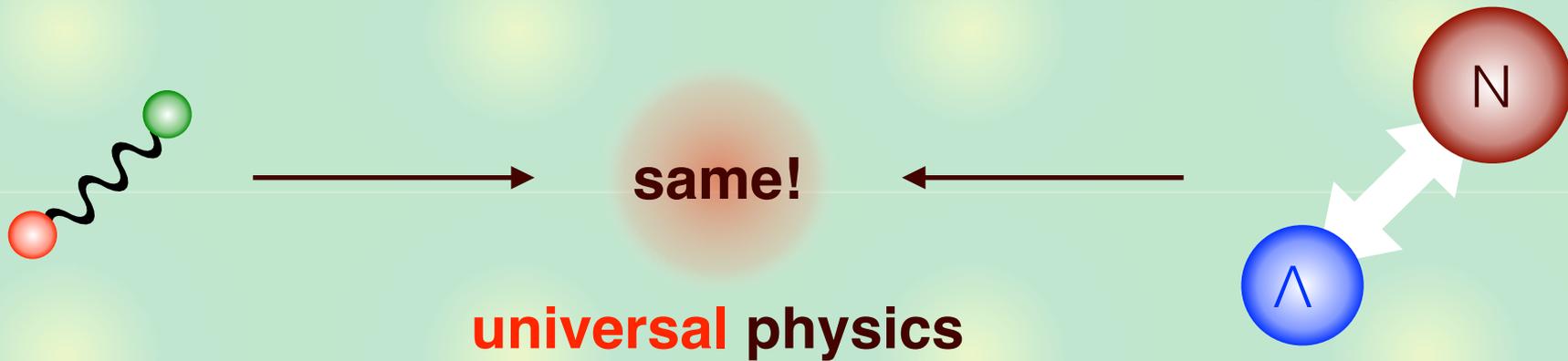
Summary

Study of few-body systems

Properties of few-body systems \longleftrightarrow two-body interaction
 - c.f. hypernuclei



In some cases, different interactions give the same physics.



E. Braaten, H.-W. Hammer, *Phys. Rept.* 428, 259 (2006);

P. Naidon, S. Endo, *Rept. Prog. Phys.* 80, 056001 (2017)

Two-body universal physics

Universal two-body physics: **unitary limit**

1) s-wave short range interaction

2) scattering length : $|a| \gg r_s$: interaction range

- system is scale invariant

- a shallow bound state exists if $a > 0$

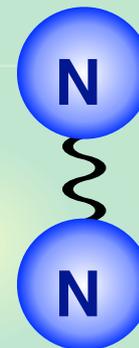
$$B_2 = \frac{1}{ma^2} \left[1 + \mathcal{O}\left(\frac{r_s}{a}\right) \right]$$

vdW

Examples: nucleons and ^4He atoms

	N [MeV]	^4He [mK]
B_2	2.22	1.31
$1/ma^2$	1.41	1.12

strong



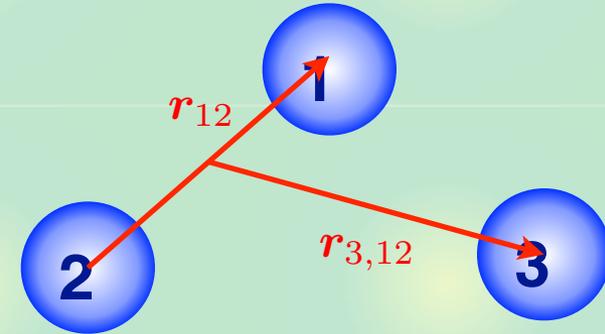
^4He

Three-body universal physics

Three-body system in hyperspherical coordinates

$$(\mathbf{r}_{12}, \mathbf{r}_{3,12}) \leftrightarrow (\underline{R}, \alpha_3, \hat{\mathbf{r}}_{12}, \hat{\mathbf{r}}_{3,12})$$

hyperradius hyperangular variables Ω
(dimensionless)



If $|a| \rightarrow \infty$, system is scale invariant.

