

Collapses of rotating supermassive stars and associating transient

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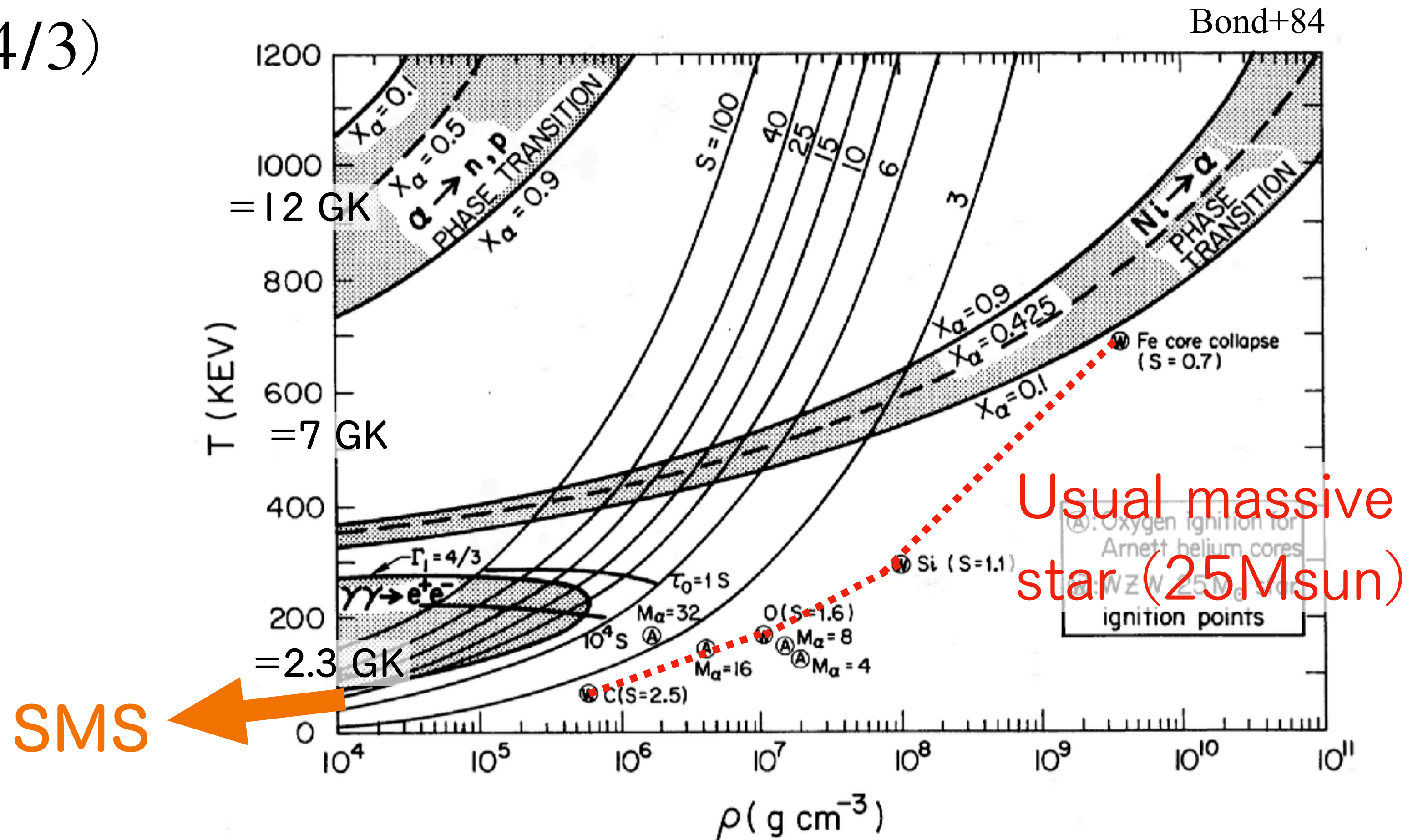


Outline

- Collapse and explosion of supermassive stars ($\sim 10^5 M_{\text{sun}}$)
 - Introduction
 - Numerical setup
 - Results
- Collapse of usual massive star ($\sim 10 M_{\text{sun}}$)

Supermassive star

- Hypothetical very massive ($\gtrsim 10^4 M_\odot$) star
- High- s , P_{rad} -dominant ($\Gamma \approx 4/3$)
- Dies likely by GR instability



Supermassive star

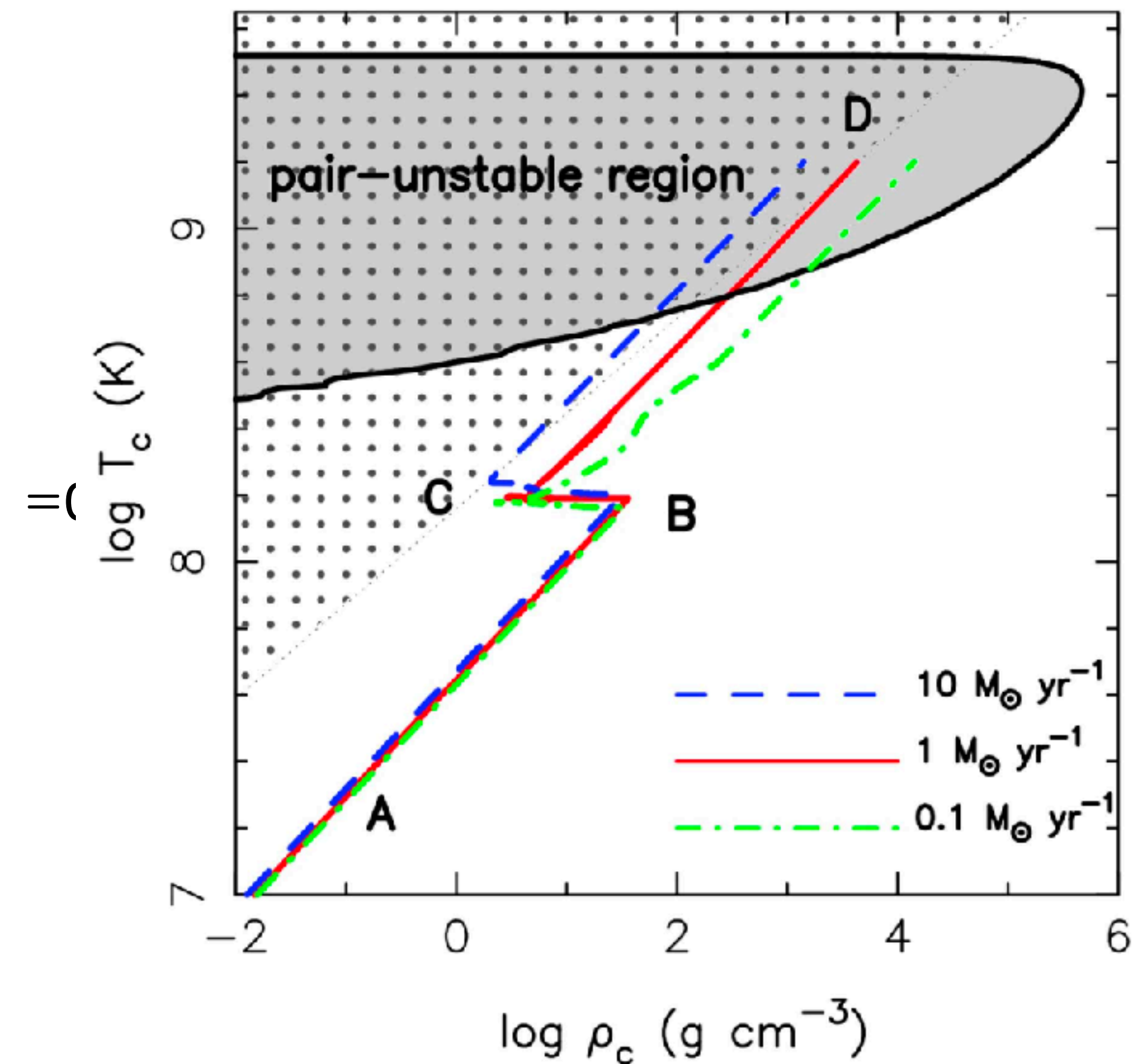
- Hypothetical very massive ($\gtrsim 10^4 M_\odot$) star
- High- s , P_{rad} -dominant ($\Gamma \approx 4/3$)
- Dies likely by GR instability

$$\rho_{\text{crit}} \approx 1.994 \times 10^{18} \left(\frac{0.5}{\mu}\right)^3 \left(\frac{M_\odot}{M}\right)^{7/2} \text{ g cm}^{-3}$$

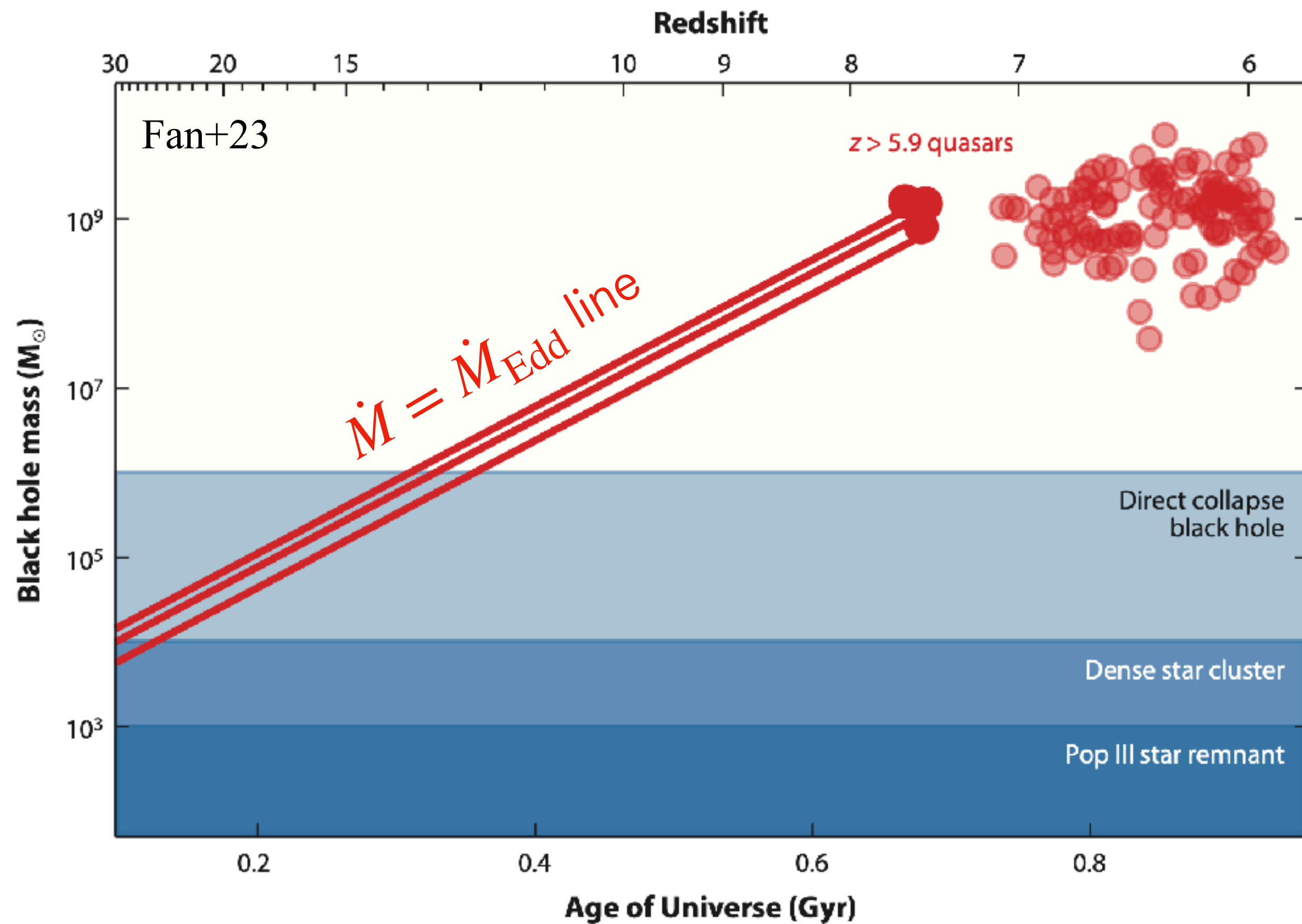
$$\approx 4 \text{ g/cm}^3 (M/10^5 M_\odot)^{-3.5} (\mu/0.59)^{-3}$$

(Shapiro-Teukolsky 83, Fuller+86)

Umeda+16



SMBHs in early universe



Distant (high- z) SMBH with $M \sim 10^9 M_{\odot}$

How are they formed?

Basic idea: $\dot{M} \lesssim \dot{M}_{\text{Edd}} \propto M$.

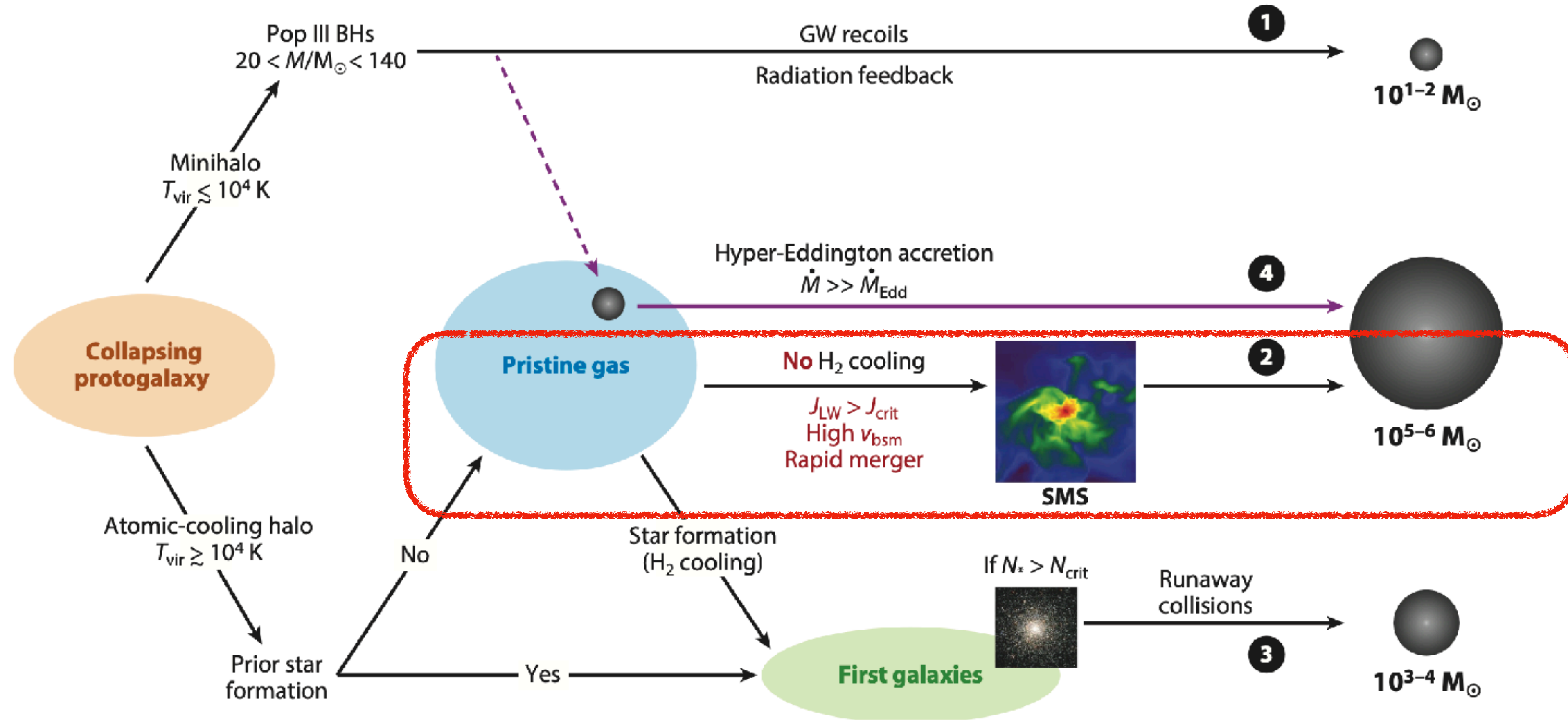
(Super/hyper-Eddington accretion may be possible)

Keeping a high Eddington ratio $\dot{M}/\dot{M}_{\text{Edd}}$

from $10^2 M_{\odot}$ to $10^9 M_{\odot}$ may not be easy

Initial high mass of BH may make it easier!

Direct collapse scenario for SMBHs in early universe



Inayoshi+20, see also Rees 1978

Explosion of supermassive stars: thermonuclear case

Chen+14

Nagele+20, 22

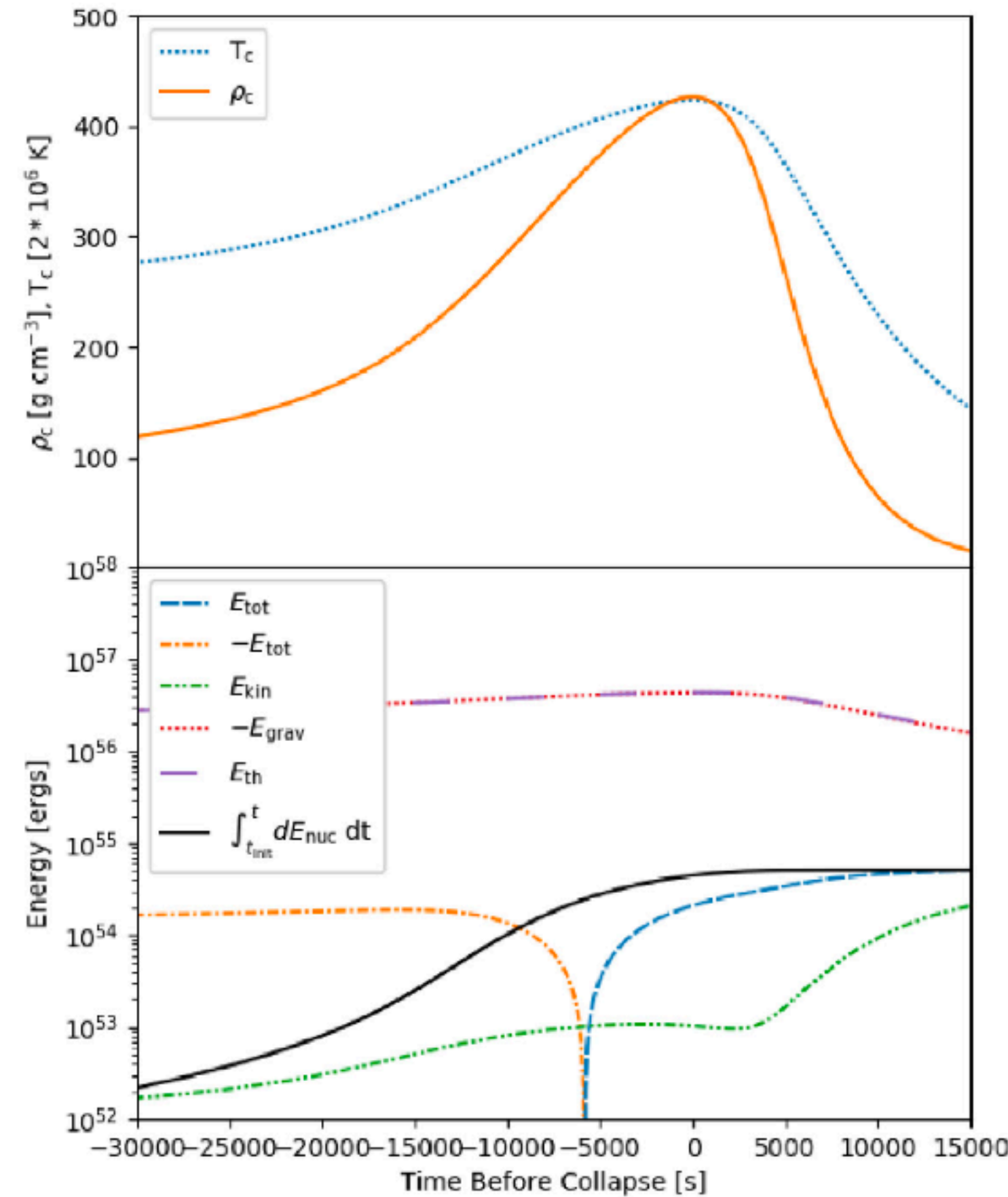
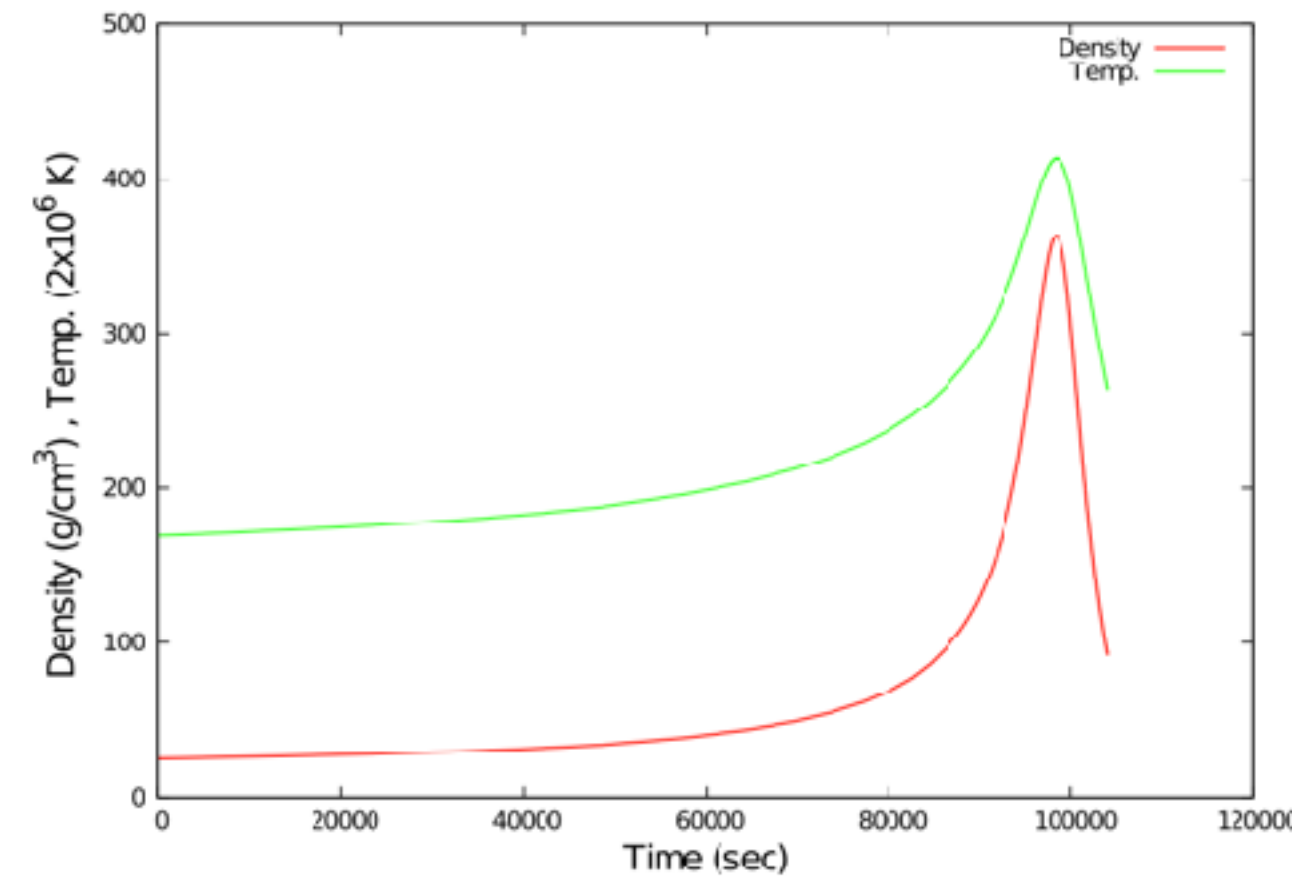
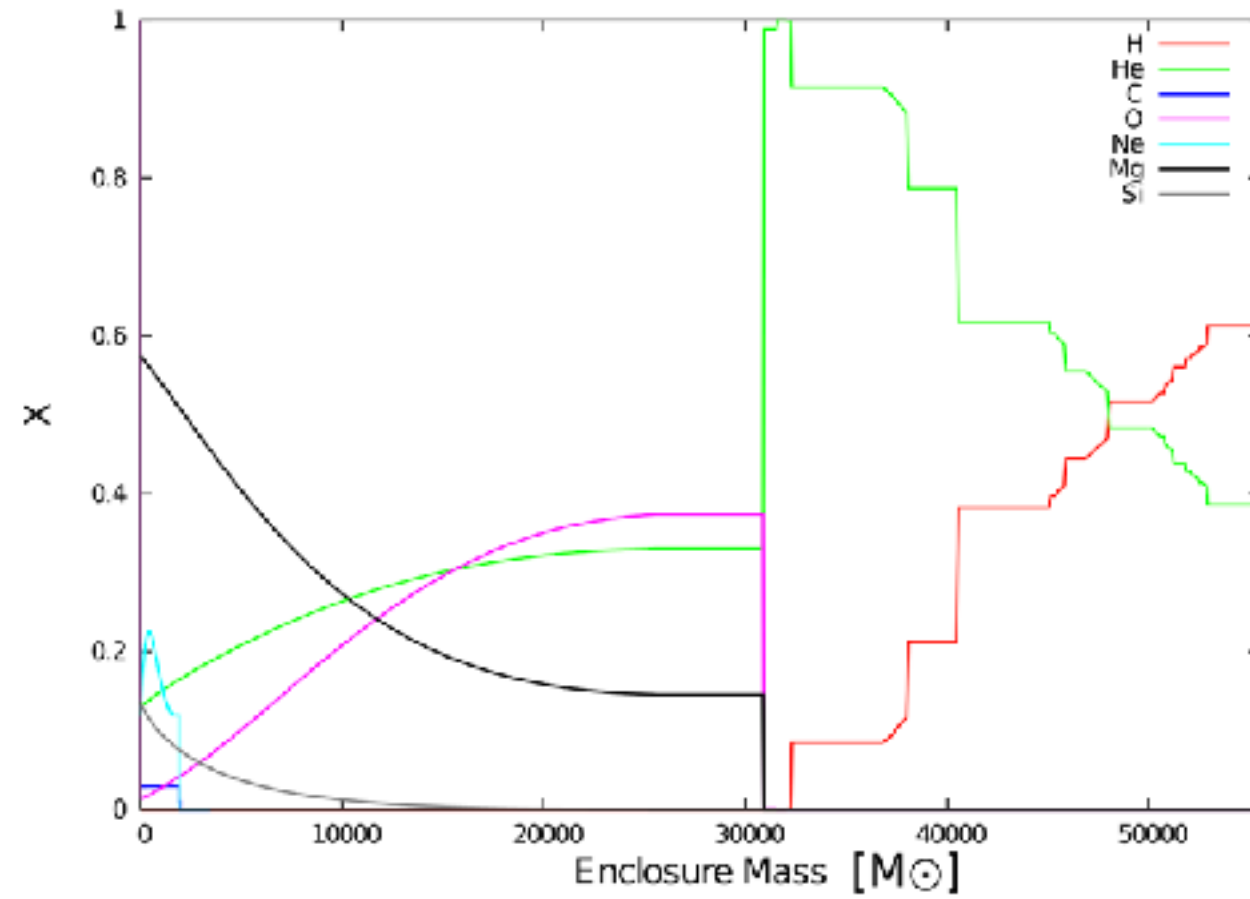


