# ブラックホール (or 中性子星) 降着流と噴出流 の数値シミュレーション 〜最近の成果と今後の発展〜

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## Accretion Power

Black hole (BH) accretion system is one of the most powerful energy production mechanisms in the Universe. The luminous compact objects, like active galactic nuclei, X-ray binaries, Gamma-ray bursts, are believed to be powered by the BH accretion system.

X-ray binary (Cyg X-1) quasar (3C273)





## Three Accretion Modes

Different spectral states imply the existence of different accretion modes



**Mass accretion rate (disk luminosity)**





# Three Accretion Modes



### Importance of Radiation and Magnetic Fields



Radiation-HD/MHD Simulations are necessary for high accretion rate.

#### **Magnet Fields;**

- Angular momentum is transported by MRI, leading to the mass accretion onto BHs.
- $\blacksquare$  Magnetic dissipation heats the gas (heating)

Magnetic force drives outflows

### **Radiation Fields;**

- Disk loses the energy by emitting photons (cooling).
- Radiation pressure determines the thickness of the disk.
- Radiation force drives outflows

## Development of Simulations of Super-Edd. disks

### **1D approach** Slim disk model has been established (Abramowicz et al. 1988) **Radiation-HD sim. Radiation-MHD sim. General Relativistic Radiation-MHD Sim.** Quasi steady inflow-outflow structure has been revealed. **Multi-dimensional Simulations** General relativistic effects (e.g., BZ effect, LT precession) has been studied 1988~  $2005~$  $2014$ ~ (**Ohsuga et al. 2005, 2009, Ohsuga & Mineshige 2011**, Jiang et al. 2014, 2019) (Sadowski et al. 2014, 1016, **Takahashi, Ohsuga et al. 2016, Utsumi, Ohsuga et al. 2022**, **Asahina et al. 2022**, Brandon 2023, **Asahina & Ohsuga submitted**).

## Numerical Methods for Radiation Field



### Radiation-MHD simulations Super-Edd. Flows

![](_page_8_Picture_1.jpeg)

Ohsuga et al. 2009; Ohsuga & Mineshige 2011

### **Setup**

- BH mass: 10Msun
- $\blacksquare$  Initial condition: equilibrium torus with embedded poloidal magnetic field (plasma-beta=100)

### **Quasi-steady structure**

- The super-Eddington disks (Mdot  $\sim$  a few  $100$ L $Edd/C<sup>2</sup>$ , Ldisk >> L $Edd$ )
- Radiatively-driven outflows

see also Ohsuga et al. 2005, Ohsuga 2007, Jiang et al. 2014, 2019

## Why is super-Eddington accretion feasible?

![](_page_9_Picture_1.jpeg)

### **Radiatively driven outflows:**

Strong radiation pressure supports the thick disk and generates the outflows above the disk.

#### Accretion:

Photons mainly escape through the less-dense region above the disk. The radiation pressure cannot prevent the accreting motion within the disk.

#### Why is super-Eddington accretion feasible? No. 2, 2007 SUPERCRITICAL DISK ACCRETION 1287

![](_page_10_Figure_1.jpeg)

# Apparent Luminosity

![](_page_11_Picture_1.jpeg)

- **The radiative flux is mildly collimated since** the disk is optically and geometrically thick.
- Thus, observed luminosity is very sensitive to the observer's viewing angle.
- **The apparent luminosity becomes highly** super-Eddington for the face-on observers.

 $ex: 22$ LEdd for  $\leq 20^{\circ}$ when Mdot $\sim$ 100L $_{\text{Edd}}/c^2$  & Ldisk $\sim$ 3L $_{\text{Edd}}$ .

Ohsuga, Mineshige <sup>2011</sup> **Large luminosity of ULXs (>1039-40erg/s) can be explained for the face-on case.**

![](_page_12_Figure_0.jpeg)

### Comparison with ULXs Kitaki et al. 2017

![](_page_13_Figure_1.jpeg)

High-energy X-ray photons are generated within the funnel region. Photons undergo down-scattering above the disk, and in some cases, they are absorbed.