$$V(R, \Omega) \propto \frac{1}{R^2}$$

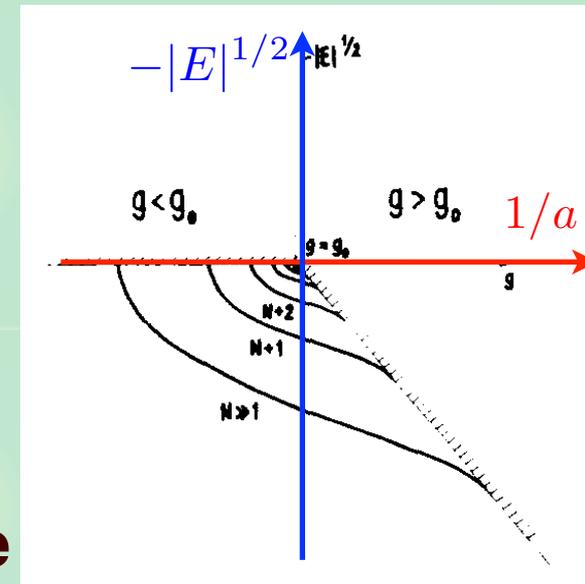
Efimov effect: attractive $1/R^2$

V. Efimov, Phys. Lett. B 33, 563 (1970)

$$B_3^n / B_3^{n+1} \approx 22.7^2$$

- infinitely many bound states
- discrete scale invariance: RG limit cycle

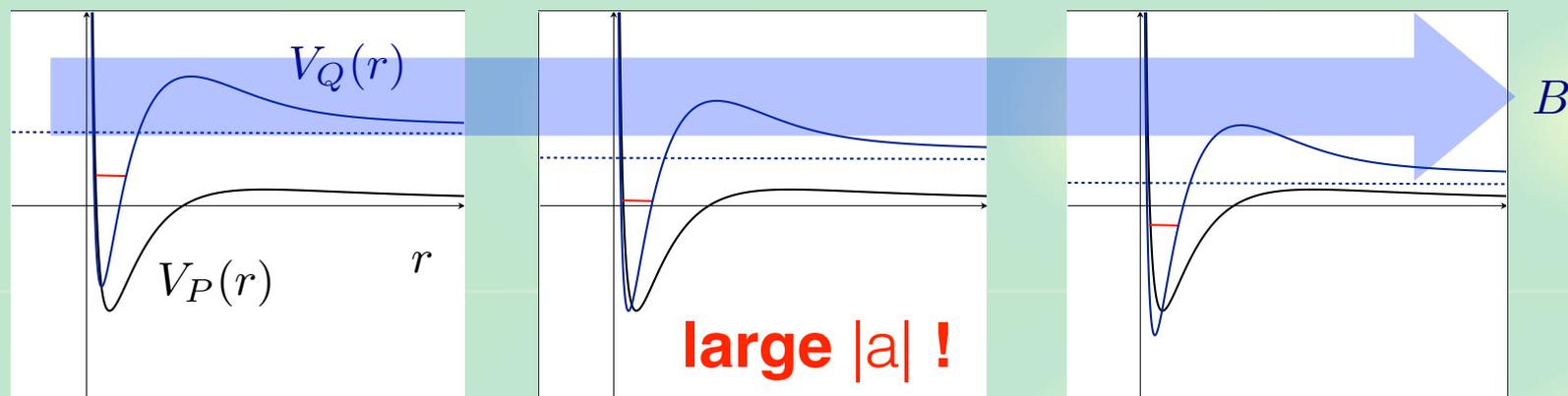
P.F. Bedaque, H.-W. Hammer, U. van Kolck, Phys. Rev. Lett. 82, 463 (1999)



Tuning two-body interactions

Large $|a|$ is achieved by **tuning** two-body interaction

- Atomic physics: Feshbach resonance



In hadron physics, interactions are basically fixed.

- quark mass term in QCD: external scalar field

$$\mathcal{L}_{\text{QCD}} = \mathcal{L}_{\text{QCD}}^{(0)} - m\bar{q}q$$

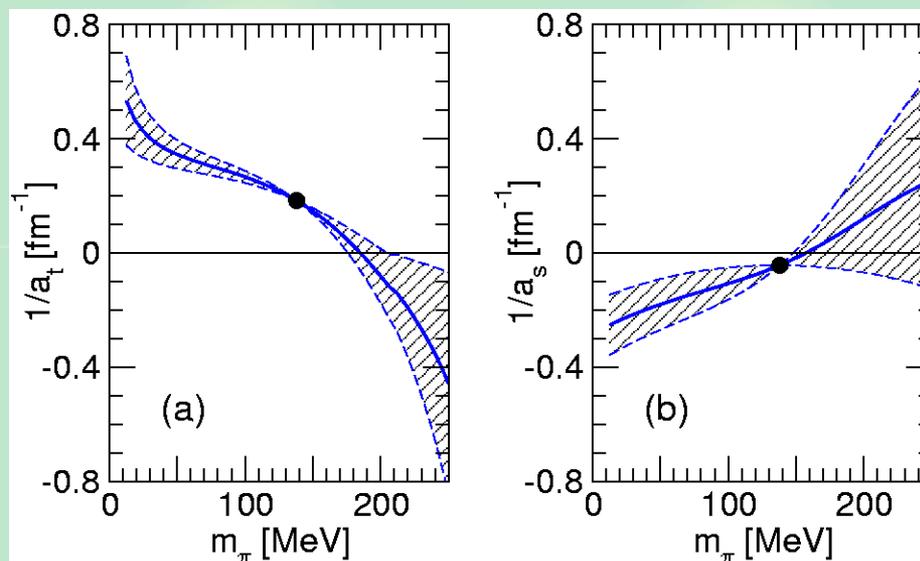
- variation of **quark mass** \rightarrow tuning hadron interaction?