Table 1

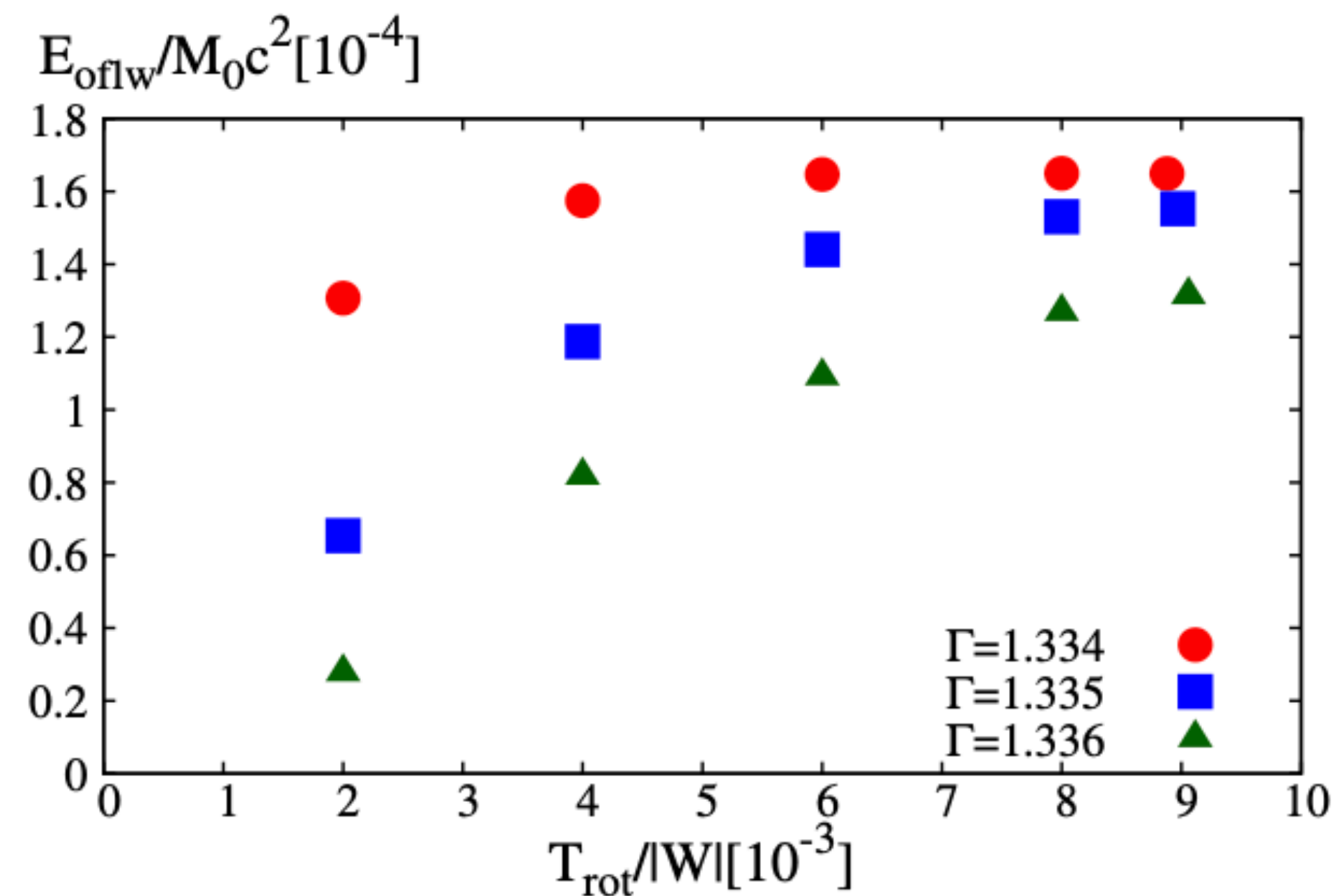
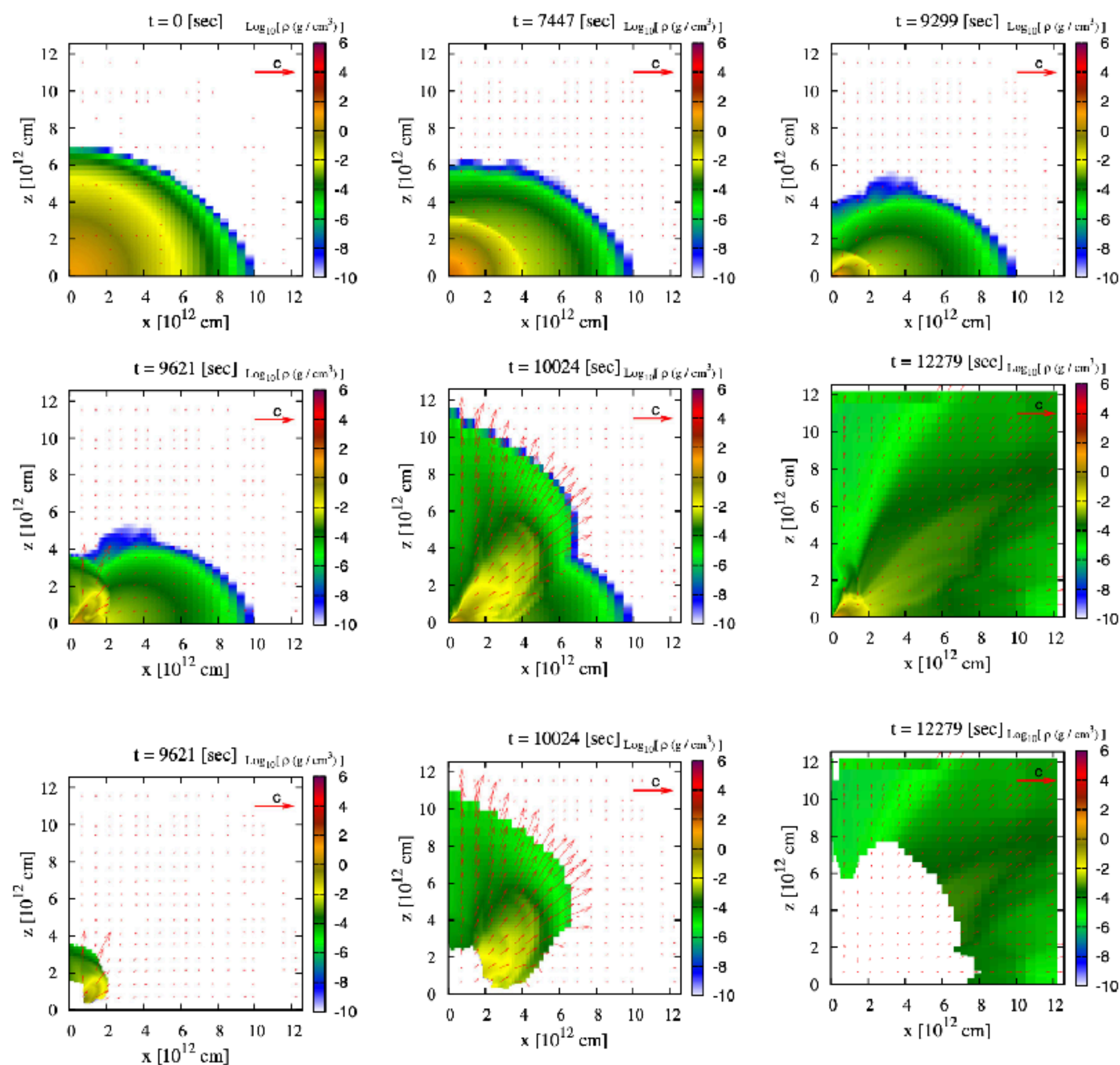
Elemental Masses Before and After the Explosion

Isotope	1H (M_{\odot})	4He (M_{\odot})	^{12}C (M_{\odot})	^{16}O (M_{\odot})	^{20}Ne (M_{\odot})	^{24}Mg (M_{\odot})	^{28}Si (M_{\odot})	Total (M_{\odot})
Before	8336	24902	922	7972	5110	7748	515	55505
After	8335	24145	919	6856	4819	8943	1485	55502
ΔM	-1	-757	-3	-1116	-291	1195	970	-3

$^{16}O(\alpha, \gamma)^{20}Ne(\alpha, \gamma)^{24}Mg$ is responsible for explosion

Effects of rotation: bounce-induced explosion

Uchida+17



Method: Numerical setup

General relativistic gravity

Hydrodynamics with nuclear reaction

$H \rightarrow He \rightarrow C$ (Only forward reaction)

CNO cycle triple- α

Equation of state

Composite of ions(H, He, C), photons, electrons and positrons

Neutrino radiation

Only for neutrinos emitted by CNO cycle $\sim 8\%$ of heating rate

Method: Initial supermassive star models

Marginally stable SMS with rotation.

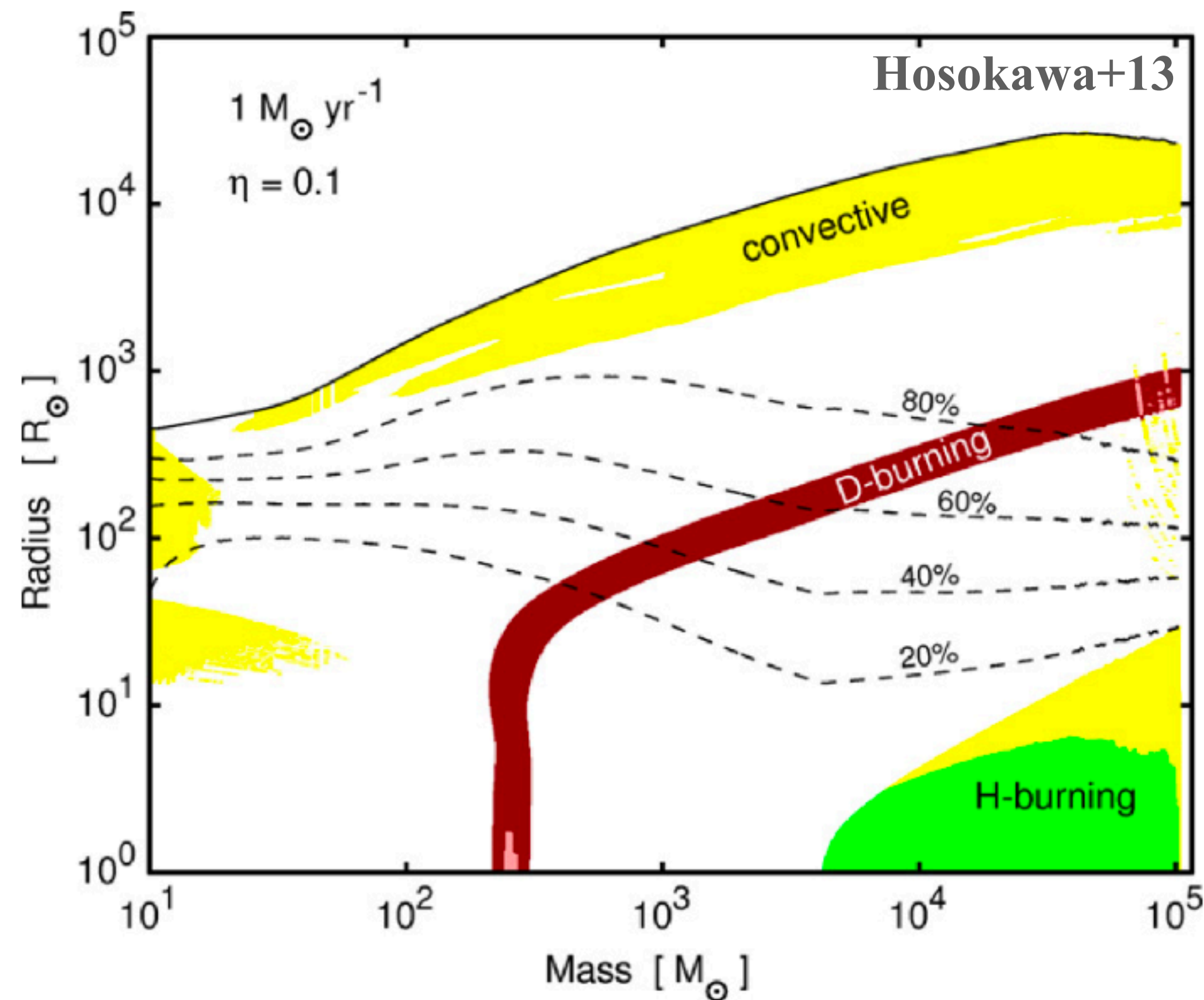
model	$M_0 (M_\odot)$	$R_{e0} \text{ (cm)}$	$T_{\text{kin}}/ W $	$\alpha_{c,0}$	$\gamma_{c,0} - 4/3$	\hat{A}	s/k_B
H1	2.1×10^5	1.7×10^{13}	0.002	0.992	0.0026	∞	450
H2	3.2×10^5	2.3×10^{13}	0.004	0.990	0.0021	∞	550
H3	4.3×10^5	2.7×10^{13}	0.006	0.988	0.0018	∞	630
H4	6.9×10^5	4.4×10^{13}	0.009	0.985	0.0014	∞	800
Hdif1	9.2×10^5	5.0×10^{13}	0.011	0.983	0.0012	2	920
Hdif2	1.1×10^6	5.3×10^{13}	0.013	0.981	0.0012	1.5	1000
Hdif3	1.9×10^6	7.4×10^{13}	0.018	0.976	0.0009	1.0	1300
He1	5.0×10^4	4.3×10^{12}	0.002	0.992	0.0023	∞	210
He2	7.1×10^4	5.1×10^{12}	0.004	0.990	0.0019	∞	250
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He4	1.6×10^5	1.0×10^{13}	0.009	0.985	0.0013	∞	380

Primordial composition
 $X(\text{H})=0.25, X(\text{He})=0.75$

Purely He star
 $X(\text{He})=1$

Method: Initial supermassive star models

Caution: here are only the isentropic “core” of SMS



Realistic SMS may have inflated envelope

Method: Initial supermassive star models

Marginally stable SMS with rotation.

rotation

model	$M_0 (M_\odot)$	$R_{e0} \text{ (cm)}$	$T_{\text{kin}}/ W $	$\alpha_{c,0}$	$\gamma_{c,0} - 4/3$	\hat{A}	s/k_B
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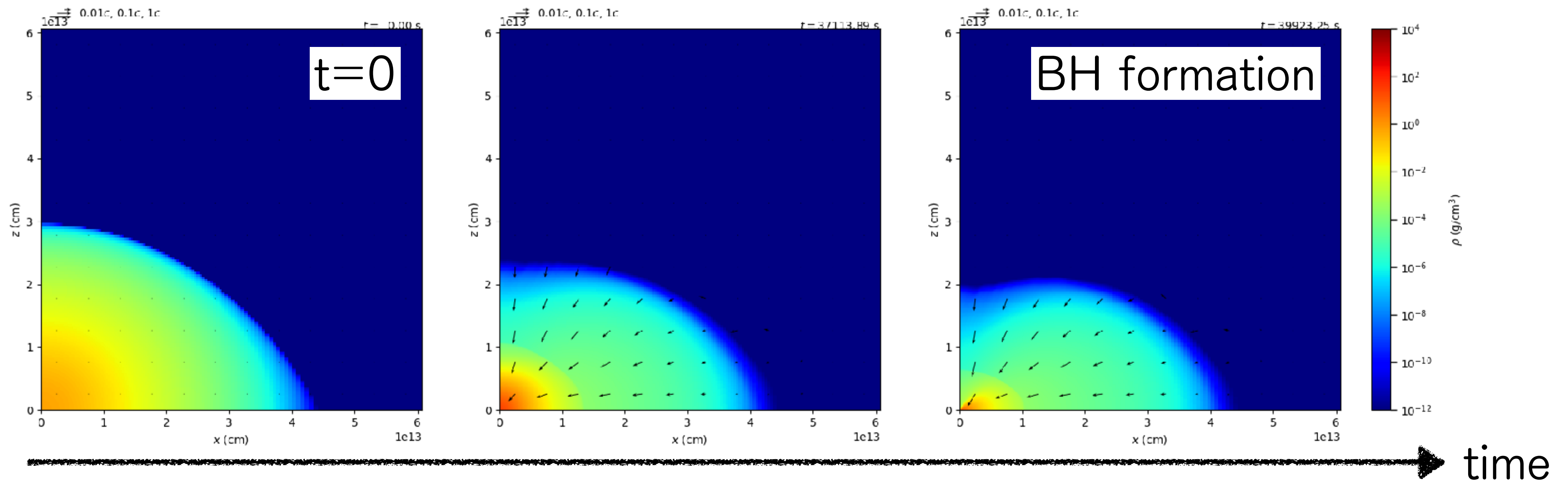
More massive