## Radiatively-driven Jets

![](_page_14_Figure_1.jpeg)

## Clumpy Outflows

![](_page_15_Figure_1.jpeg)

Takeuchi, Ohsuga, Mineshige 2013

Clumpy outflows: Radiative Winds fragment into many gas clouds

### RT instability

![](_page_16_Figure_1.jpeg)

![](_page_17_Picture_0.jpeg)

## Observations of Clumpy outflows

Some ULXs exhibit the time variations of X-ray luminosity, implying the launching of clumpy outflows.

Launching of clumpy winds is also reported by observations of NLS1s or V404 Cyg.

![](_page_18_Figure_3.jpeg)

Middleton+11 Jin+17 see also Motta+17

![](_page_18_Figure_5.jpeg)

# Comparison with ULXs

### Absorption lines Outflow velocity of  $\sim 0.1$ -0.2c agrees with the observations of blueshifted lines.

![](_page_19_Figure_2.jpeg)

### Time variation Timescale of the luminosity variation (100Rs/0.3Vkep) is

$$
\sim 2.5 \left(\frac{M_{\rm BH}}{10\,M_{\odot}}\right) \left(\frac{\ell^{\theta}_{\rm cl}}{10^2\,r_{\rm S}}\right) \left(\frac{r}{10^3\,r_{\rm S}}\right) {\rm s}
$$

Our result is consistent with the observations of ULXs (Middleton+11) and V404 Cyg (Motta+17) in the case of MBH~10-100Msun.

# Overall structure of the super-Edd. flows

Schematic picture of the overall structure

![](_page_20_Figure_2.jpeg)

Super-Eddington flows consist of three components;  $\square$  radiation pressure-dominated disk  $\square$  radiatively-driven high-velocity outflow around the rotation axis (jet)  $\Box$  radiatively-driven clumpy wind (&

failed wind).

## General Relativistic Radiation-MHD sim.

Utsumi, Ohsuga et al. 2022

![](_page_21_Figure_2.jpeg)

### **Setup**

- BH mass: 10Msun
- Initial condition: equilibrium torus with embedded poloidal magnetic field (plasma-beta=100)
- $\blacksquare$  Spin parameter: -0.9, -0.7, -0.5, -0.3, 0, 0.3, 0.5, 0.7, 0.9

### **Quasi-steady structure**

- In all models, the super-Eddington disks (Mdot  $\sim$  a few 100L $_{\text{Edd}}/c^2$ ) and strong outflows are formed.
- \* Magnetic field is not so strong (SANE)

# Energy Conversion Efficiency

![](_page_22_Figure_1.jpeg)

For the case of  $a^* \sim 0$ , energy is mainly released by the radiation. When  $|a^*|$  is large, the energy released by the Poynting flux (Magnetic Luminosity) exceeds the Radiation Luminosity.

Radiation luminosity accounts for 80% when  $a^* \sim 0$ . But the magnetic luminosity is three times larger than the radiation luminosity for the case of  $a^*$  > 0.5.

See also Sadowski et al 2014, 2019, Bandon 2023

# Enhancement of Poynting Flux

![](_page_23_Figure_1.jpeg)

The Poynting flux around the rotation axis is stronger for larger |a\*|. This is probably caused by Blandford-Znajek (BZ) effect.

# Enhancement of Poynting Flux

![](_page_24_Figure_1.jpeg)

実線:磁気フラックス  $-M_t^r = -(b^2 u_t u^r - b_t b^r)$ 

破線 : BZフラックス  $F_{\rm BZ}|_{r_{\rm H}} = 2(B^r)^2 \omega r_{\rm H} (\Omega_{\rm H} - \omega) \sin^2 \theta$ 

# Kinetic luminosity vs X-ray luminosity

![](_page_25_Figure_1.jpeg)

\*Isotropic X-rau Luminosity: Radiation luminosity observed by face-on observer.

In our results, the ratio of the kinetic luminosity to isotropic X-ray luminosity tends to increase with  $|a^*|$ .