Tuning two-hadron interactions

Nuclear force @ unphysical quark masses

- Nuclear forces can reach unitary limit by $m_{ud}(m_\pi) \uparrow$

E. Braaten, H.-W. Hammer, *Phys. Rev. Lett.* **91**, 102002 (2003)



- predictable by chiral EFT for small m_{ud}
- calculable by **lattice QCD** for large m_{ud}

c.f.) N. Barnea, L. Contessi, D. Gazit, F. Pederiva, U. van Kolck,
Phys. Rev. Lett. **114**, 052501 (2015)

Two-pion interaction

$\pi\pi$ scattering length \leftarrow chiral low energy theorem

S. Weinberg, *Phys. Rev. Lett.* **17**, 616-621 (1966)

$$a^{I=0} \propto -\frac{7}{4} \frac{m_\pi}{f_\pi^2}, \quad a^{I=2} \propto \frac{1}{2} \frac{m_\pi}{f_\pi^2}$$

- $m_\pi \sim (m_{ud})^{1/2} \sim$ **explicit breaking of chiral symmetry**

Physical $\pi\pi$ ($I=0$) is unbound but has a resonance “ σ ”
(resonance: unstable eigenstate above threshold)

$f_0(500)$ or σ
 was $f_0(600)$

$I^G(J^{PC}) = 0^+(0^{++})$

A REVIEW GOES HERE – Check our WWW List of Reviews

$f_0(500)$ T-MATRIX POLE \sqrt{s}

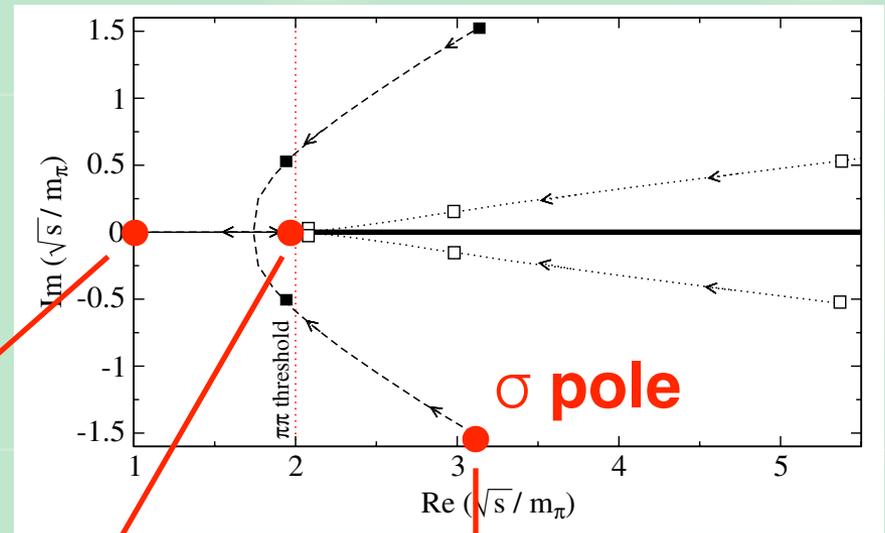
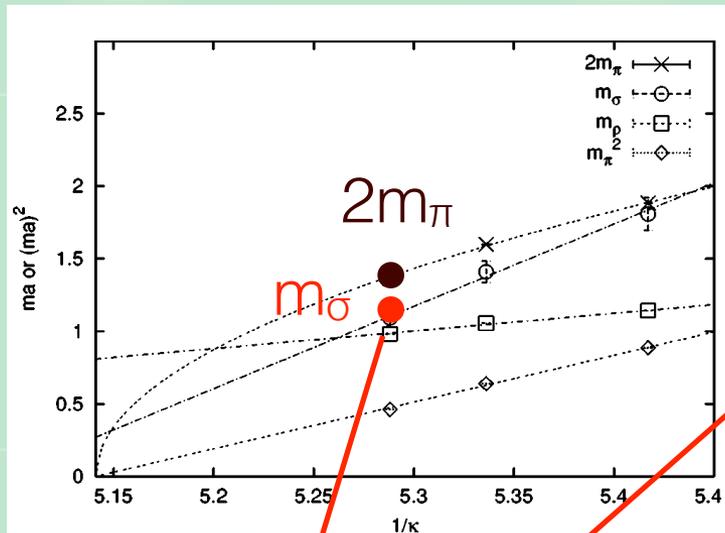
Note that $\Gamma \approx 2 \operatorname{Im}(\sqrt{s_{\text{pole}}})$.