- more radiation-dominant
- $\gamma \rightarrow 4/3$
- higher entropy

More rapid rotation - more stable
larger mass for being unstable

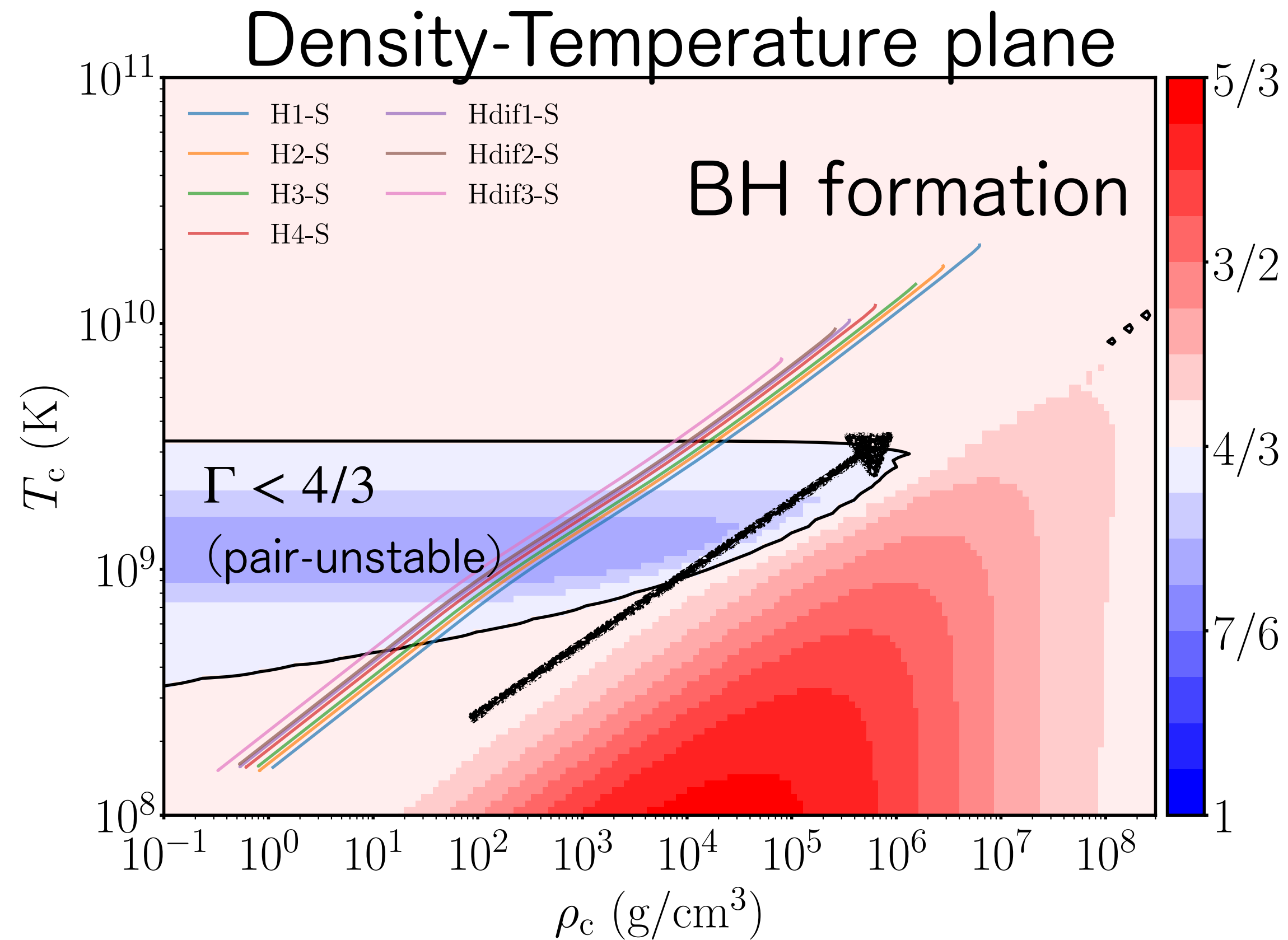
Result: Outline of evolution

Primordial composition, mass-shedding case



Collapsing motion is \sim coherent (characteristic of GR instability)

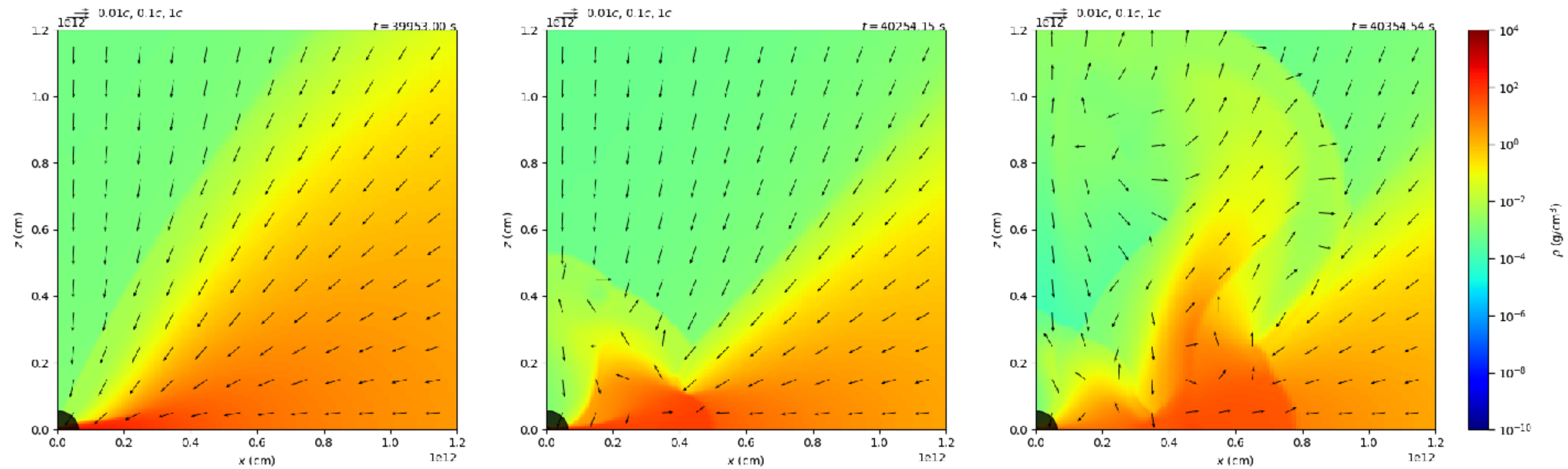
Result: Outline of evolution



Collapse proceeds
outside pair-unstable region.
→ GR instability

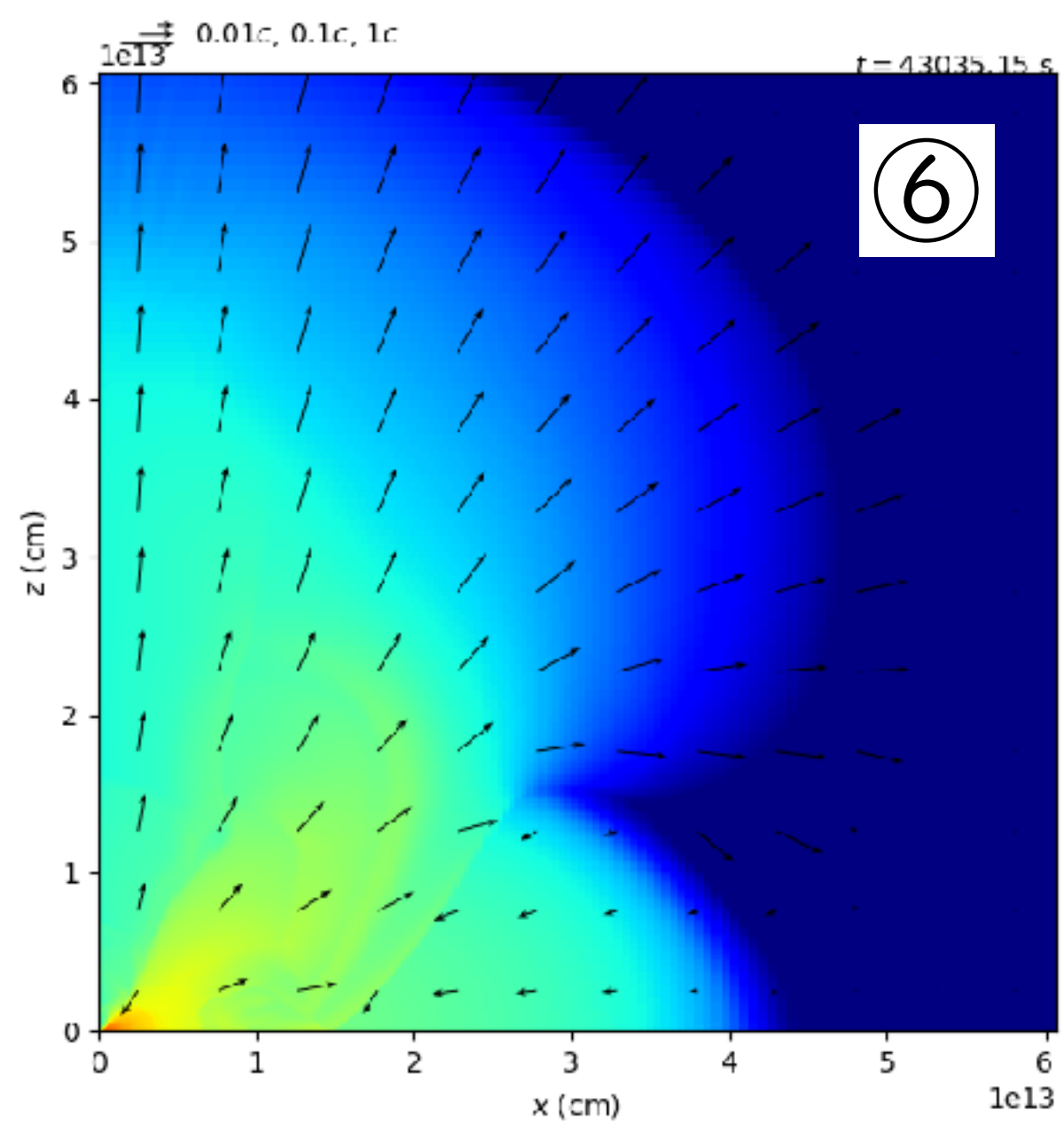
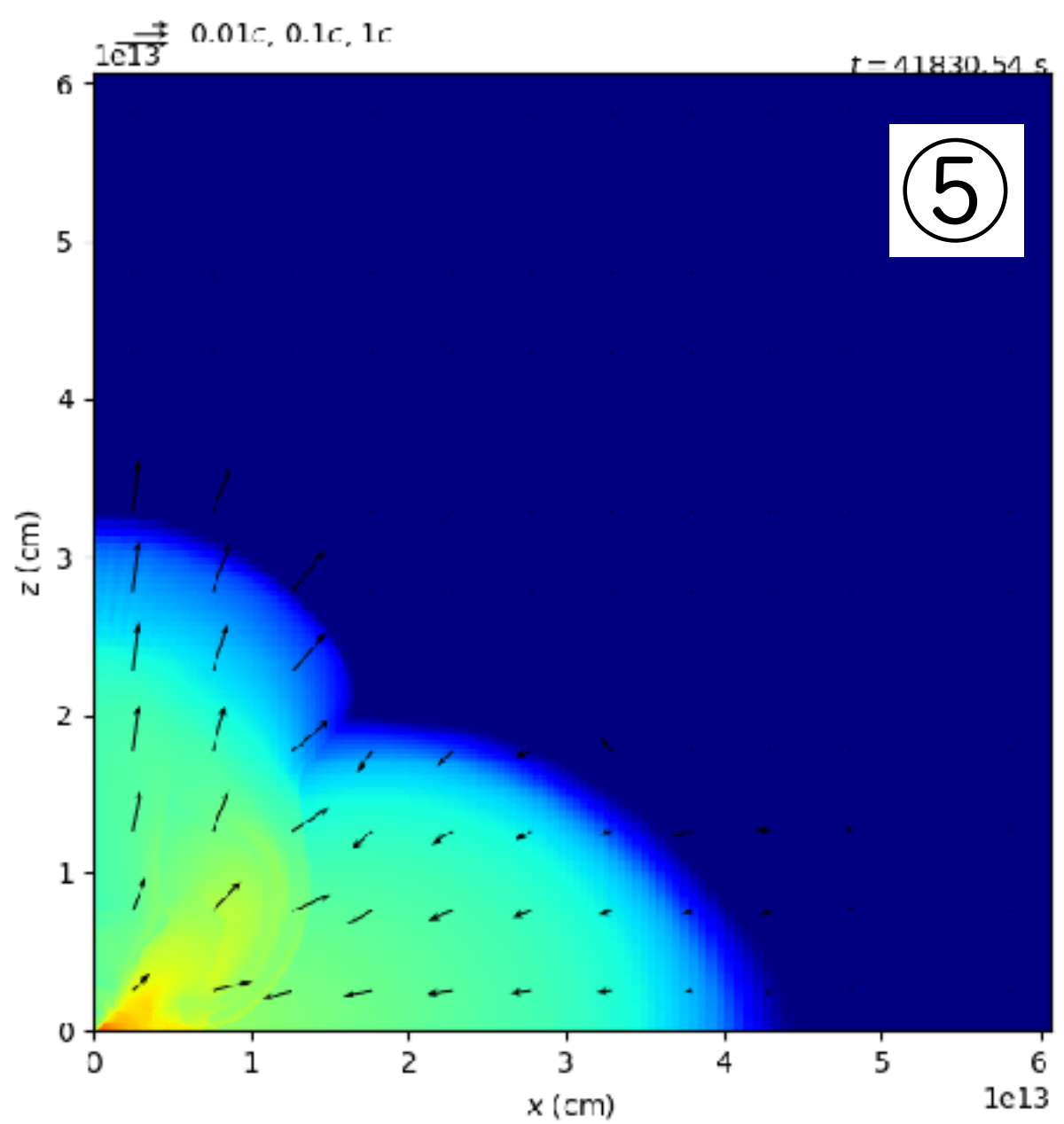
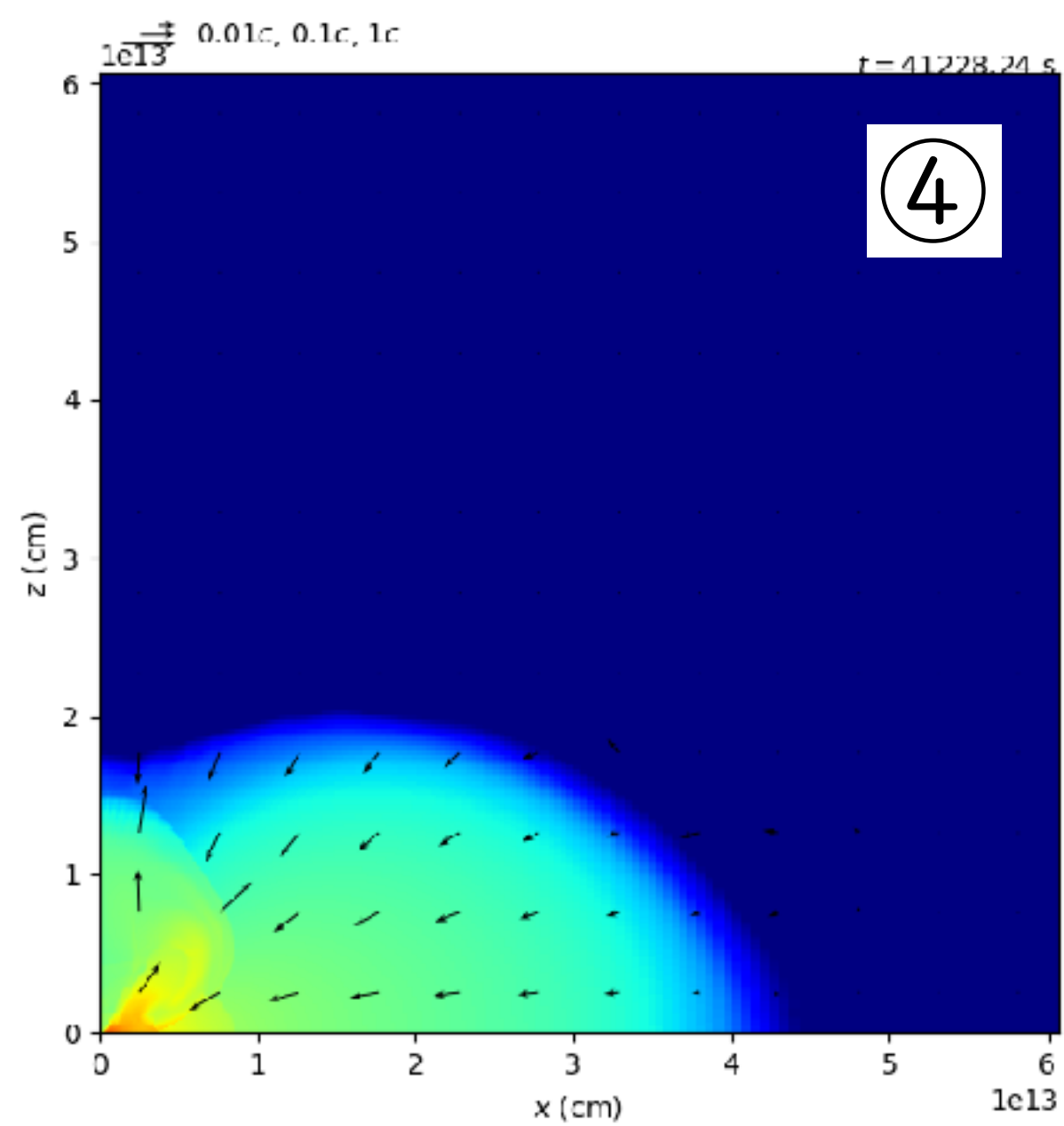
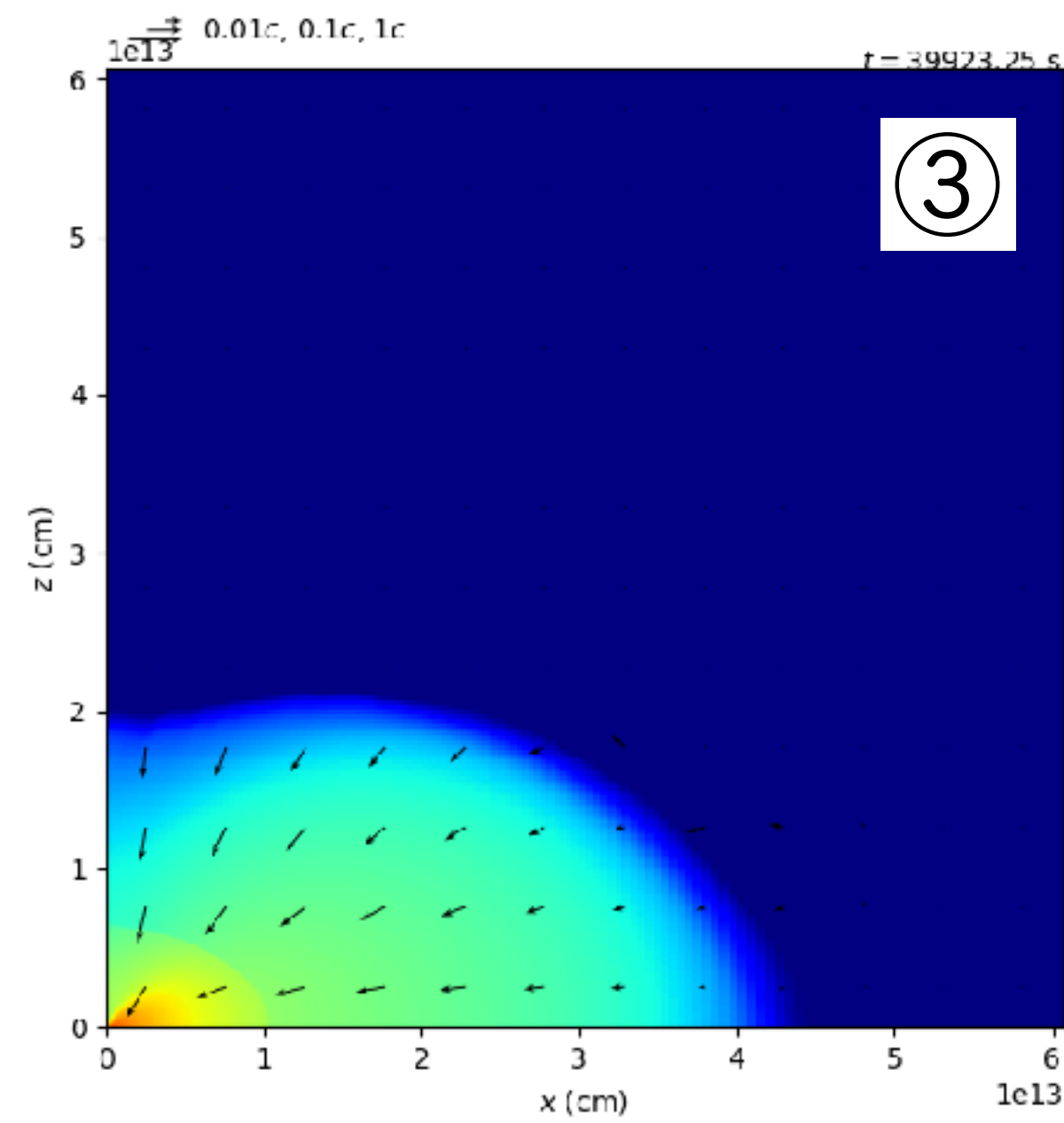
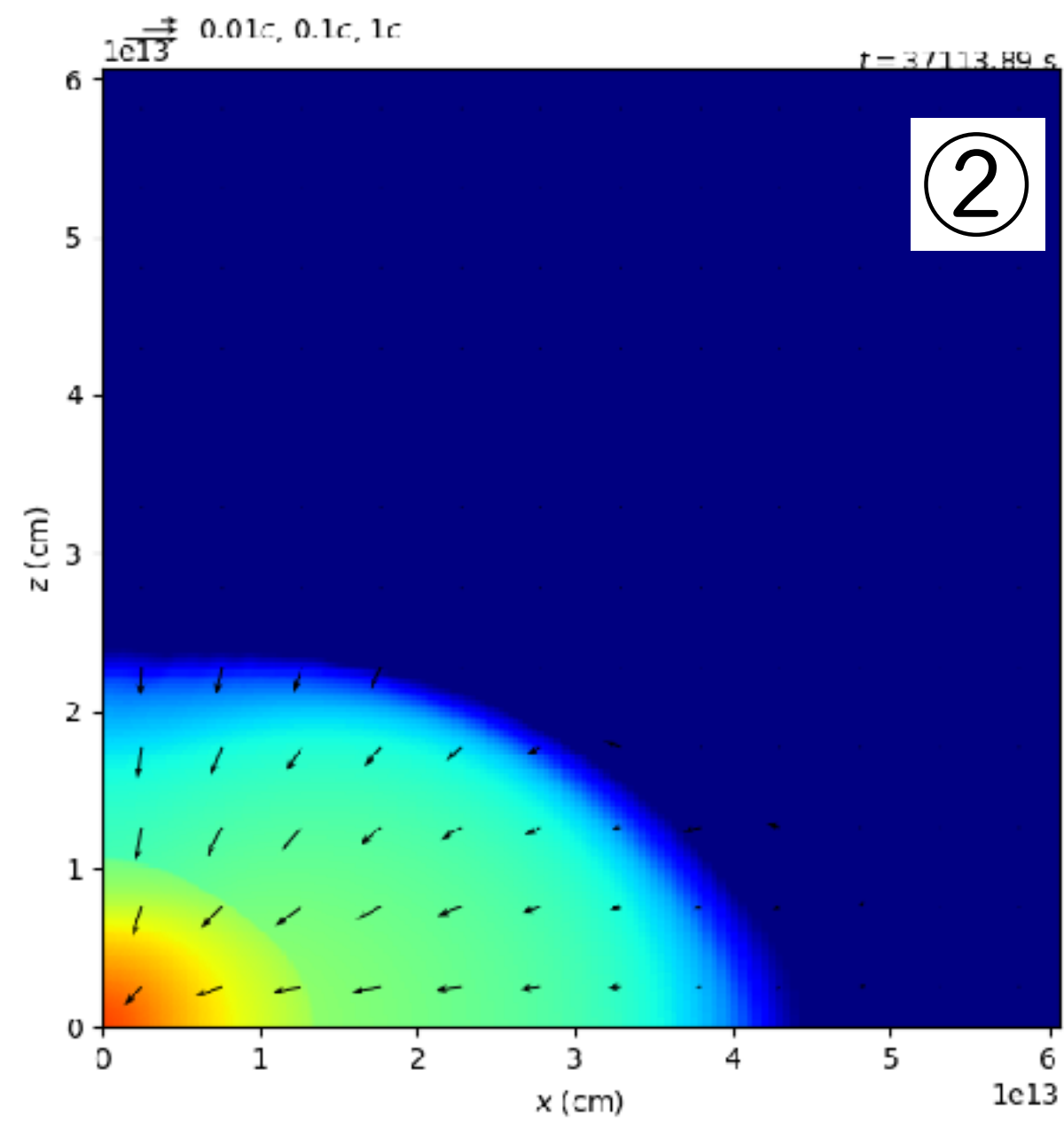
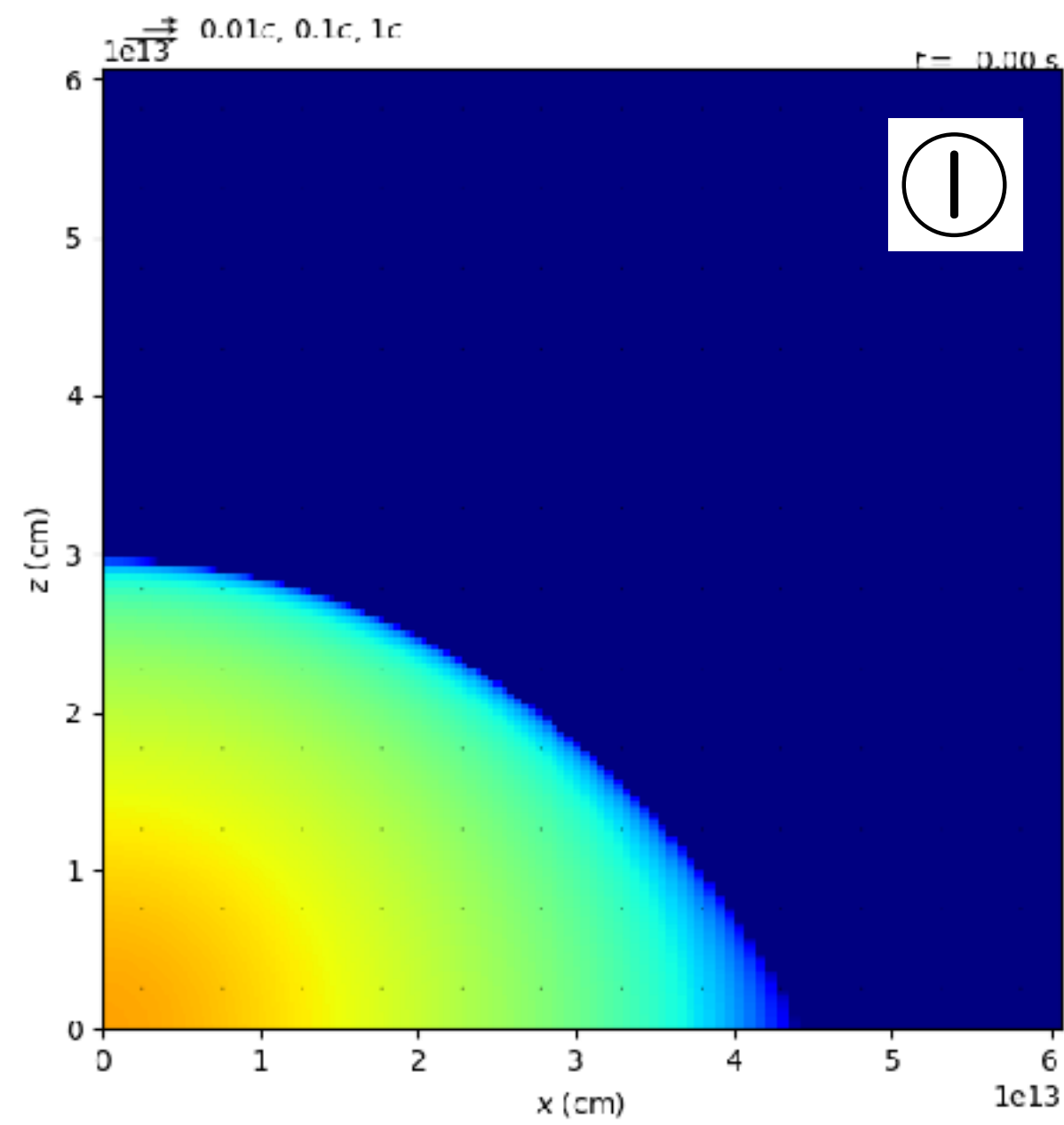
Result: Bounce-shock-induced ejecta