Thus, rapidly (slowly) rotating black hole probably exist in IC342 X-1 (Holmberg II

## Lense-Thirring Precession of Super-Edd. disk

![](_page_26_Picture_1.jpeg)

Asahina & Ohsuga accepted yesterday

### **Setup**

#### BH mass: 10Msun

**Initial condition: equilibrium torus with** embedded poloidal magnetic field (plasma-beta=100) tilted **30 degree**.

■ Spin parameter: 0.9

### **Inflow-outflow structure**

- The super-Eddington disk, which is tilted and twisted, forms.
- Strong outflows are also formed.
- p Accretion rate: several **100 LEdd/c<sup>2</sup>** p Radiation Luminosity: **several LEdd** p Kinetic Luminosity: **several LEdd**

# Tilted and twisted super-Edd. disk

![](_page_27_Figure_1.jpeg)

- The tilt angle of the outer region is  $\sim$ 30°, which is determined by the initial setting of the torus.
- $\blacksquare$  The disk is gradually tilted as it approaches the black hole, except for the region of  $r < 5r<sub>g</sub>$ .
- **The precession angle increases as it** approaches the BH.

![](_page_27_Figure_5.jpeg)

# Precession of disk, outflow, radiation

![](_page_28_Figure_1.jpeg)

- $\blacksquare$  The super-Eddington disk exhibits the precession
- $\blacksquare$  The gas and radiation is mainly ejected around the rotation axis of the disk ( $\sim$ 30°), rather than around the spin axis of the BH  $(0^{\circ})$ .
	- $\blacksquare$  The direction of outflow and radiation also changes according to the precession motion of the disk.

# Comparison with observations

[1] Quasi periodic oscillations of ULXs: The typical timescale of the precession is  $\sim$  several The typical timescale of the precession is  $\sim$  several  $\frac{1}{8}$ <br>sec for the case of stellar mass BH and disk size is  $\frac{1}{8}$ few 10rg. This timescale is consistent with the QPOs observed in some ULXs (0.01-1Hz, Atapin 2019).

![](_page_29_Figure_2.jpeg)

[2] Precession of jets in V404 Cygni:

The direction of jet is changing with time in V404 Cygni (a few min  $\sim$  a few hours, Miller-Jones et al. 2019).

Such behavior maybe reproduced if the disk size is a few 100rg.

![](_page_29_Figure_6.jpeg)

## Lense-Thirring Precession of RIAF Cui et al. 2023

![](_page_30_Figure_1.jpeg)

# Super-Edd. Flows around magnetized NSs

Inoue et al. 2023, Inoue et al. in prep.

### **Code**

2D General relativistic Radiation-MHD simulation code

### **Setup**

- NS: 1.4Msun, 10km
- **Magnetic field:** Bdip/(Bdip+Bqua)=1, 0.75, 0.5, 0.25, 0
- Initial condition: equilibrium torus with embedded poloidal magnetic field (plasma $beta=100$

![](_page_31_Picture_8.jpeg)

# ULX Pulsars

![](_page_32_Figure_1.jpeg)

#### **NS + Super-Eddington flow**

If the central objects of ULXs are NSs, super-Eddington is necessary because the mass of NSs is a few Msun.

![](_page_32_Picture_4.jpeg)

Basko & Sunyaev 76; Ohsuga 07; Mushtukov+15, 18; King & Lasota 16; Kawashima et al. 16; Takahashi & Ohsuga 17, 18; Chashkina+17

# Accretion column

The radiation energy escapes from the side face of the accretion columns. The side face is likely the origin of the X-ray pulse.

![](_page_33_Figure_2.jpeg)

![](_page_33_Figure_3.jpeg)

Radiation energy is released from the side surface of the column.