VALUE (MeV)	DOCUMENT ID	TECN	COMMENT
(400–550)–i(200–350) OUR ESTIMATE			

Low energy theorem is valid only for small m_{ud}
How about large m_{ud} ? \rightarrow lattice QCD

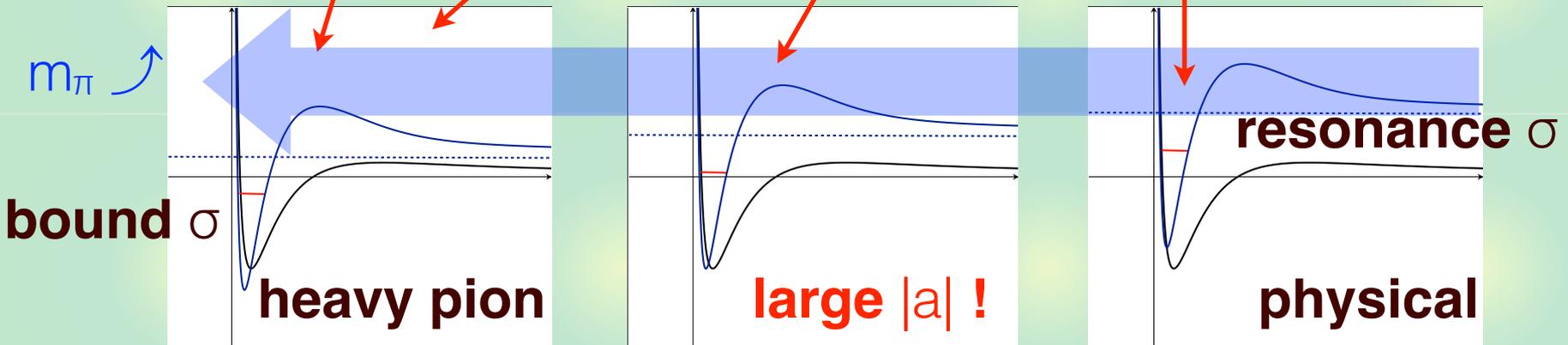
Increasing ud quark mass

Lattice QCD/chiral EFT can tune the $\pi\pi\pi$ interaction



T. Kunihiro *et al.* (SCALAR Collaboration), *Rev. Rev. D* **70**, 034504 (2004)

C. Hanhart, J.R. Pelaez, G. Rios, *Phys. Rev. Lett.* **100**, 152001 (2008)



$\pi\pi\pi$ scattering can reach unitary limit by increasing m_{ud}

Universal physics of three pions

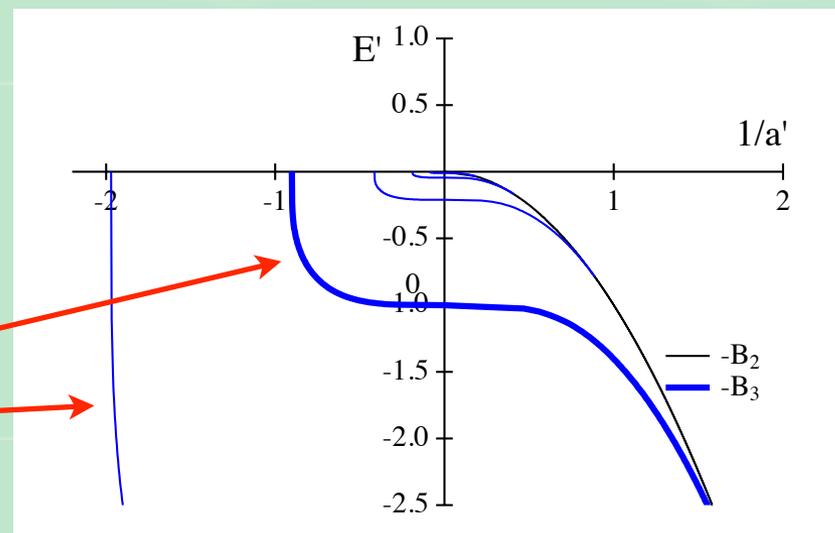
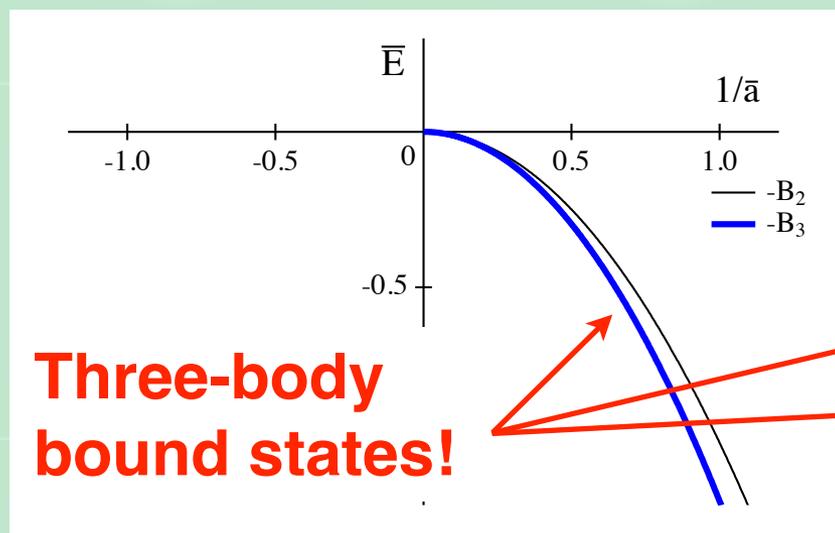
Universal physics of $\pi\pi\pi$ with large $\pi\pi(l=0)$ scattering length

T. Hyodo, T. Hatsuda, Y. Nishida, Phys. Rev. C89, 032201(R) (2014)

- Let $l=0$ $\pi\pi$ scattering length large by **changing** m_q

isospin sym. ($\pi^0\pi^0\pi^0$ - $\pi^0\pi^-\pi^+$)

isospin breaking ($\pi^0\pi^0\pi^0$)



- Universal physics of pions @ unphysical m_{ud}
- Coupled-channel effect reduces the attraction.