Density snapshots around torus formation time.

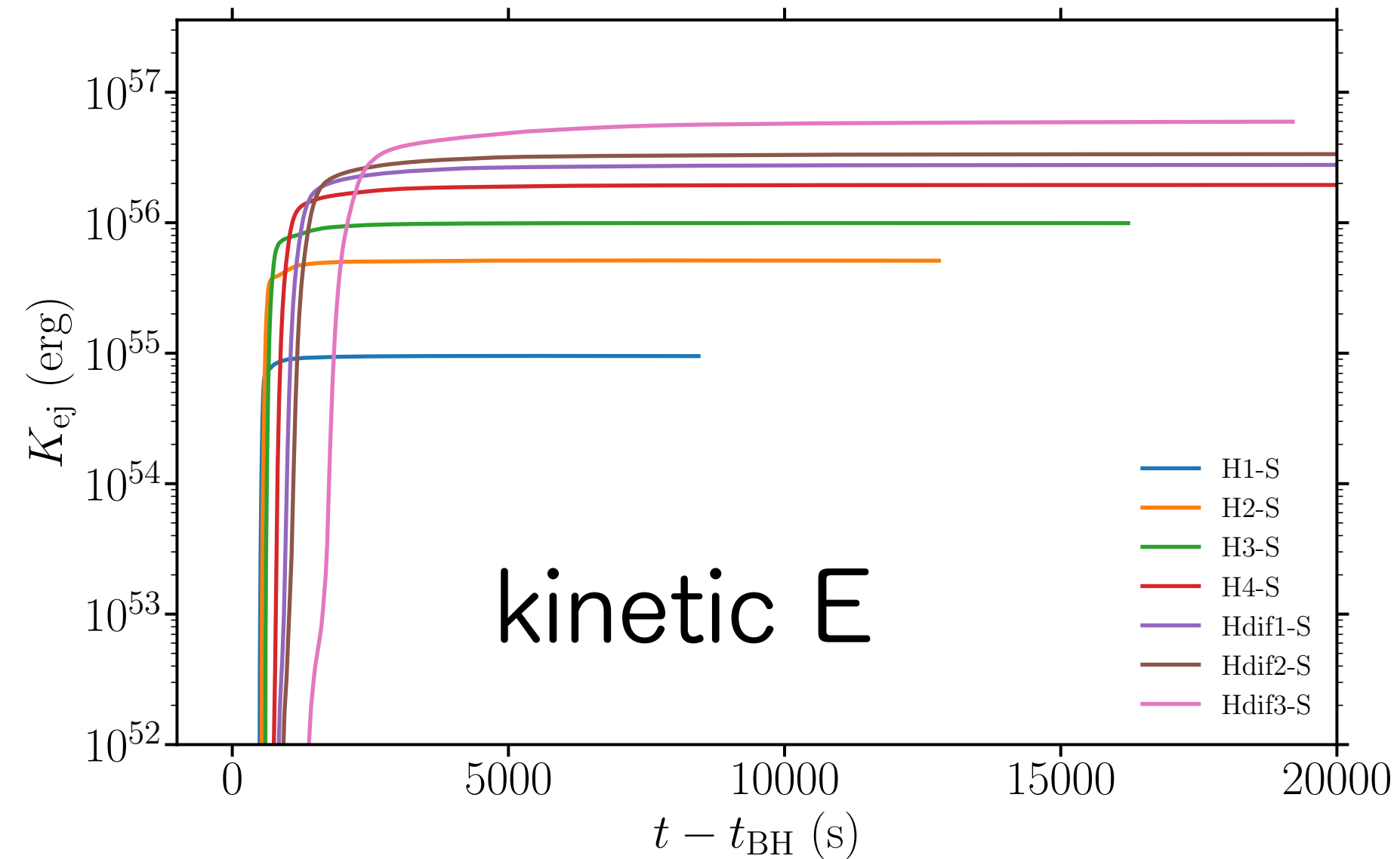
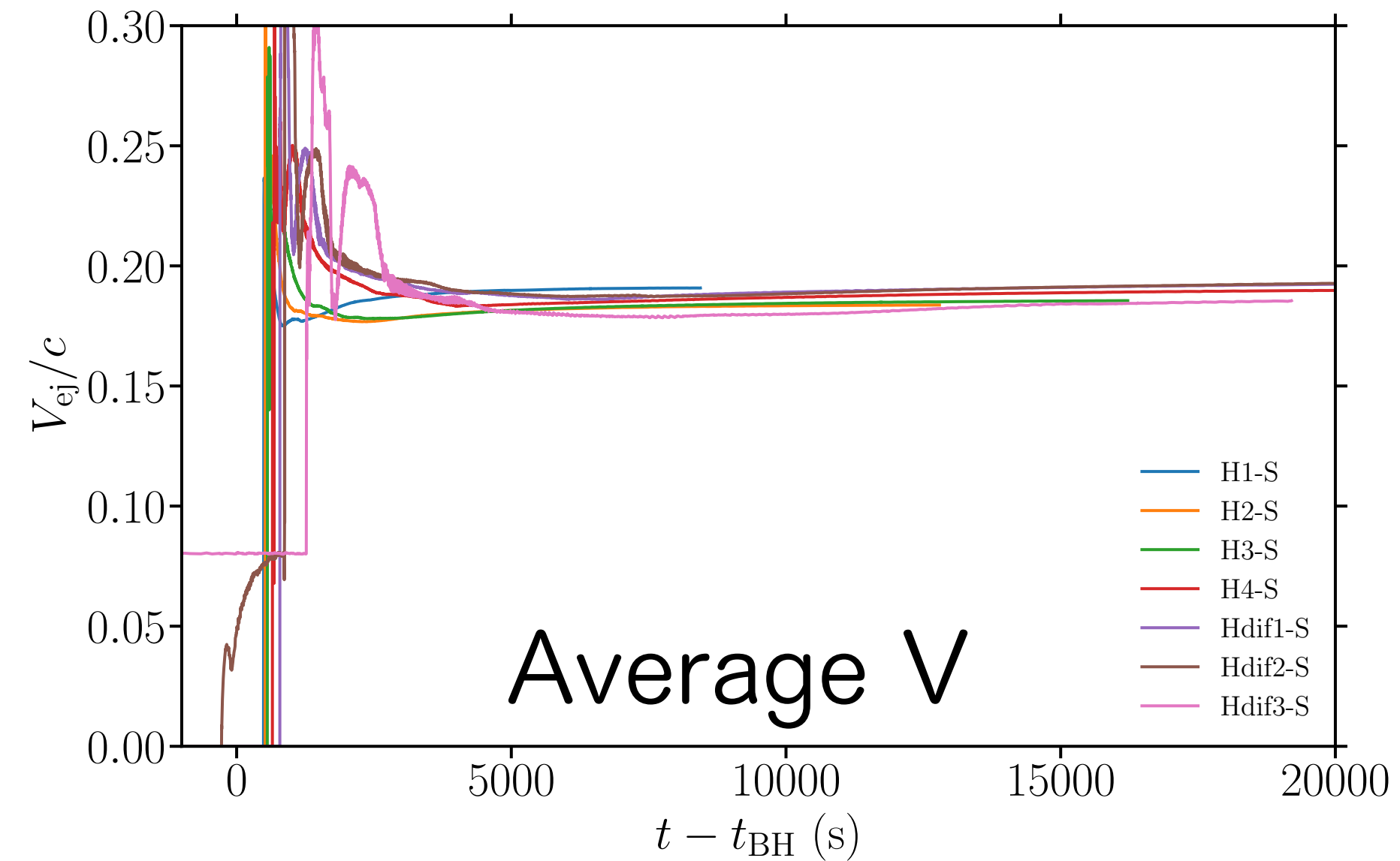
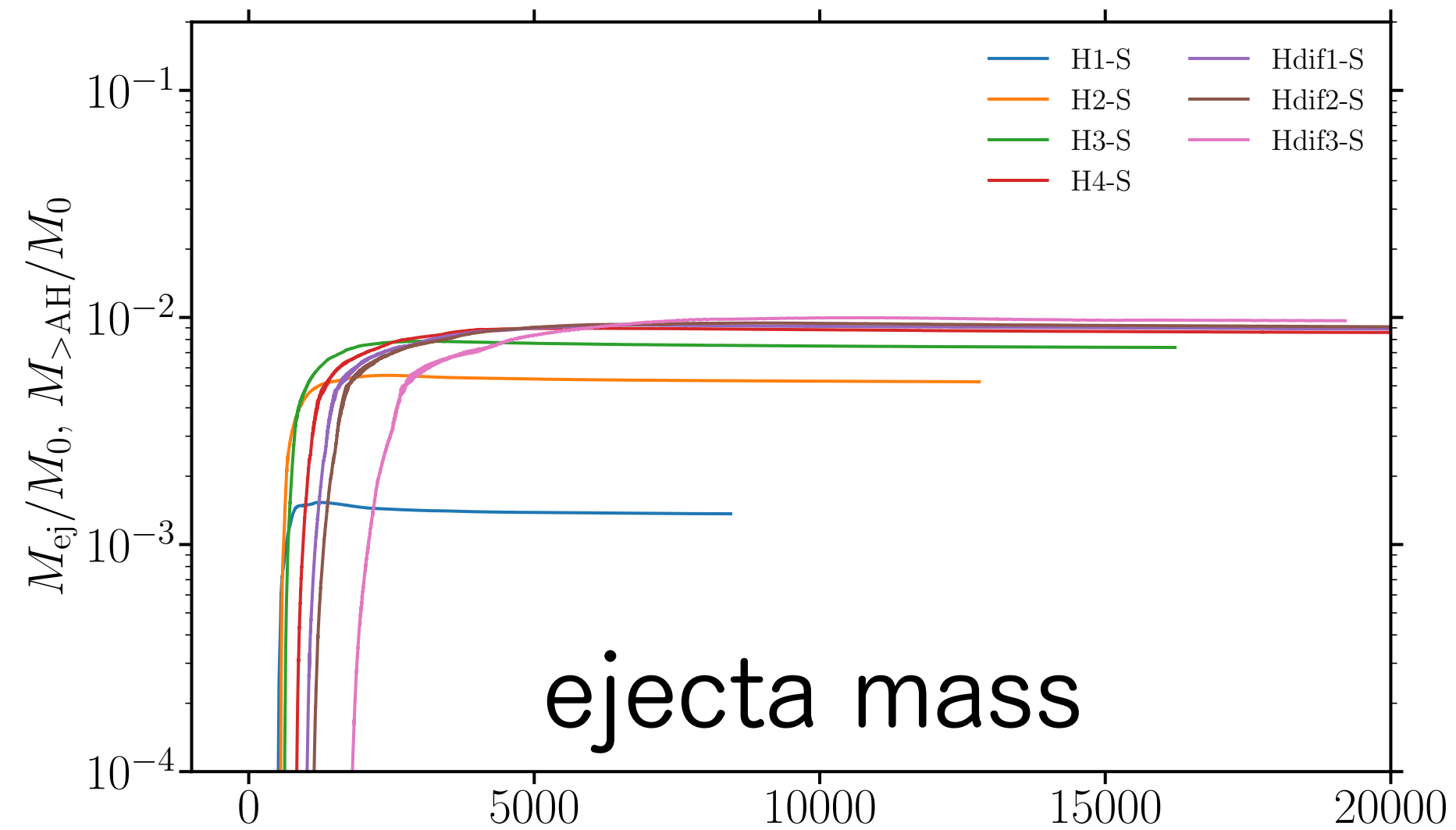


time

Sudden formation of centrifugally supported torus induces its bounce

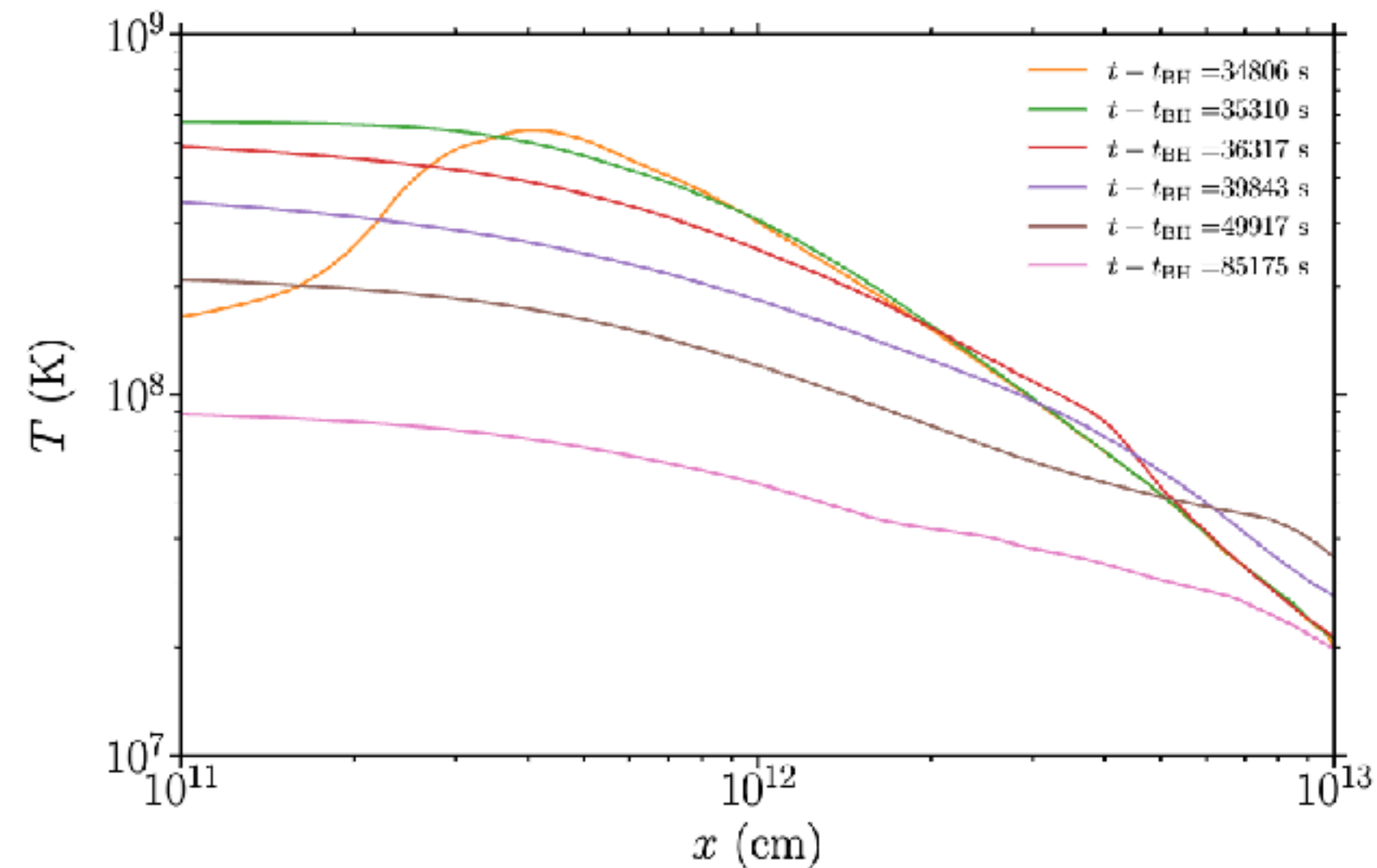
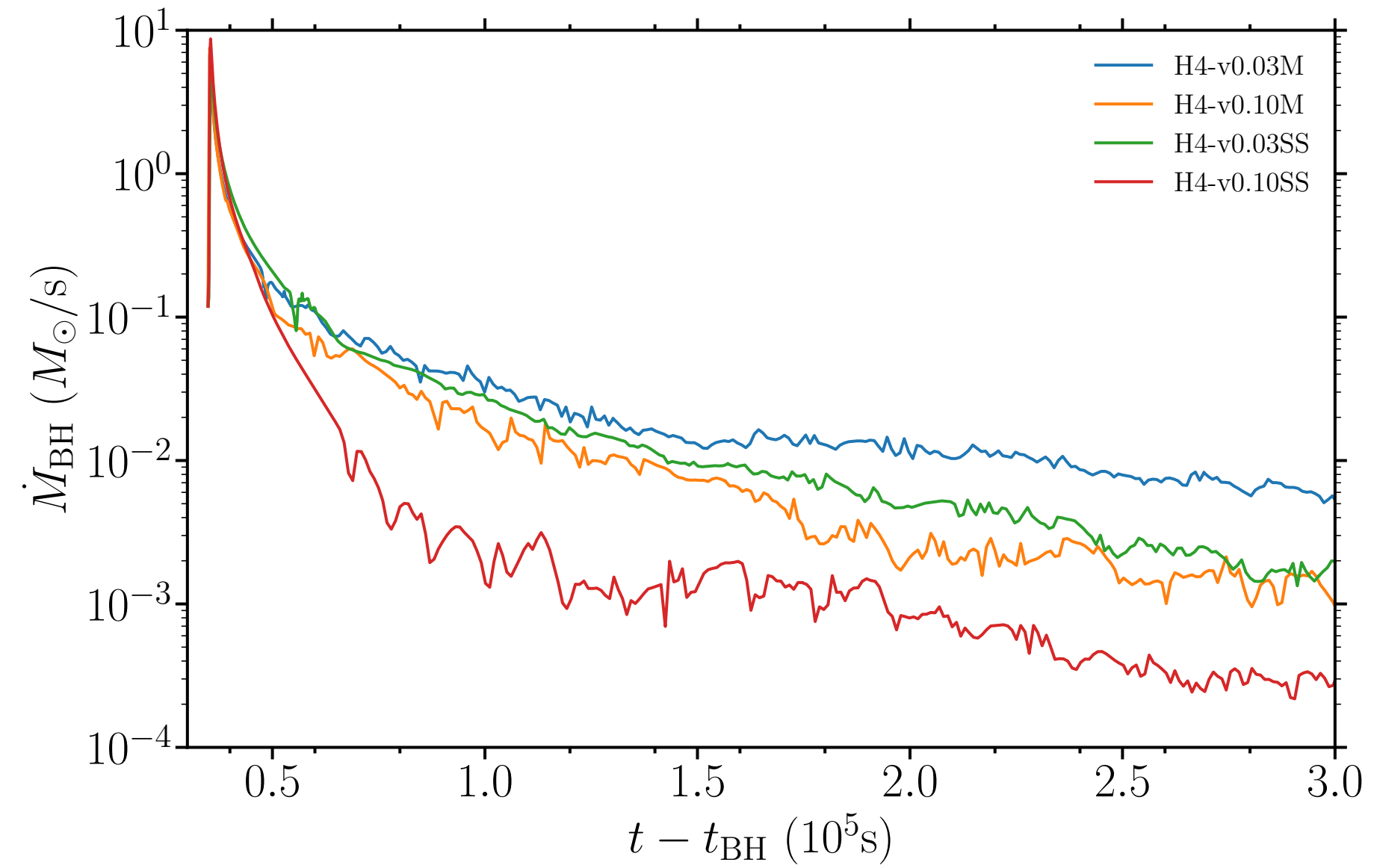
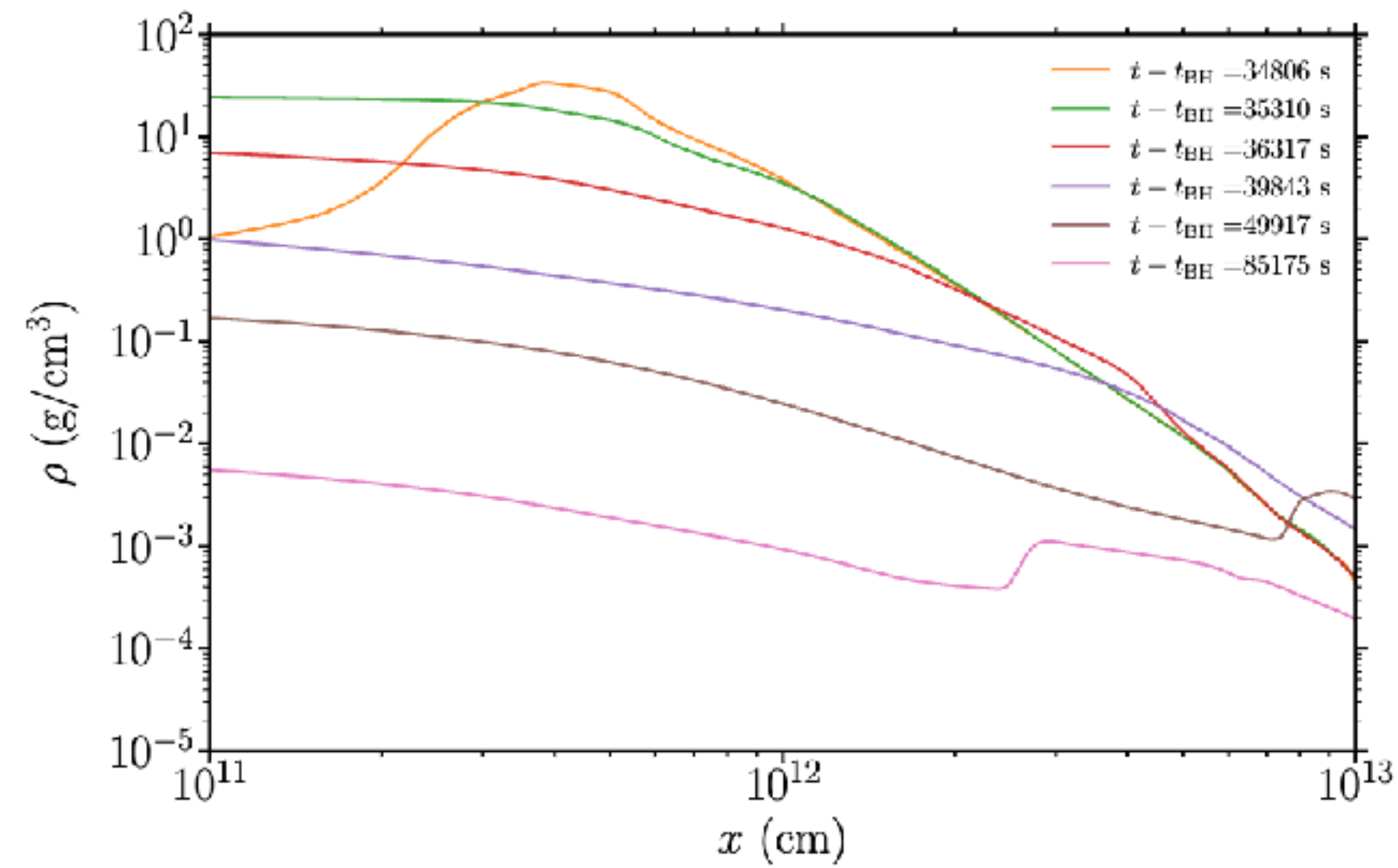