Kawashima et al. 2016

see also Basko & Sunyaev (1976) Meszaros (1998), Mushtukov et al. (2015), Zhang et al. (2023), Abolmasov et al. (2023)

# Dipole-dominated case

![](_page_34_Figure_1.jpeg)

## **ULXP Swift J0243.6+6124** Transient ULXP

 $z$  [km]

 $\mathcal{D}$  $\dot{M}_{\text{in}} \sim 700 \dot{M}_{\text{Edd}}$   $T_{\text{gas}}$  [K] Counts cm<sup>-2</sup> s<sup>-1</sup> (15-50 keV)  $1.5$  $10<sup>8</sup>$ AstroSat 1500  $Obs-2$ Photosphere of the outflows:  $\mathbf{1}$ Size: a few 100km 750 AstroSa Temperature: ~0.5 keV  $0.5$  $Obs 10<sup>7</sup>$  $\bf{0}$  $\mathsf{O}$ 50 100 150 Time (MJD-58029) The observed blackbody 10 radiation can be explained. 90302319006  $-750$  $\cos$  cm<sup>-2</sup> s<sup>-1</sup> keV<sup>-1</sup>) cutoffpl (colu bb (outflow bb (column top)  $10<sup>6</sup>$  $-1500$ 750  $\Omega$  $\overline{O}$ keV (Pt

> Black body emission

Tao et al. 2019, Beri et al. 2021

bb (hot soc

## ULXP Swift J0243.6+6124: magnetic fields

Pulse period:  $\sim$ 9.8s

 $\blacksquare$  Spin-up rate: 6.8x10<sup>-9</sup>ss<sup>-1</sup> (obs 1), 1.8x10<sup>-8</sup>ss<sup>-1</sup> (obs 2), 2.2x10<sup>-8</sup>ss<sup>-1</sup> (obs 3) Black Body emission (T~0.5keV and R~100-500km) is detected in two observations and not in one observation

 $r_{\rm M}$   $r_{\rm sph}$ 

Dipole fields  $\sim$  3x10<sup>11</sup>G - 4x10<sup>12</sup>G Cyclotron resonance scattering feature:  $\sim$ 2x10<sup>13</sup>G Accretion column NS **Outflow** Slim

## Quadrupole-dominated case

![](_page_37_Figure_1.jpeg)

## ULXP Swift J0243.6+6124: magnetic fields

Pulse period:  $\sim$ 9.8s

Spin-up rate:  $6.8x10^{-9}$ ss<sup>-1</sup> (obs 1),  $1.8x10^{-8}$ ss<sup>-1</sup> (obs 2),  $2.2x10^{-8}$ ss<sup>-1</sup> (obs 3) Black Body emission (T~0.5keV and R~100-500km) is detected in two observations and not in one observation

Cyclotron resonance scattering feature:  $\sim$ 2x10<sup>13</sup>G

![](_page_38_Picture_4.jpeg)

Quadrupole fields  $\sim 2x10^{13}$ G  $\leftarrow$  Cyclotron resonance scattering Dipole fields  $< 4 \times 10^{12}$ G

> ↑ If the dipole magnetic field exceeds this value, the magnetospheric radius becomes too large, causing the disappearance of the super-Eddington disk region and preventing the formation of radiatively driven thick winds.

#### Nomura et al. 2016, 2017, 2020, 2021 (see also Proga et al. 00, 04)<br>log density From the standard disk, the disk  $-15$ wind is launched by the radiation  $time = 0.463yr$  $-16.$ force for spectral lines (line-force). 17  $\log p(g \text{ cm}^{-3})$ BH(108Msun)  $\overline{U}$  $-18$ eacceleration Line winds  $-19$ **Nucleus**  $-20$ Line Force: Metal ions get  $-21$ momentum by the line absorption  $10<sup>3</sup>Rg$ (bound-bound transition) of UV. Standard disk  $(L_{disk} \sim 0.5$ LEdd)

## Simulations of Line Winds

野村真理子さん (弘前⼤)

### Simulations of Line Winds

![](_page_40_Figure_1.jpeg)

detected in some Sy galaxies. Line-driven winds  $(\sim 0.1c)$  are one of the plausible model.

![](_page_40_Figure_3.jpeg)

## Comparison with UFOs (PG 1211+143)

![](_page_41_Figure_1.jpeg)

## Future observations

![](_page_42_Figure_1.jpeg)

A more detailed comparison with observations by XRISM provides a more accurate understanding of the disk wind structure.

Good news. Our proposal was accepted last week

Absorption lines from H-like and He-like iron are resolved by XRISM.