Implication for real world

Universality -> a $\pi\pi\pi$ bound state @ heavy m_{ud}

- Heavy m_{ud} is continuously connected to physical point.
- Existence of a pole (eigenstate) is stable against the continuous change of parameters (such as m_{ud}).

Y. Kamiya, T. Hyodo, arXiv:1711.04558 [hep-ph]

$$\frac{1}{2\pi} \oint_C dz \frac{d \arg \mathcal{F}(z)}{dz} = (\# \text{ of zeros}) - (\# \text{ of poles})$$

-> $\pi\pi\pi$ state may exist as a resonance at physical point.

Possible candidates $\sim | = 1, J=0$ state : excited state of π ?

$\pi(1800)$	$I^G(J^P)$	$\pi(1300)$	$I^G(J^{PC}) = 1^-(0^-+)$
See also minireview under non- $q\bar{q}$ candid: Physics G33 1 (2006).			
$\pi(1800)$ MASS		$\pi(1300)$ MASS	
VALUE (MeV)	EVTS	DOCUMENT ID	TECN COMMENT
1300 ± 100	OUR ESTIMATE		
VALUE (MeV)	EVTS	DOCUMENT ID	TECN COMMENT

Two neutrons and one flavored meson

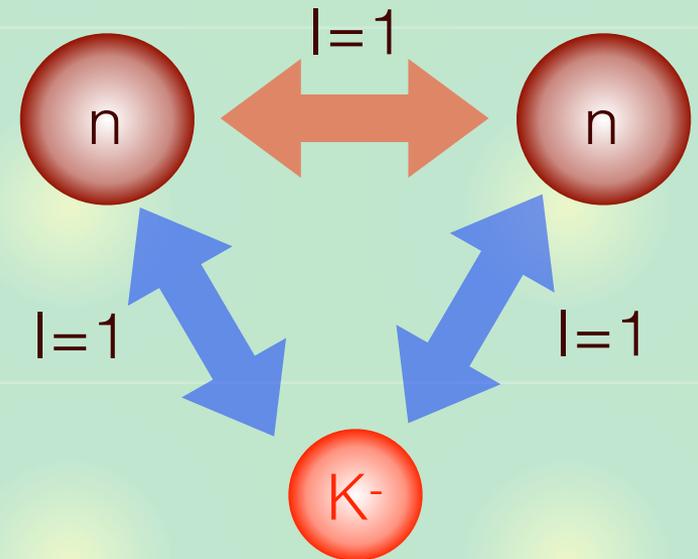
Flavored meson: K^- ($s\bar{u}$, strangeness), D^0 ($c\bar{u}$, charm)

K -nn/ D^0 nn **system with** $J=0$, $l=3/2$, $l_3=-3/2$

- **different from** $J=0$, $l=1/2$ (**so-called** K -pp- \bar{K}^0 np)
- **all interactions: isospin** $l=1$ (**no** $\Lambda(1405)$)

Desirable features for Efimov effect

- **no coupled channels**
- **no Coulomb interaction**
- $|a_{nn}| \sim 20 \text{ fm} \gg r_s \sim O(1) \text{ fm}$



Two-body meson-neutron scattering length?

Meson-neutron interaction

K-n system: $\bar{K}N$ scattering data

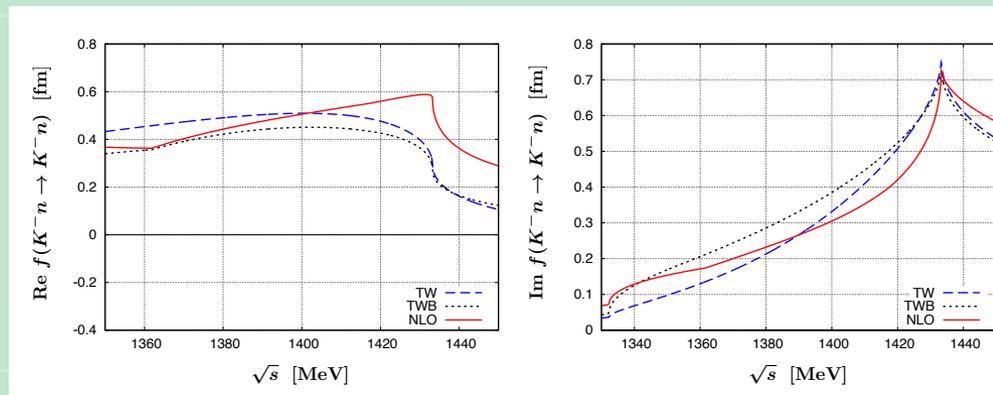
Y. Ikeda, T. Hyodo, W. Weise, Phys. Lett. B706, 63 (2011); Nucl. Phys. A881, 98 (2012)