Result: Properties of ejecta



- Ejecta mass $\sim 1\%$ of initial SMS mass for fast-rotating SMS.
- Velocity $\sim 0.2 c$
- Kinetic E $\sim 10^{-4} M c^2$

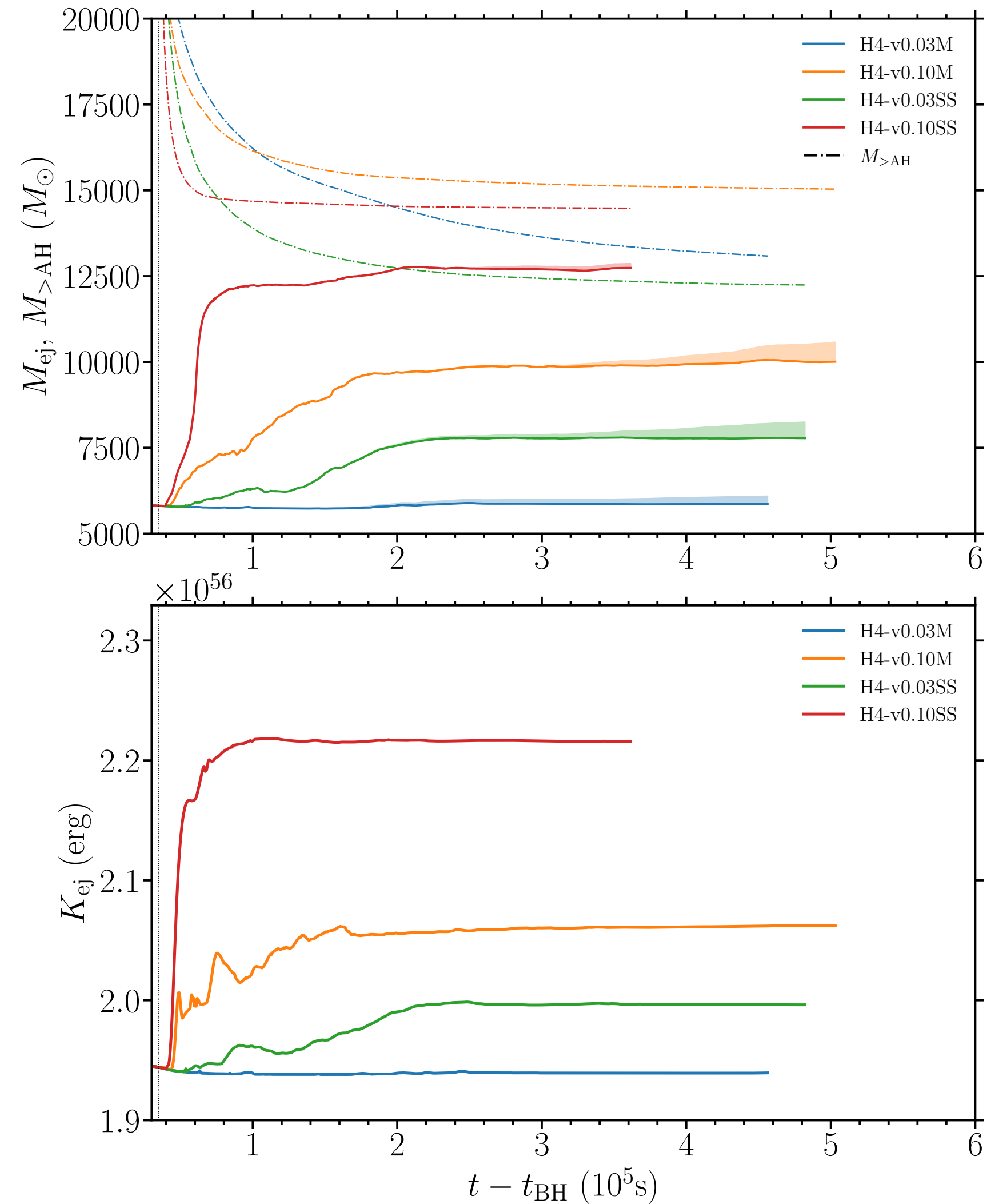
Result: Viscous evolution of the disk



Accretion timescale $\sim 10^4$ s

Upto $\sim 10 M_{\odot}/\text{s} \sim 10^{13} \dot{M}_{\text{Edd}}$ (Hyper-Eddington)

Result: Viscous evolution of the disk

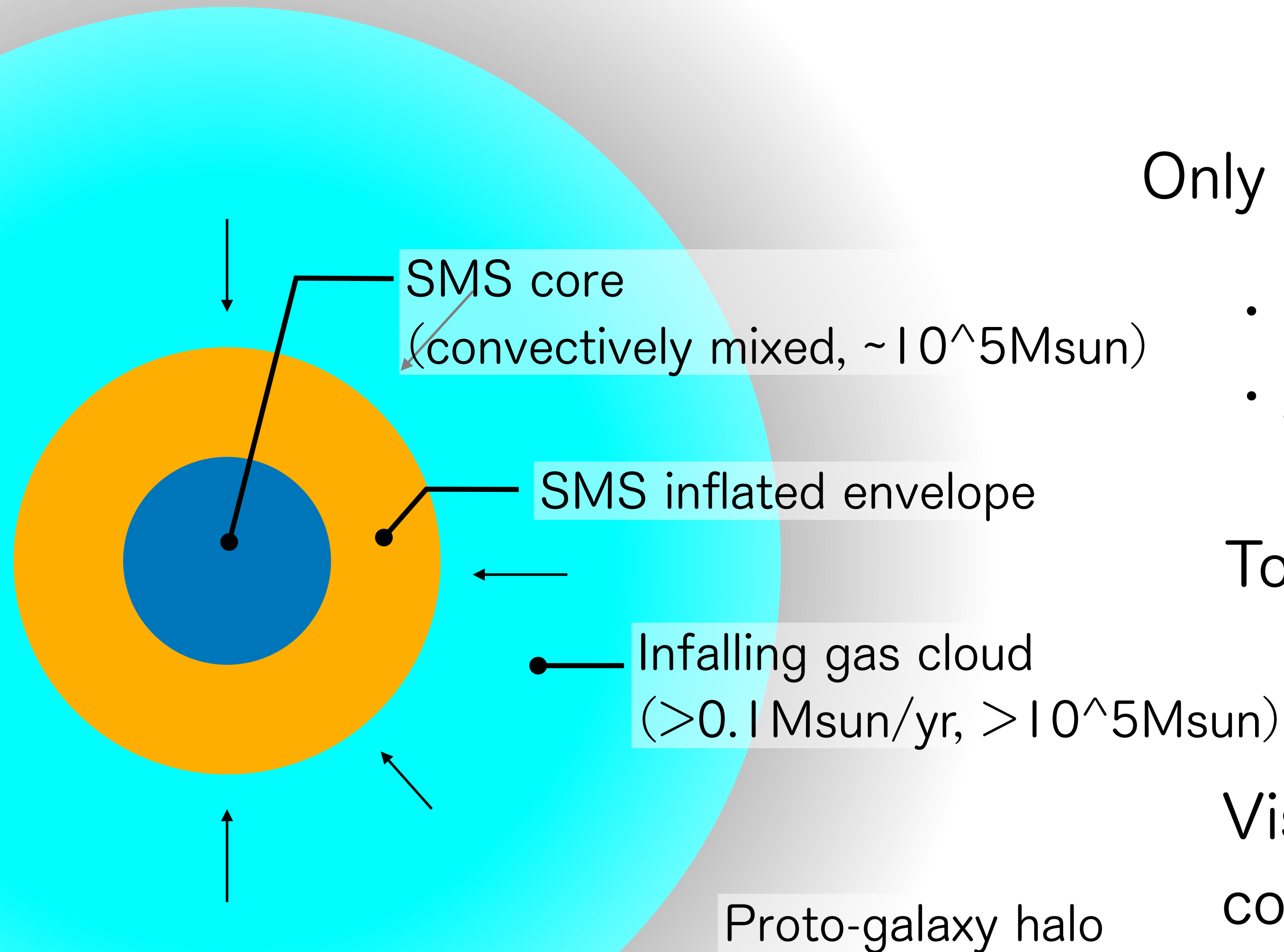


Viscosity-driven ejecta
(with different prescription & strength)

Ejecta mass can be $\sim 3 \times$ bounce-driven ejecta

Velocity $\sim 0.05 c \approx 1/4 \times$ bounce-driven ejecta
→ effect is minor in kinetic energy

Discussion: Realistic environment



Only the collapse of SMS core is simulated.

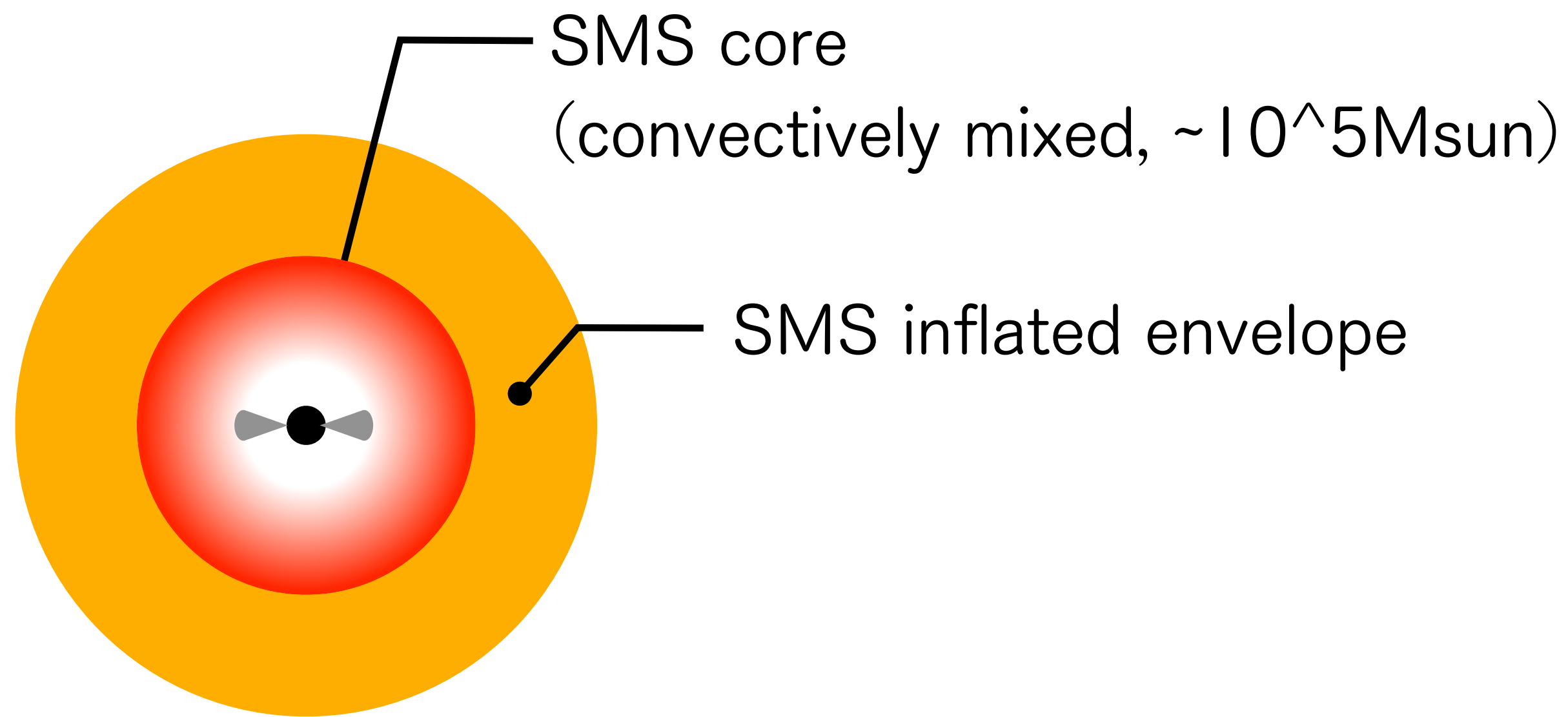
- Envelope
- Atomic cooling cloud \sim SMS mass

Total ejecta mass $\sim 10^5 M_{\text{sun}}$

Viscosity-driven ejecta does not contribute much to total ejecta property

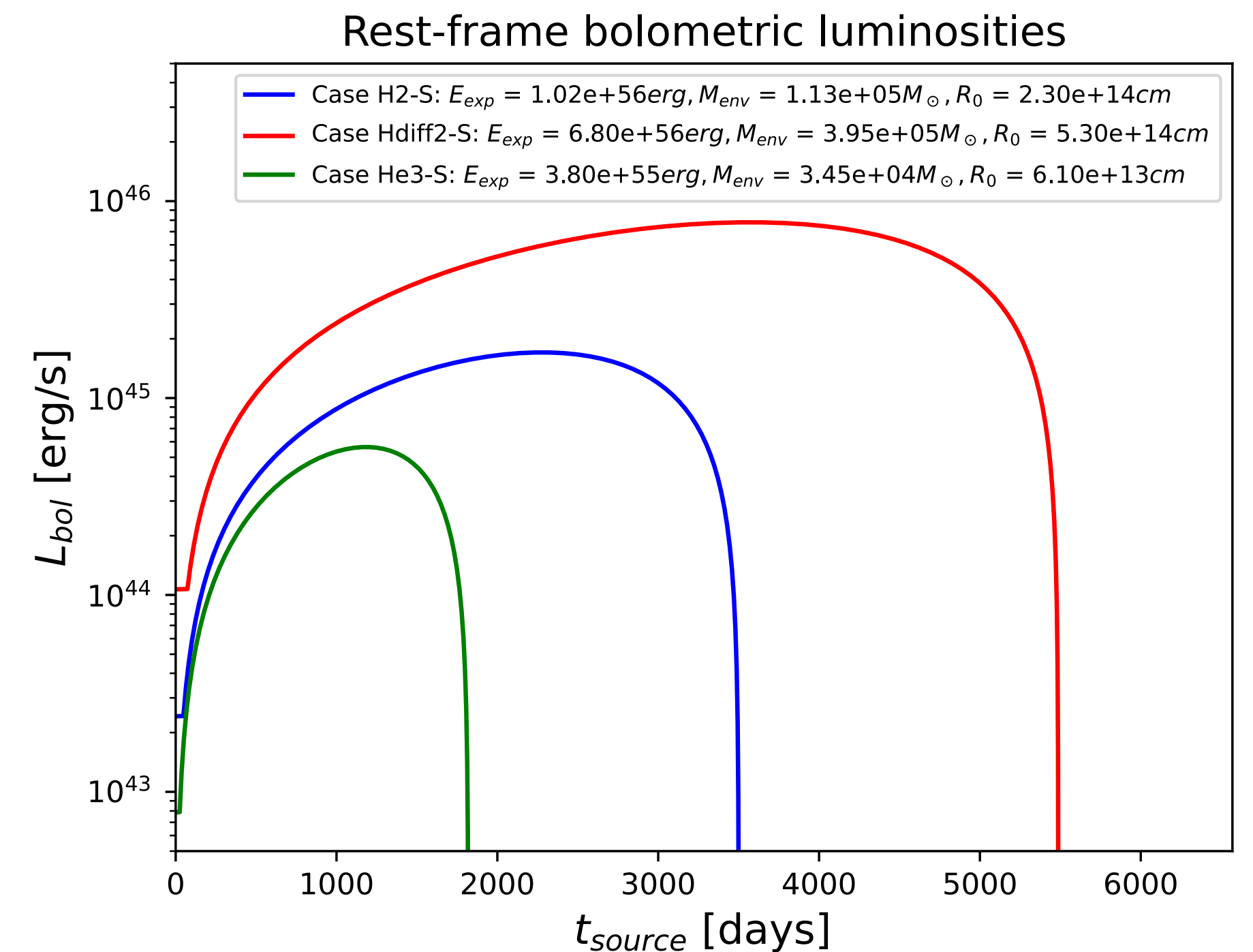
Discussion: Realistic environment

Only the collapse of SMS core is simulated.