乞うご期待!

## Does a line wind not occur !? Higgnbottom et al. 2014

Ionization parameter16.5 16.5 PK04 New calculation Central source, bremstrahlung SED 16.0 16.0 15.5 15.5  $\widehat{\mathbb{E}}_{\substack{\text{N}\\ \text{on}}}$  15.0<br> $\frac{\text{N}}{2}$  14.5 15.0 Low 14.5 ionization 14.0 14.0 region Low ionization region 13.5 13.5  $15.5$  $16.0$ 15.5 14.5  $15.0$ 16.5 14.5 15.0 16.0 16.5  $log x (cm)$  $log x$  (cm)  $\overline{0}$  $\overline{4}$ 6 7 8 9 10  $9<sup>10</sup>$  $-2$   $-1$  $\mathbf{1}$  $\overline{2}$ 3 5  $-2$   $-1$  $\overline{0}$ 2 3  $\overline{\mathbf{4}}$  $\overline{5}$ 6  $7$ 8  $\overline{1}$ log U log U

Conventional calculation methods underestimate the degree of ionization  $\rightarrow$  overestimate the force multiplier.

## Does a line wind not occur !? Higgnbottom et al. 2023

![](_page_44_Figure_2.jpeg)

WDの降着円盤でLine wind のモンテカルロ輻射流体計算を 実施

モンテカルロ輻射輸送計算を実 施したところ, 従来の⽅法より Force multiplierが小さくなり, 質量噴出率は2桁低下.

## Does a line wind not occur !? Higgnbottom et al. 2014

ISSI workshop (2024.6)でのShane Davis氏のスライド

![](_page_45_Picture_3.jpeg)

X線領域の(ちょっとマイナーな) lineを組み込んでみたら, AGNの 場合に限っては質量噴出率が2桁 増加

Line windは強いのか弱いのか? 私も小高氏(阪大)と組んで参 戦予定

### 「富岳」成果創出加速プログラム

![](_page_46_Figure_1.jpeg)

森脇可奈 (東京大)

### Machine Learning of the Edd. tensor

#### 第一原理計算計算の結果を教師データとして機械学習モデル (CWNN) を構築

rr-component of Eddington tensor

![](_page_47_Figure_3.jpeg)

超新星爆発のみなさんは, エディントンテンソルの推定に成功したようです. 我々は・・・

## Machine Learning of the Edd. tensor

Variable Eddington Tensor法 by 朝比奈 の結果を教師データとして機械 学習モデルを構築

密度分布. 左:VET 右:機械学習 インスコントリンク 質量降着率

![](_page_48_Figure_4.jpeg)

![](_page_48_Figure_5.jpeg)

![](_page_48_Figure_6.jpeg)

だいたいOKに⾒えるが・・・

上野君(筑波大M2)

## Machine Learning of the Edd. tensor

Variable Eddington Tensor法 by 朝比奈 の結果を教師データとして機械学習モデルを構築

![](_page_49_Figure_2.jpeg)

![](_page_49_Picture_3.jpeg)

光学的に薄いところで精度が悪い。。。 \*現在はこれより少しマシになってます

#### X-ray polarization of Cyg X-1 Krawczynski et al. 2022

![](_page_50_Figure_1.jpeg)

ただし,竹林晃大君の修論によると, ジェットが散乱体だったとしてもジェットに平行な偏光を説明できる

### Radiation-MHD for super-Edd. TDE

![](_page_51_Figure_1.jpeg)

Super-Edd TDEの輻射流体シミュレーション ■ 潮汐破壊された星が作るガス流を想定し, ある一 点から連続的にガスを流入 ■ 注入率は10Mdot Edd Huang et al. 2024

### Radiation-MHD for super-Edd. TDE

![](_page_52_Figure_1.jpeg)

■ 光度はEddington以下 (衝突で光る)

## Quasi-periodic eruption (QPE)

![](_page_53_Figure_1.jpeg)

X線の光度変化

G. Miniutti et al. 2019, Nature 573, 381

![](_page_53_Figure_4.jpeg)

# おわり