- data fitted as $\chi^2/\text{d.o.f.} \sim 1$

$$a_{0,K^-n} = -0.57_{-0.04}^{+0.21} - i0.72_{-0.26}^{+0.41} \text{ fm}$$

~~REMOVED~~
decay

→ Strangeness sector is unbound



D^0n system: identify $\Sigma_c(2800)$ as a $J^P=1/2^-$ state

$\Sigma_c(2800)$

$I(J^P) = 1(?^-)$ Status: ***

Seen in the $\Lambda_c^+ \pi^+$, $\Lambda_c^+ \pi^0$, and $\Lambda_c^+ \pi^-$ mass spectra.

$\Sigma_c(2800)$ MASSES

- D^0n threshold ~ 2804 MeV

→ Charm sector has a shallow **quasi-bound state**

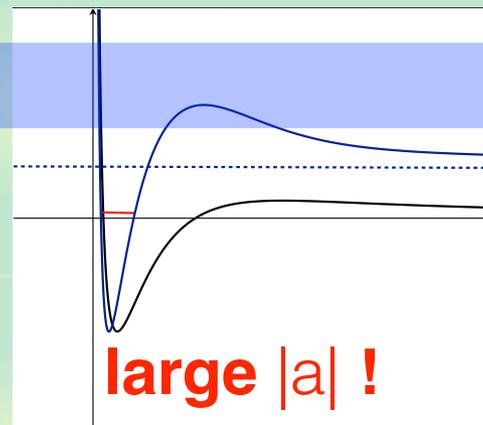
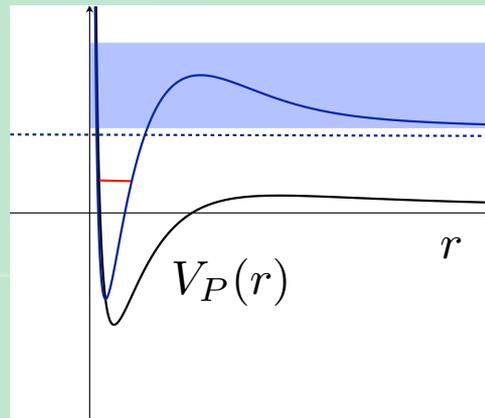
Idealization

Zero coupling limit (ZCL)

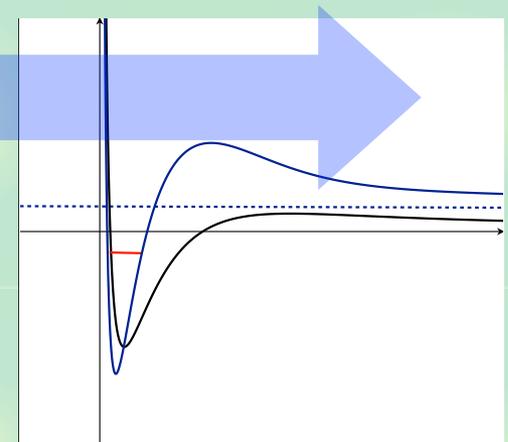
- coupling to **decay channels** are switched off
- $a_{K-n} < 0$ (**attractive**), $a_{D0n} > 0$ (**repulsive, with bound state**)