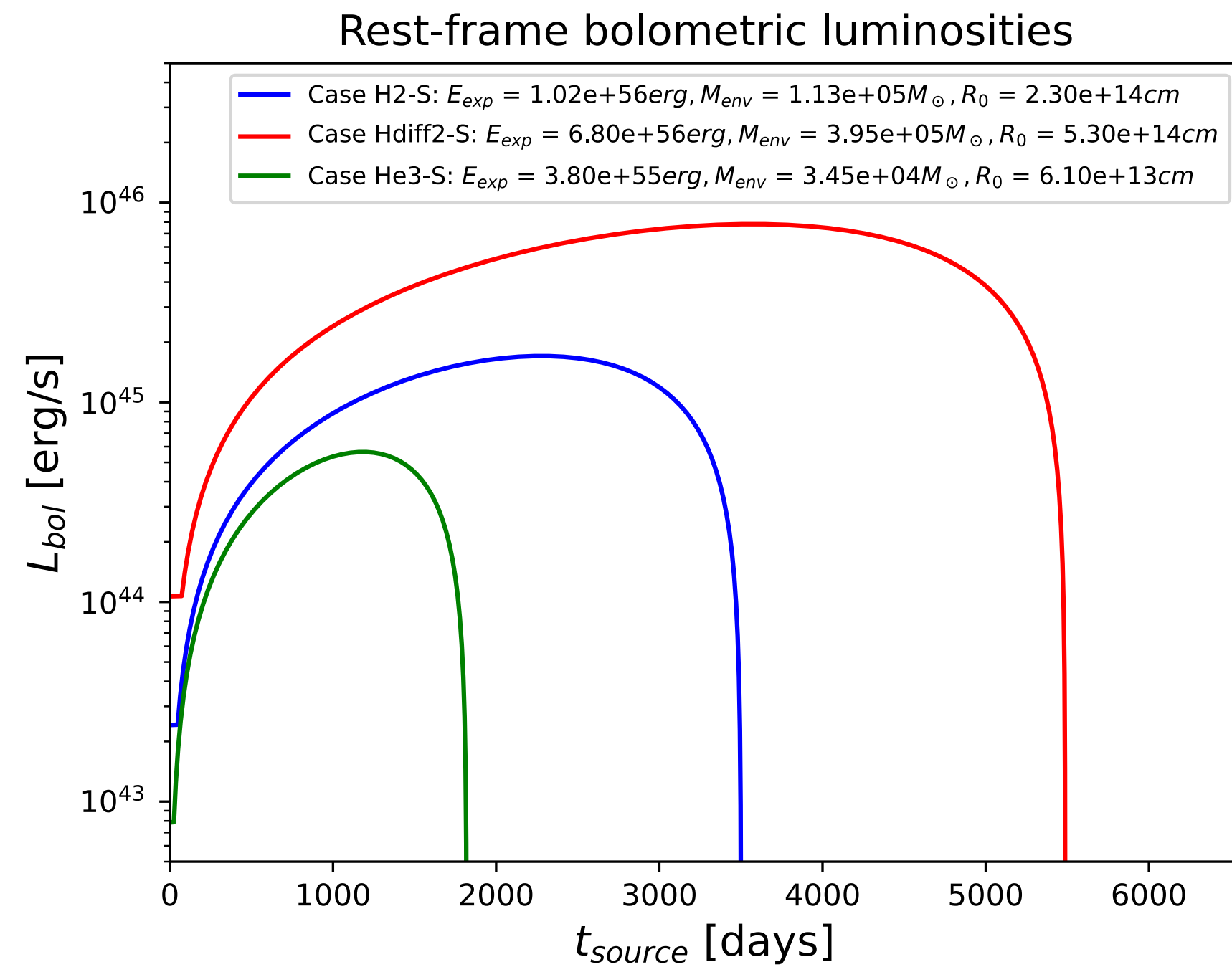
- Envelope $\sim 10\%$ of mass of SMS

Luminous H-rich SN (IIP) cf. Moriya+21

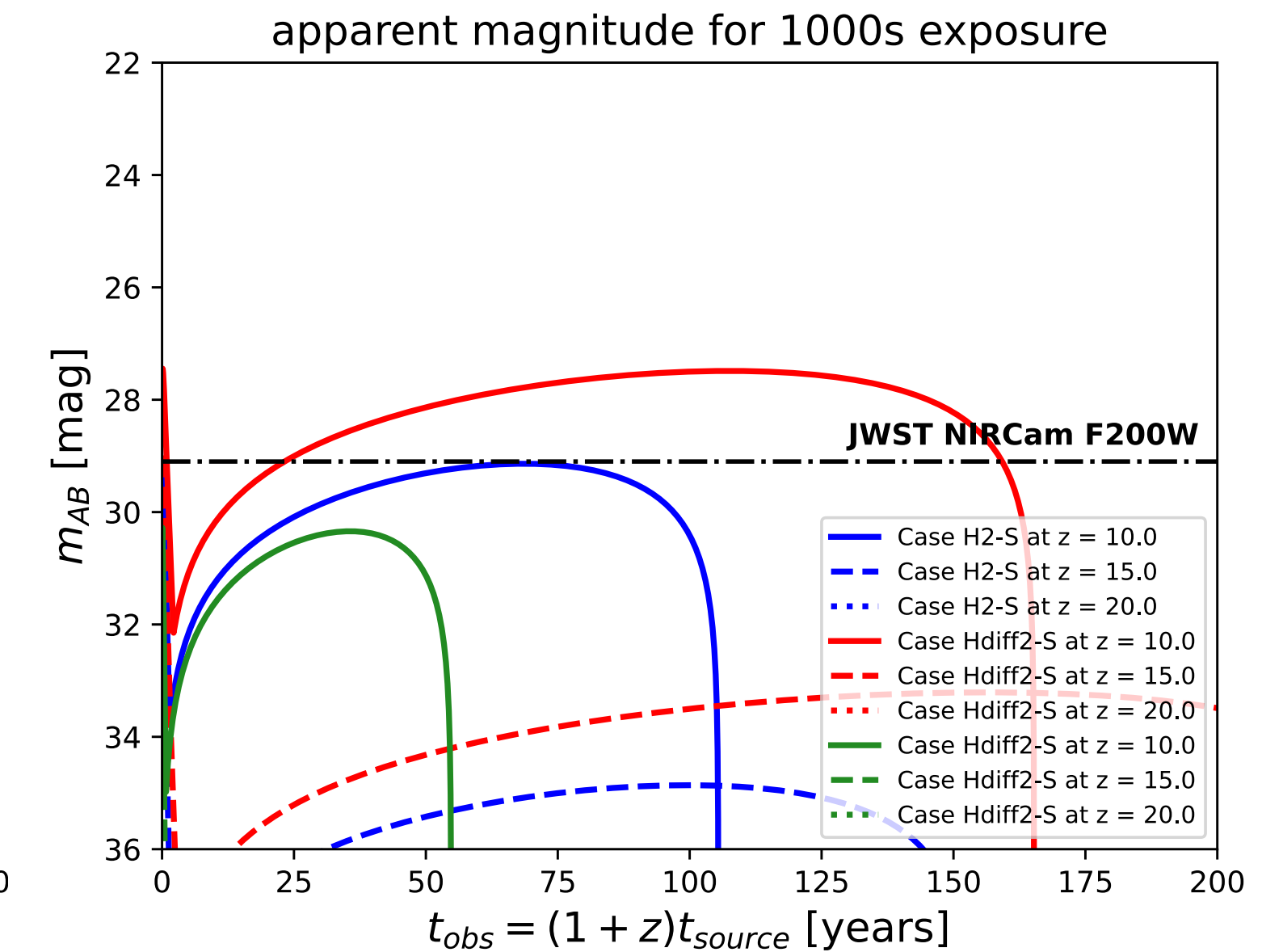
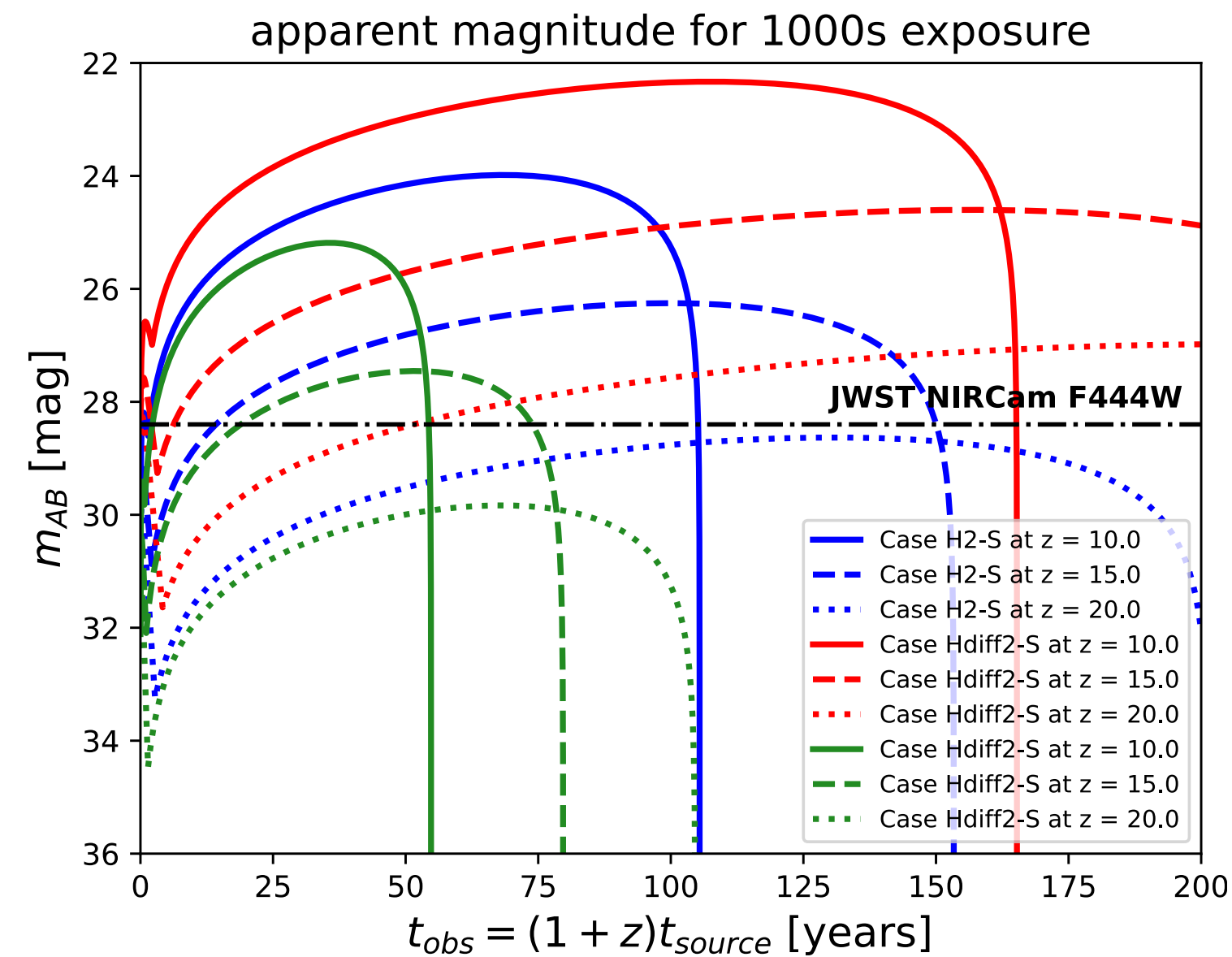


Discussion: Possible outcome of the explosion

Jockel, SF+ in prep.



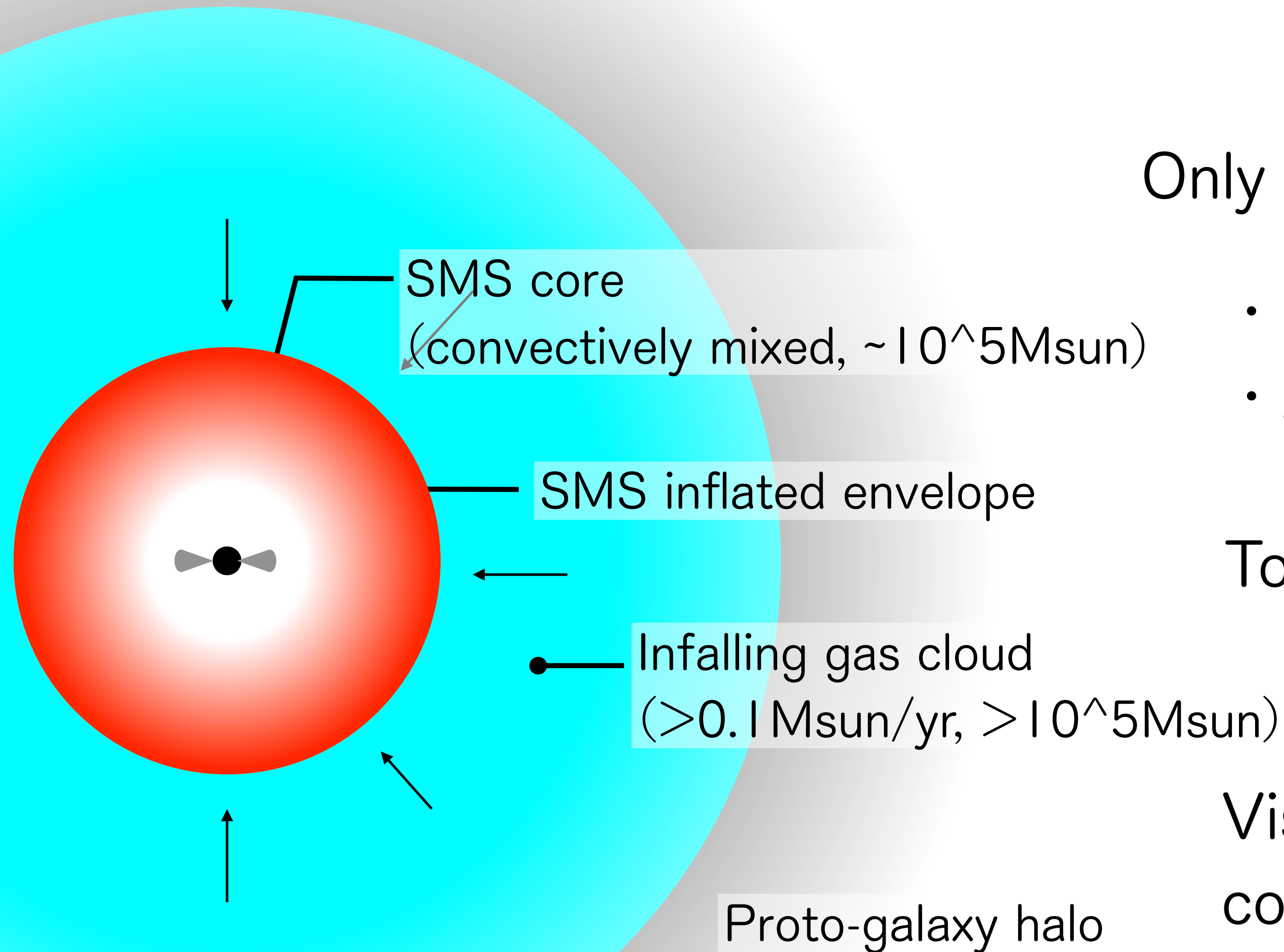
→ source time (d)



→ obs. time (yr)

Long duration (10-100 yrs), red object.

Discussion: Realistic environment



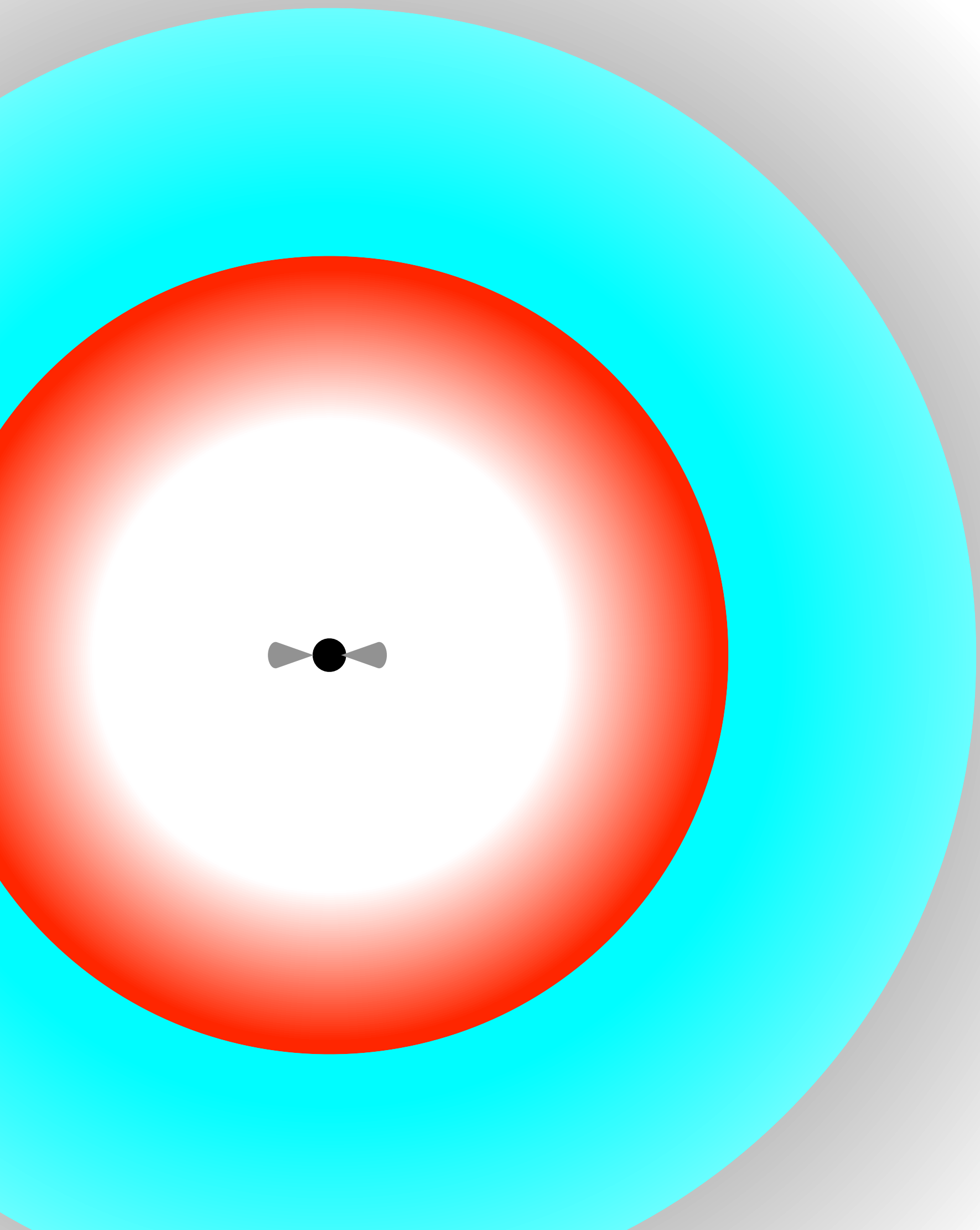
Only the collapse of SMS core is simulated.

- Envelope
- Atomic cooling cloud \sim SMS mass

Total ejecta mass $\sim 10^5 M_{\text{sun}}$

Viscosity-driven ejecta does not contribute much to total ejecta property

Discussion: Realistic environment



The ejecta have to sweep inflated envelope.

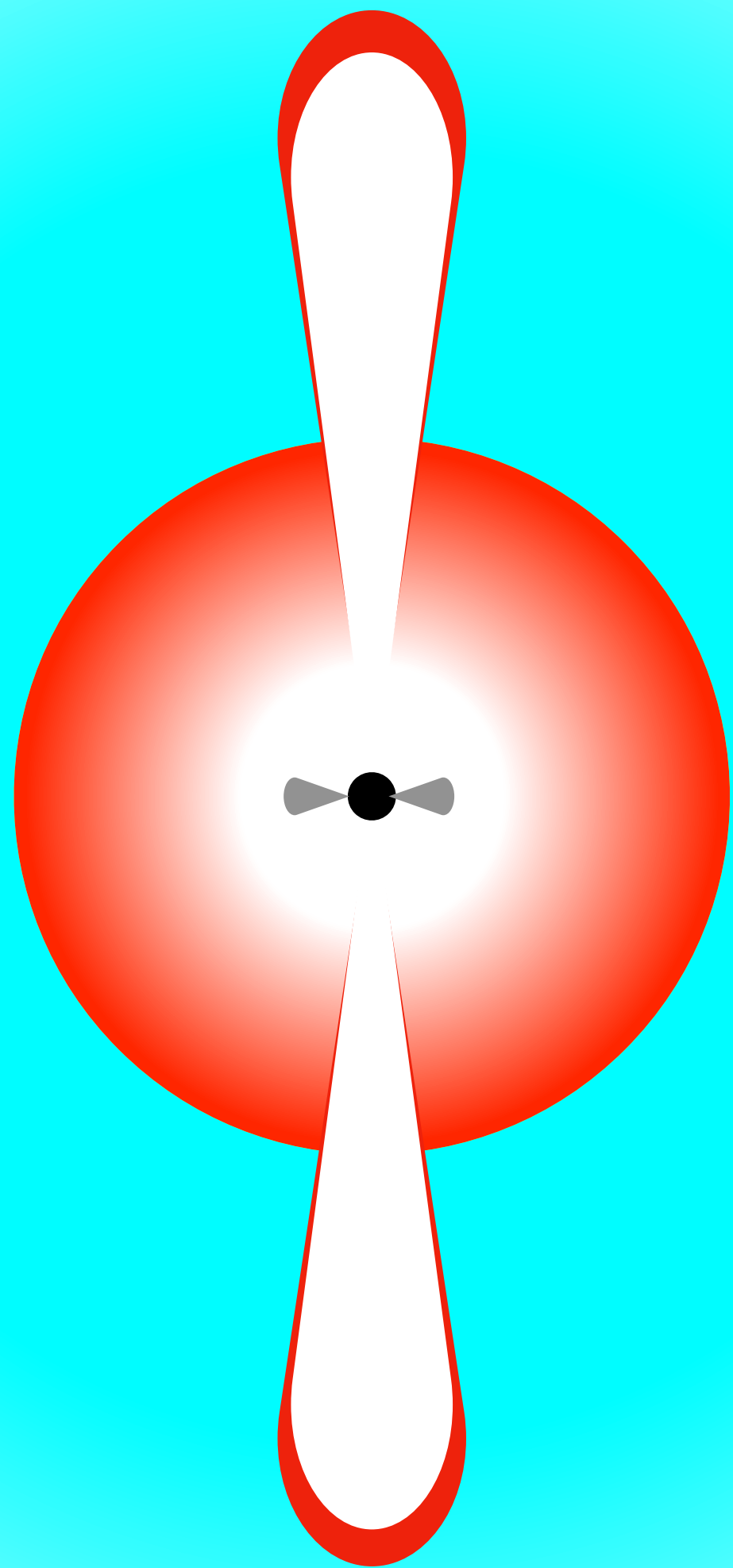
...breaks out the SMS surface (photosphere)

The ejecta then sweeps up the infalling gas cloud, with photon radiation. photosphere may locate at the shock front.