Varying $m_s \rightarrow m_c$

strange, **unbound**



charm, **bound**



Unitary limit: tuning $m_{s/c}$ in ZCL

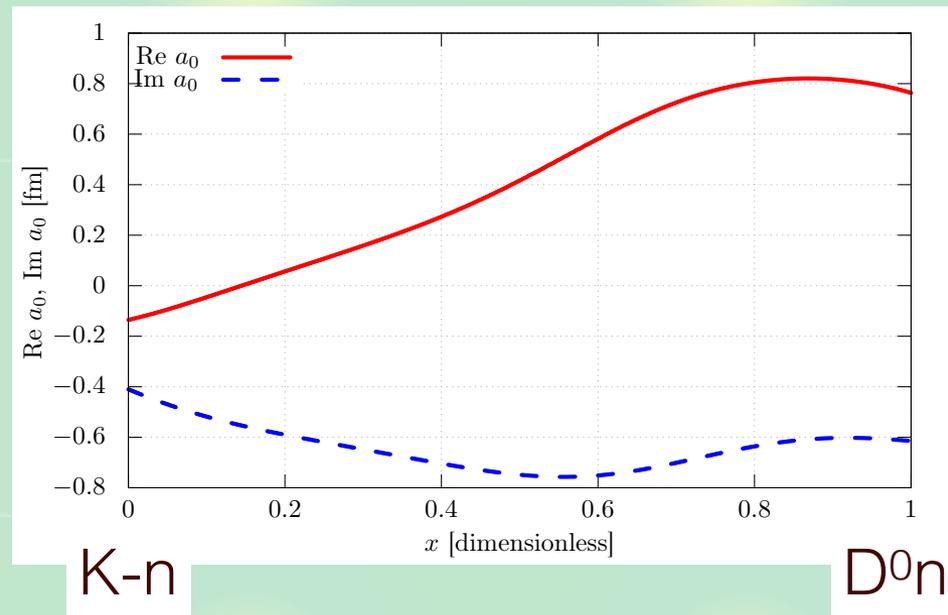
Model extrapolation

Contact interaction model with extrapolation parameter x

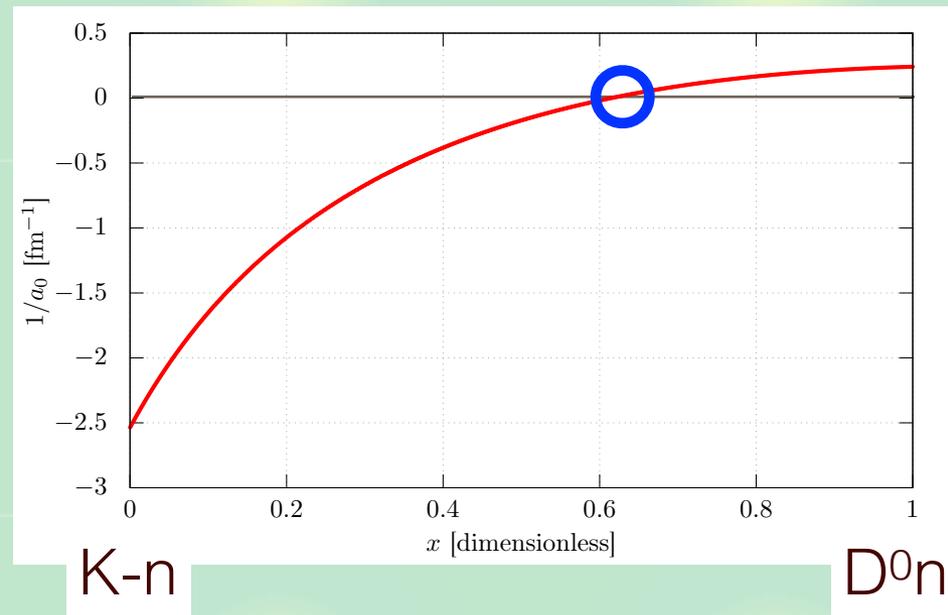
U. Raha, Y. Kamiya, S.-I. Ando, T. Hyodo, arXiv:1708.03369 [nucl-th]

- $x=0$: K^-n , $x=1$: D^0n

with decay channels



ZCL (without decay)



- In ZCL, **unitary limit** at $m_K = 1337$ MeV ($x \sim 0.6$)

- With decay channel, remnant is not very clear.

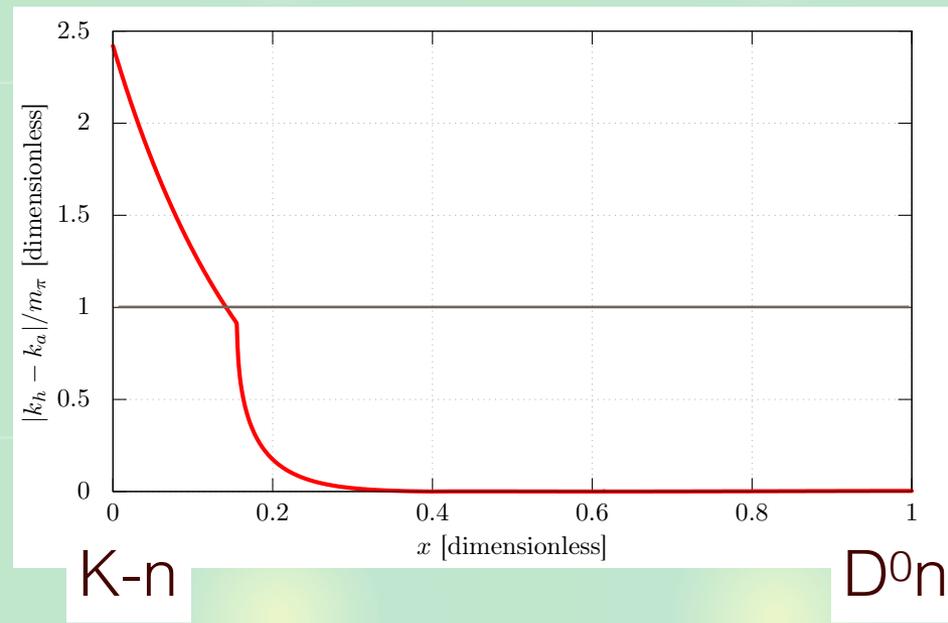
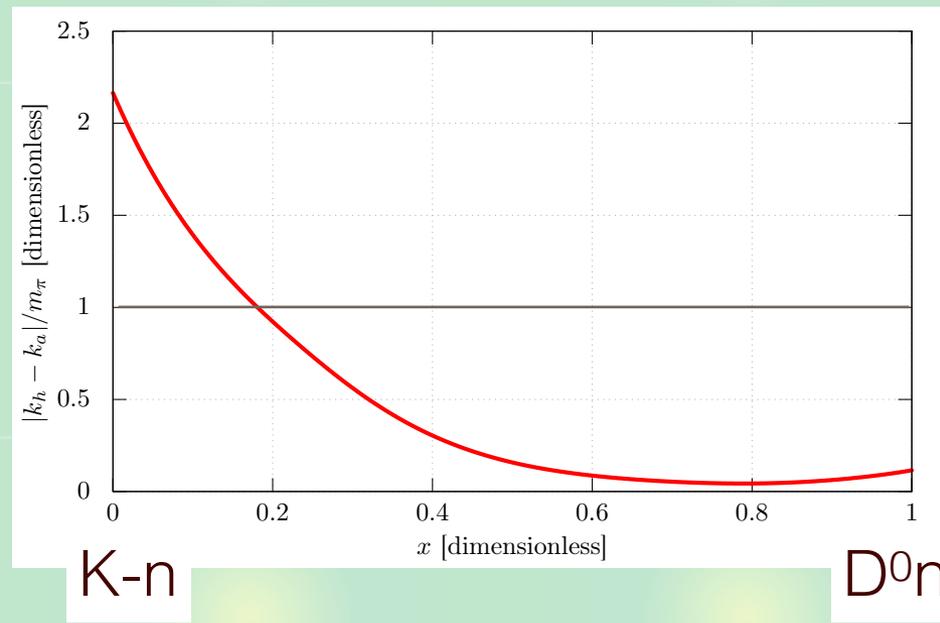
Two-body universality