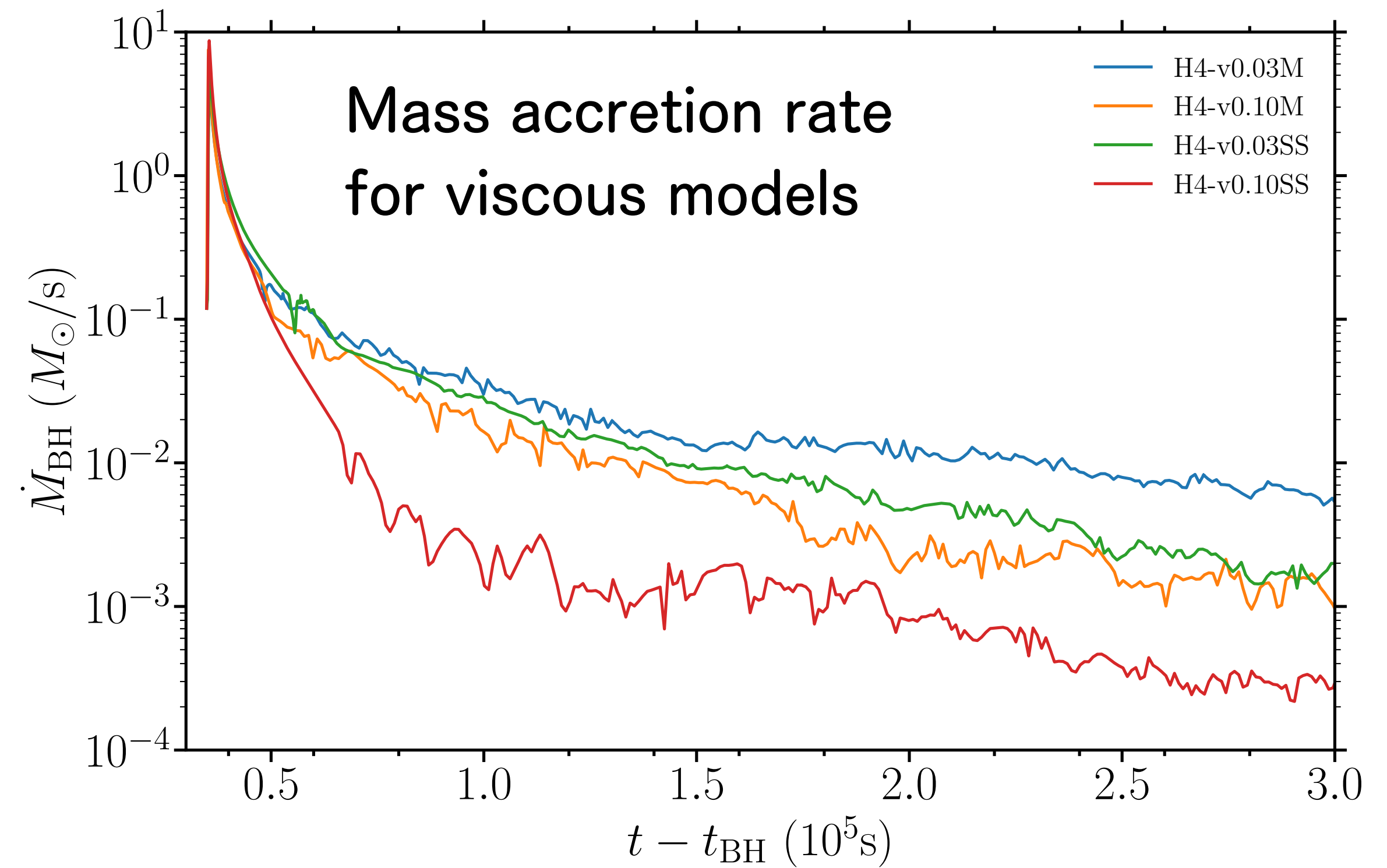
~ CSM-interacting supernova

(with $E \sim 10^{55} - 10^{56}$ erg, $M_{\text{CSM}} \sim 10^5 M_{\text{sun}}$)

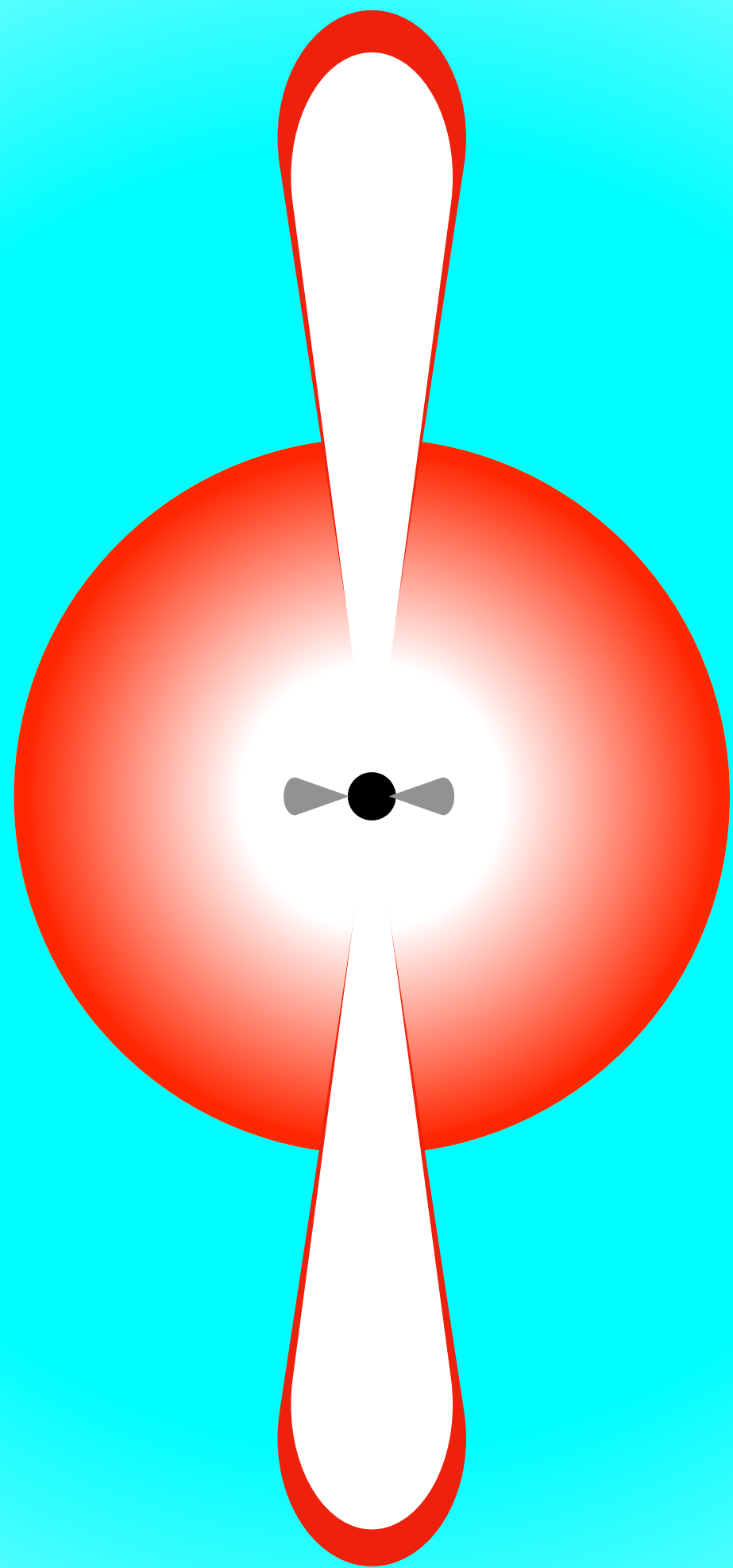
Discussion: Jet driven by BH-disk



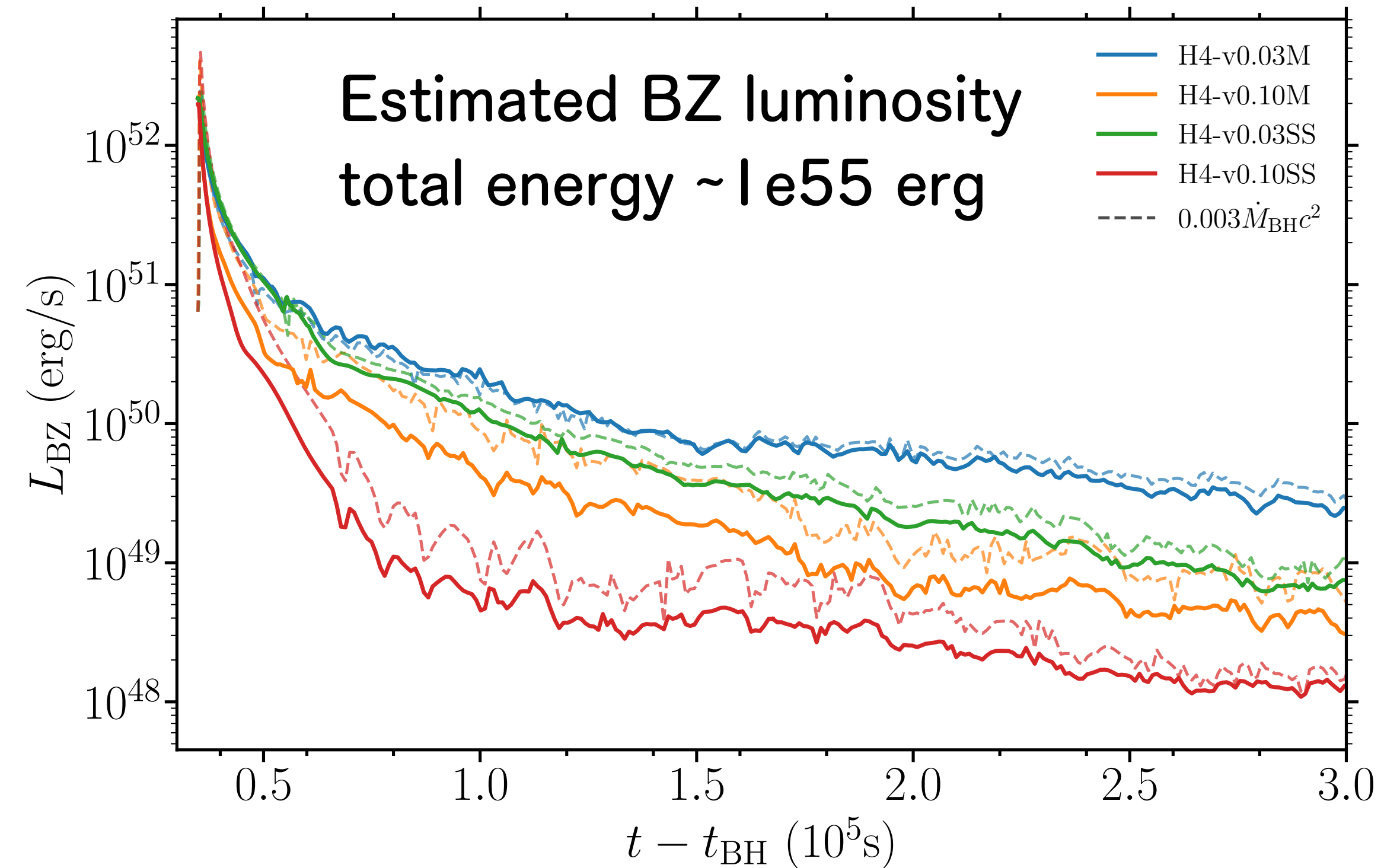
Disk accretion onto BH could drive relativistic jet cf. Matsumoto+15,16



Discussion: Jet driven by BH-disk

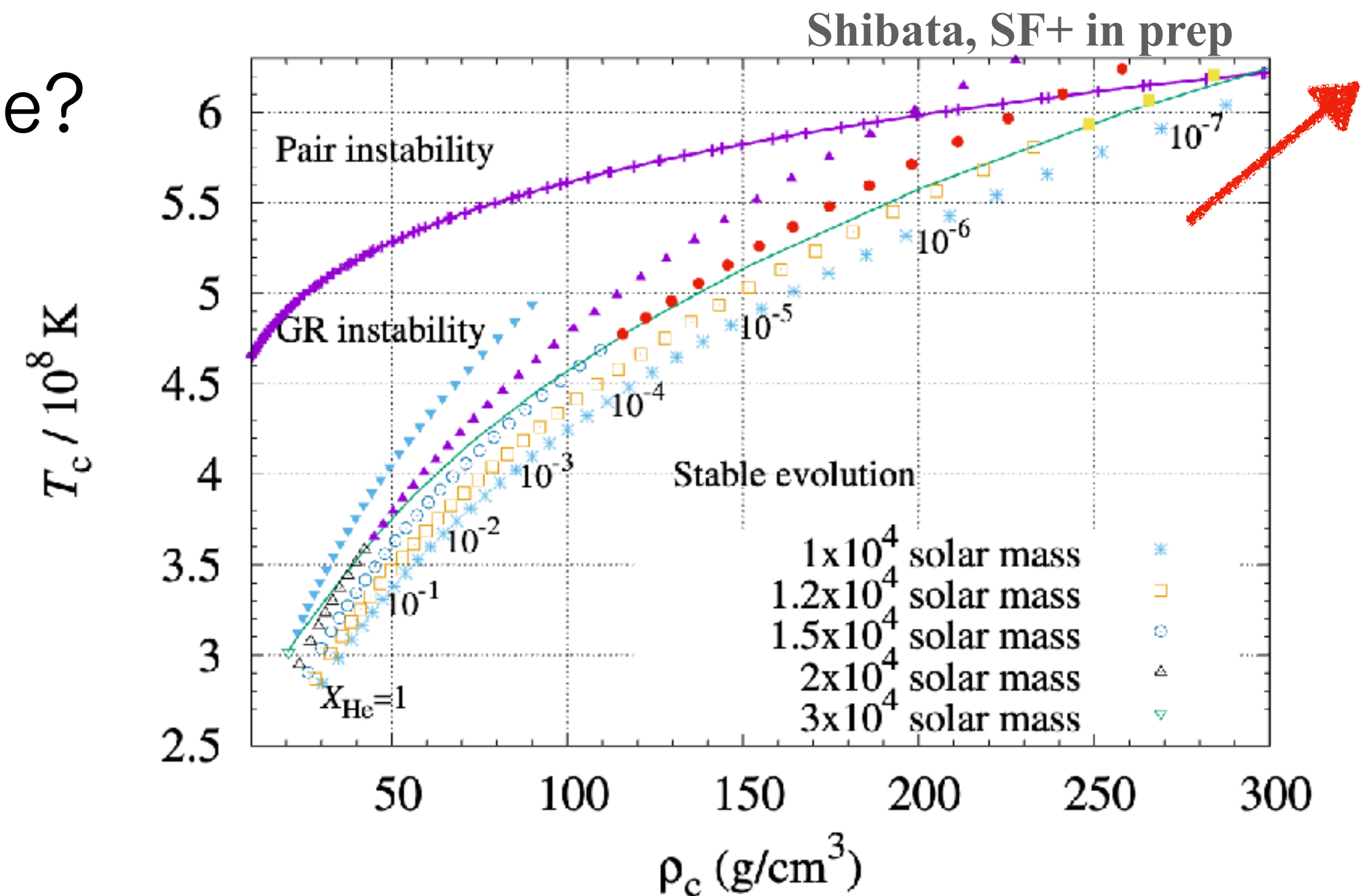


Disk accretion onto BH could drive relativistic jet cf. Matsumoto+15,16



Prospects: Lower-mass SMS

- Lower-mass SMS: higher density, temperature at the collapse
- They may be more subject to pair-instability (thermal creation of e^-e^+ pairs)
- Nuclear burning and neutrino energy loss become more significant
- Effect of rotation on pair-unstable collapse?



Can SMS fast rotate?

Balance of radiation pressure, gravitational and centrifugal force:

$$\frac{\kappa}{c} \frac{L}{4\pi R^2} = \frac{GM}{R^2} - R\Omega_{\text{crit}}^2$$

$$\frac{GM}{R^2}(1 - \Gamma_{\text{Edd}}) = R\Omega_{\text{crit}}^2, \text{ where } \Gamma_{\text{Edd}} = \frac{L}{L_{\text{Edd}}}$$

Allowed rotation is limited for $L \sim L_{\text{Edd}}$

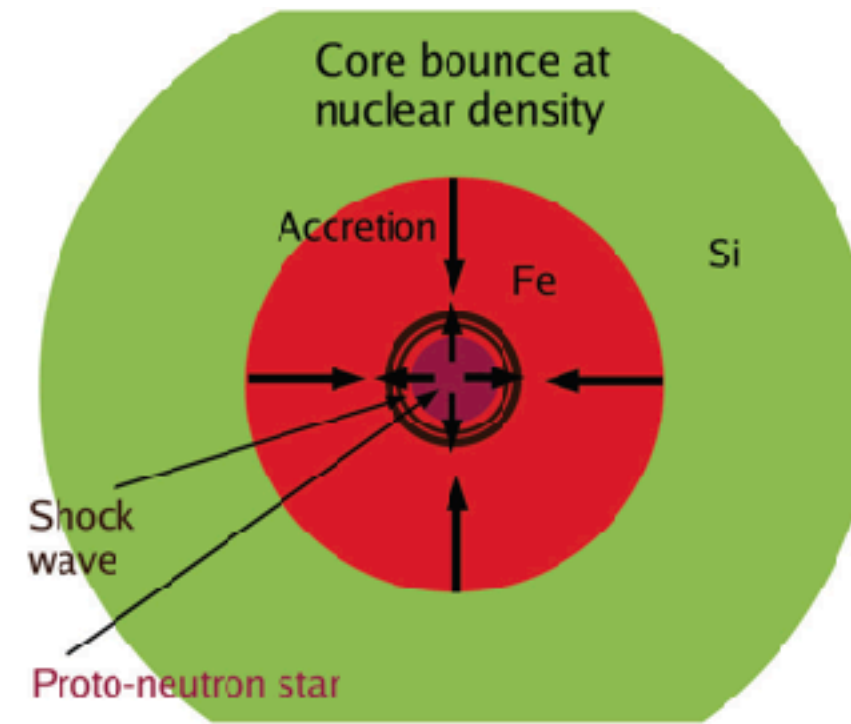
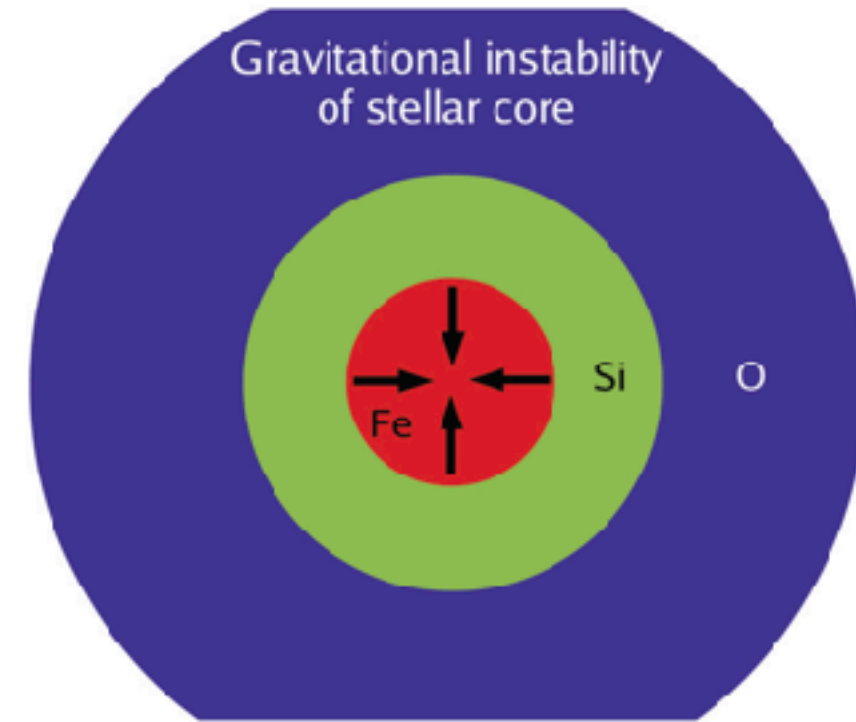
$$\Omega_{\text{crit}} = \sqrt{\frac{GM}{R^3}} \sqrt{1 - \Gamma_{\text{Edd}}}$$

“SMSs are slow rotators” Haemmerlé+2018

Core-collapse supernovae: neutrino-driven scenario

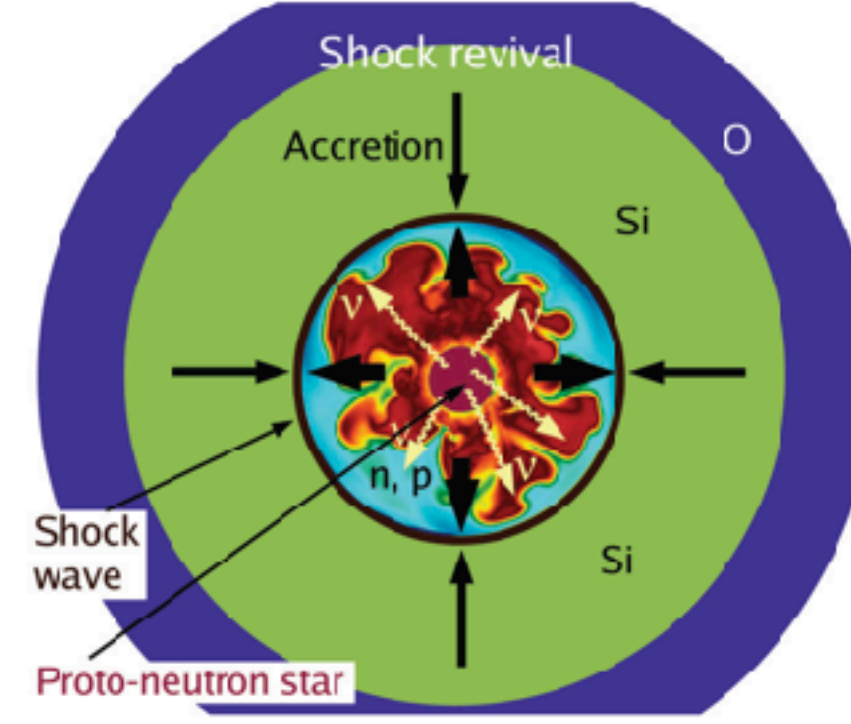
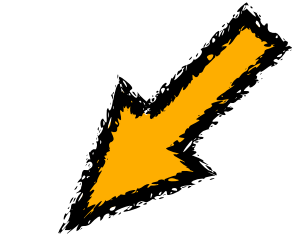
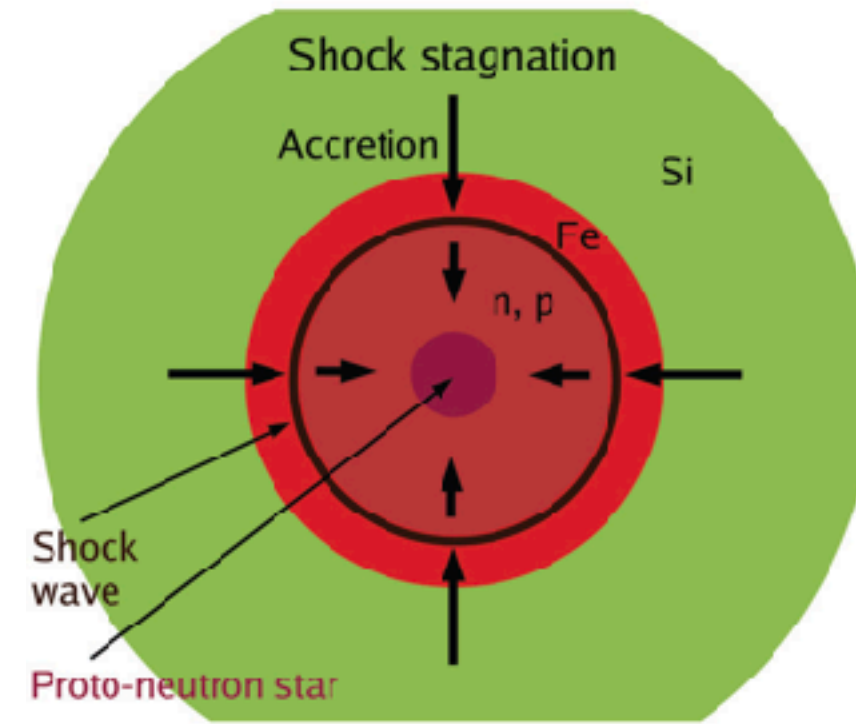
Stars with $\gtrsim 10M_{\odot}$