Check of universality: $|k_h - k_a|/m_\pi$

- full eigenmomentum: $k_h = \sqrt{2\mu E_h}$
- universality prediction: $k_a = \frac{i}{a} \left[1 + \mathcal{O}\left(\frac{r_s}{a}\right) \right], \quad r_s \sim 1/m_\pi$

with decay channels

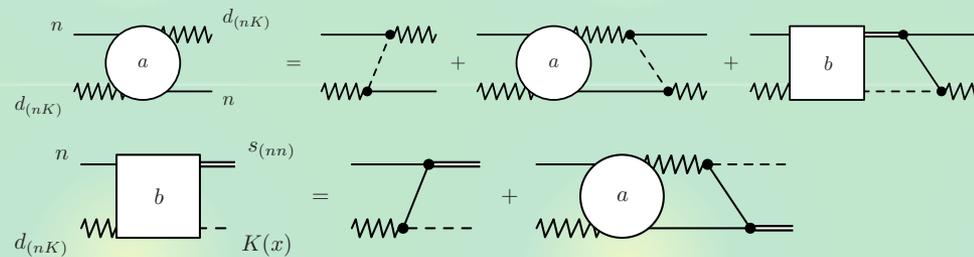
ZCL (without decay)



Universality governs the system near D^0n even with decay

Three-body system

Three-body equations for K- nn /D 0nn system



Asymptotic behavior of K- nn /D 0nn system at unitary limit:

- $a_{S(nn)} \longrightarrow \infty$ **and** $a_{d(nK)} \longrightarrow \infty$

$$1 = C_1 \frac{2\pi \sin[s \arcsin(a/2)]}{s \cos(\pi s/2)} + C_2 \frac{4\pi^2 \sin^2[s \operatorname{arccot}(\sqrt{4b-1})]}{s^2 \cos^2(\pi s/2)} \Rightarrow s_0 = 1.01156$$

- $a_{S(nn)}$ **fixed** and $a_{d(nK)} \longrightarrow \infty$

$$1 = C_1 \frac{2\pi \sin[s \arcsin(a/2)]}{s \cos(\pi s/2)} \Rightarrow s_0 = 0.327675$$

- **RG limit cycle in the asymptotic expressions: Efimov effect**

Implication for real world:

\longrightarrow D 0nn state may exist **as a resonance** at physical point.

Summary

 Hadron-hadron interactions can be tuned by changing **quark masses**.

 **Increase** m_{ud} : unitary $\pi\pi$ ($l=0$) scattering
—> Efimov effect in $\pi^0\pi^0\pi^0$

[T. Hyodo, T. Hatsuda, Y. Nishida, Phys. Rev. C89, 032201\(R\) \(2014\)](#)

 **Decrease** m_c (+ZCL) : unitary D^0n scattering
—> Efimov effect in D^0nn

[U. Raha, Y. Kamiya, S.-I. Ando, T. Hyodo, arXiv:1708.03369 \[nucl-th\]](#)

 Remnant of Efimov physics may be observed as a three-body resonance in nature.