Collapse of iron core



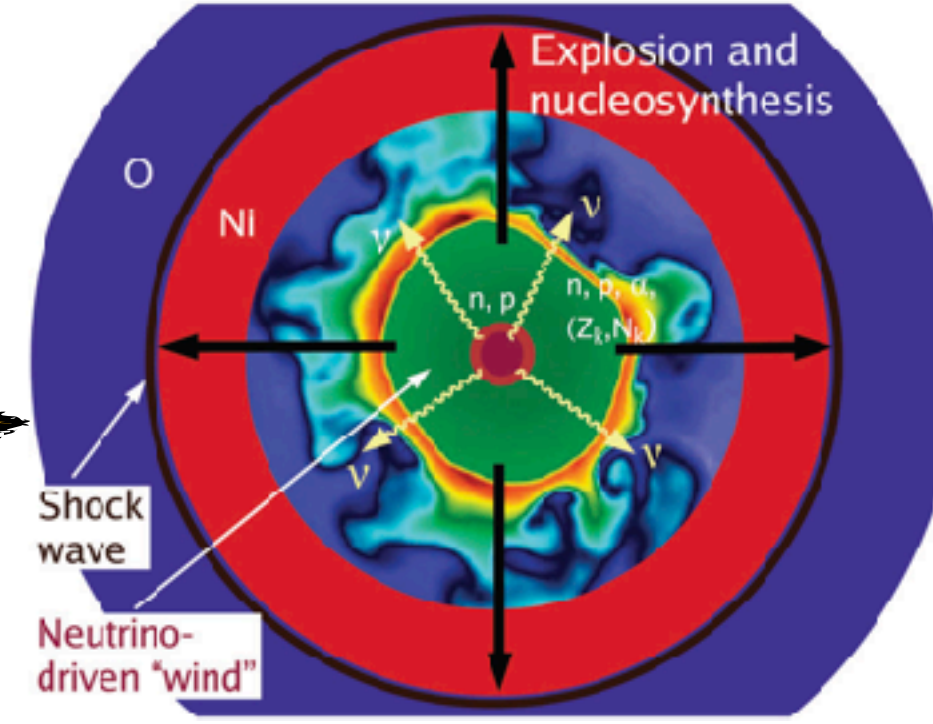
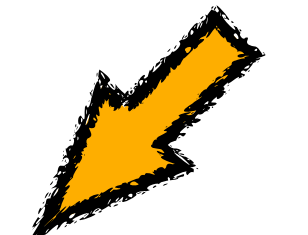
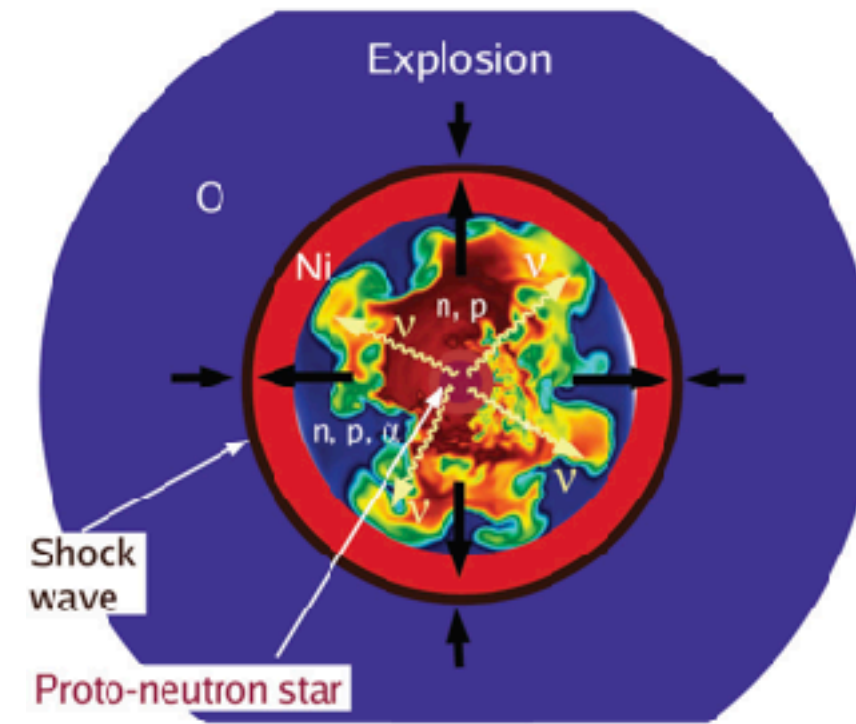
Proto-neutron star formation

Shock developed
→ stall



Shock re-energized
by neutrino heating

Successful breakout
→ observed as SN



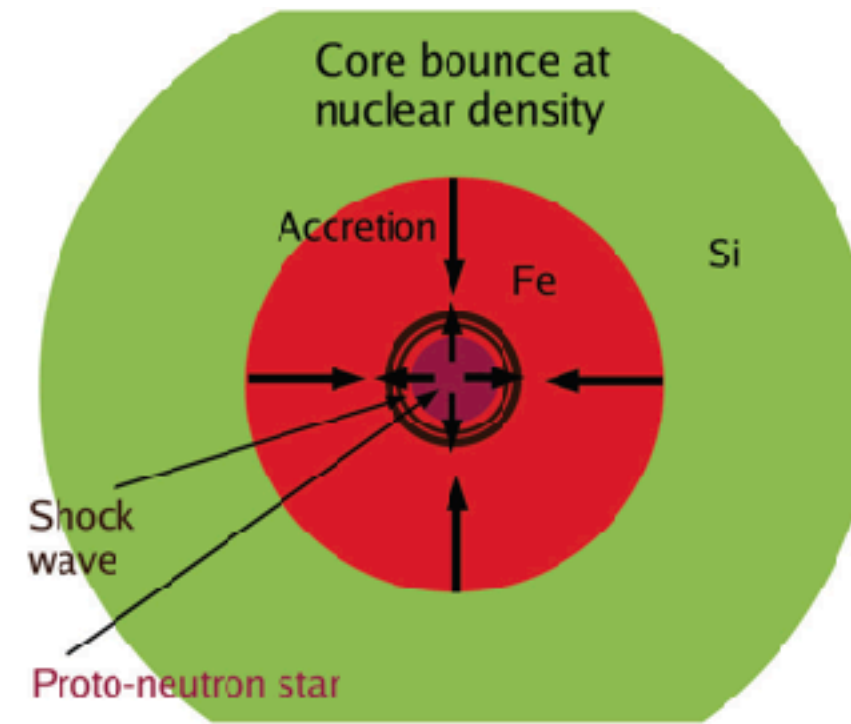
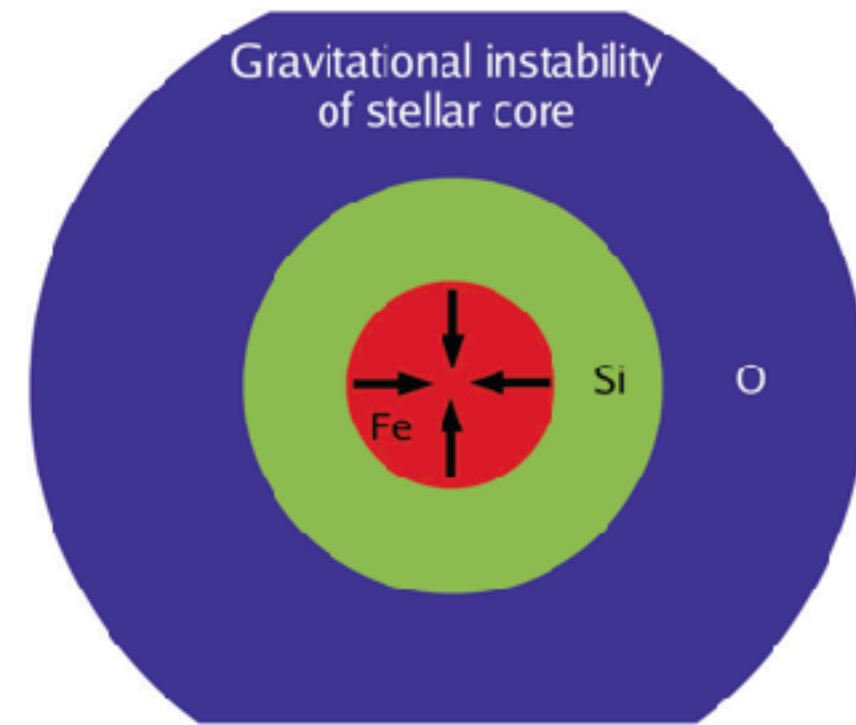
Figures taken from Janka (2012)
(layers not drawn to scale)

BH formation

e.g. O'Connor & Ott 2011

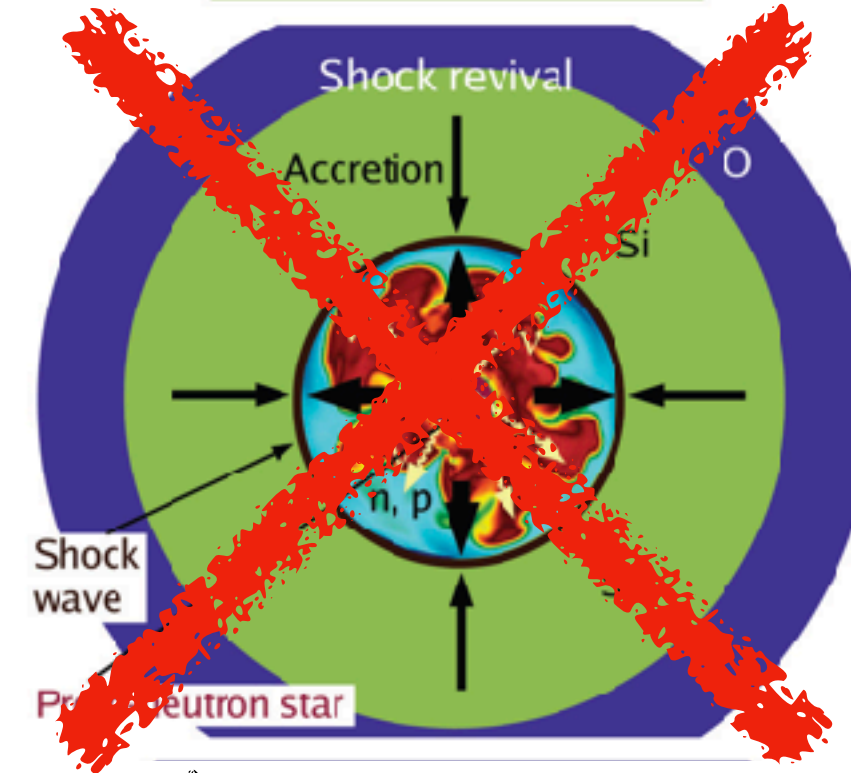
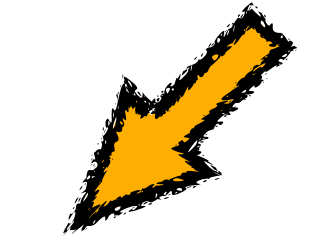
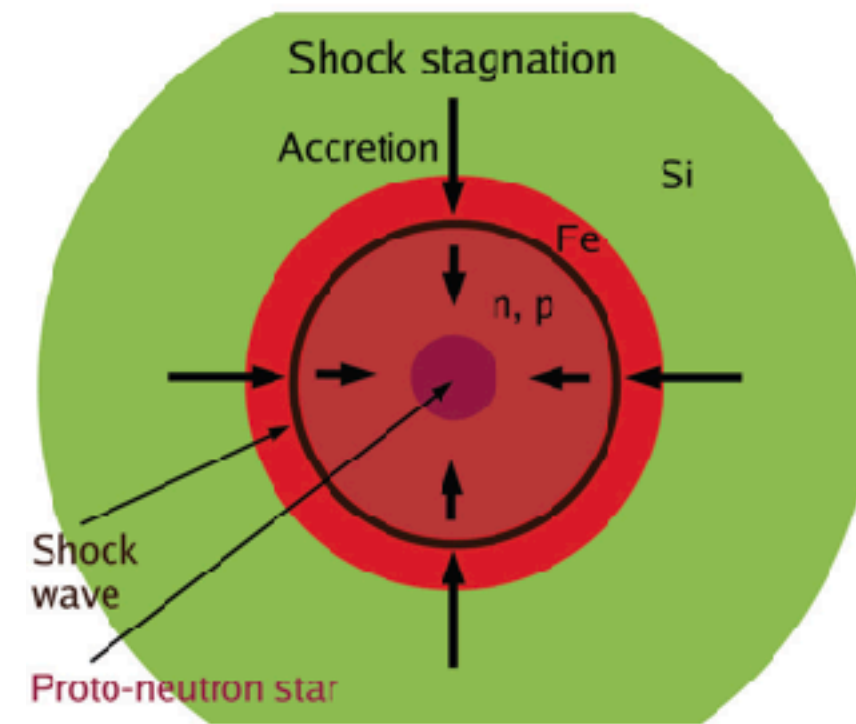
e.g., those with
too compact cores

Collapse of iron core



Proto-neutron star
formation

Shock developed
→ stall

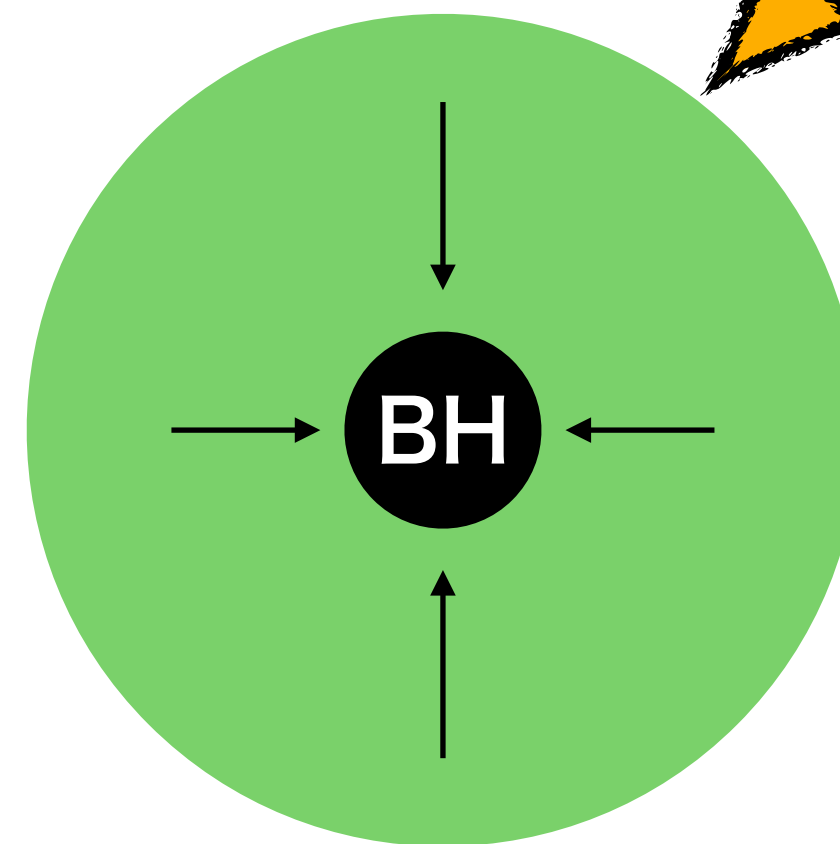


Fail to re-energize
the shock

*B-field can also help explosion

e.g. Obergaulinger & Aloy 2021

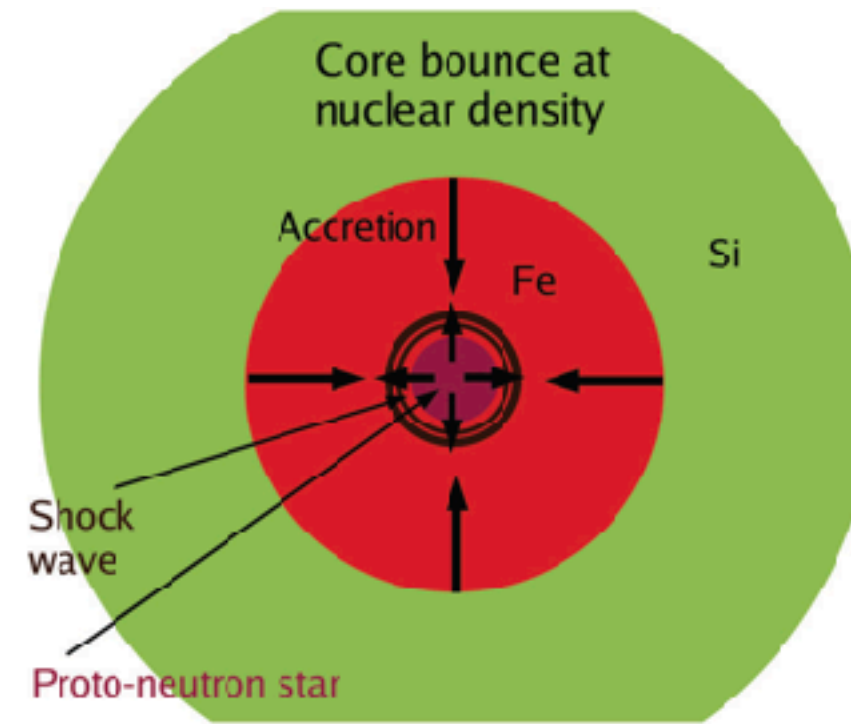
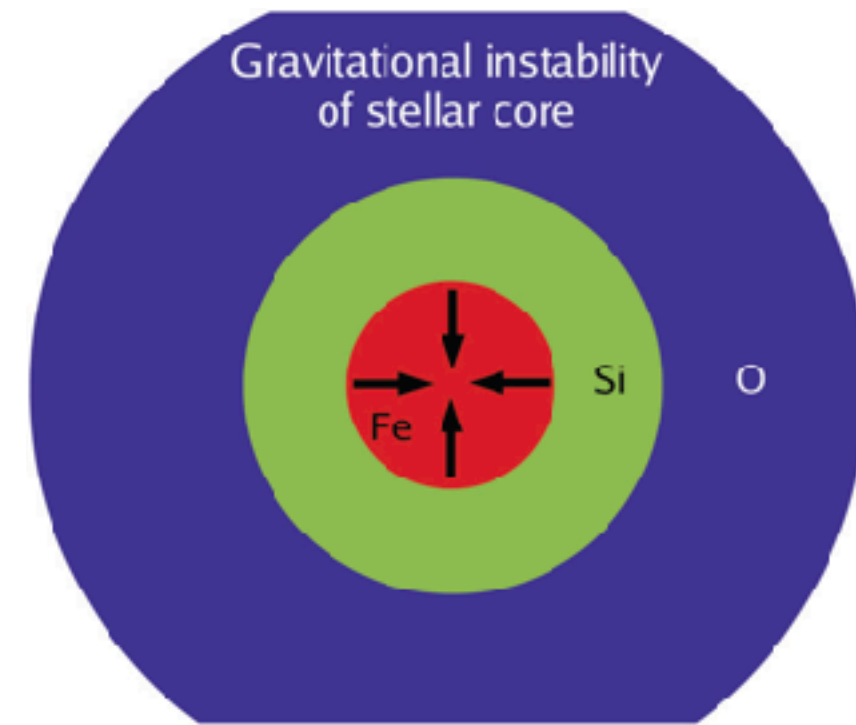
Black hole formation
No transient?



Collapse of rotating star

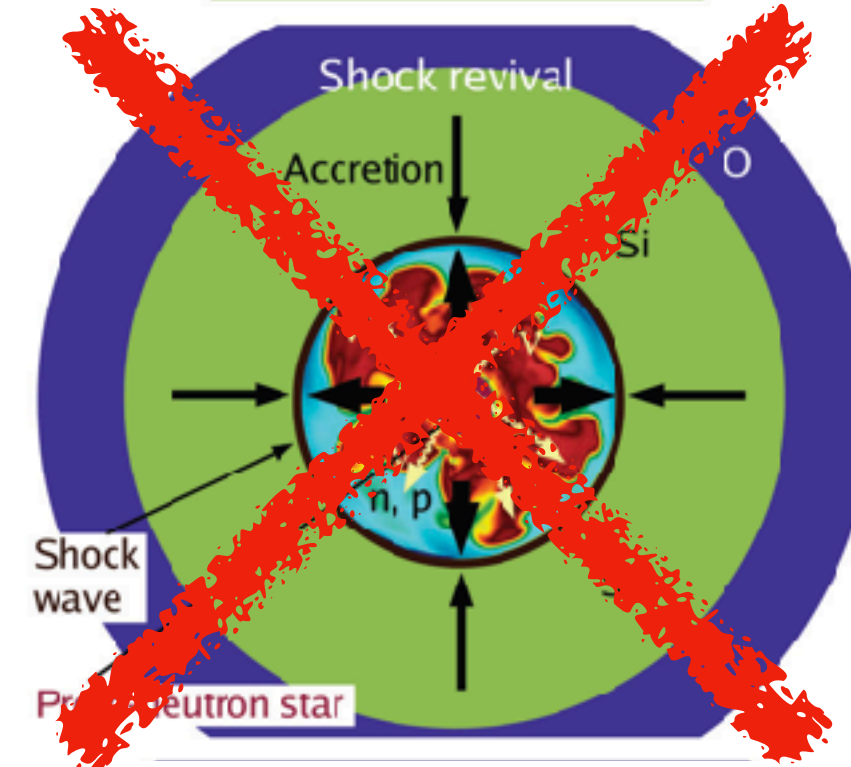
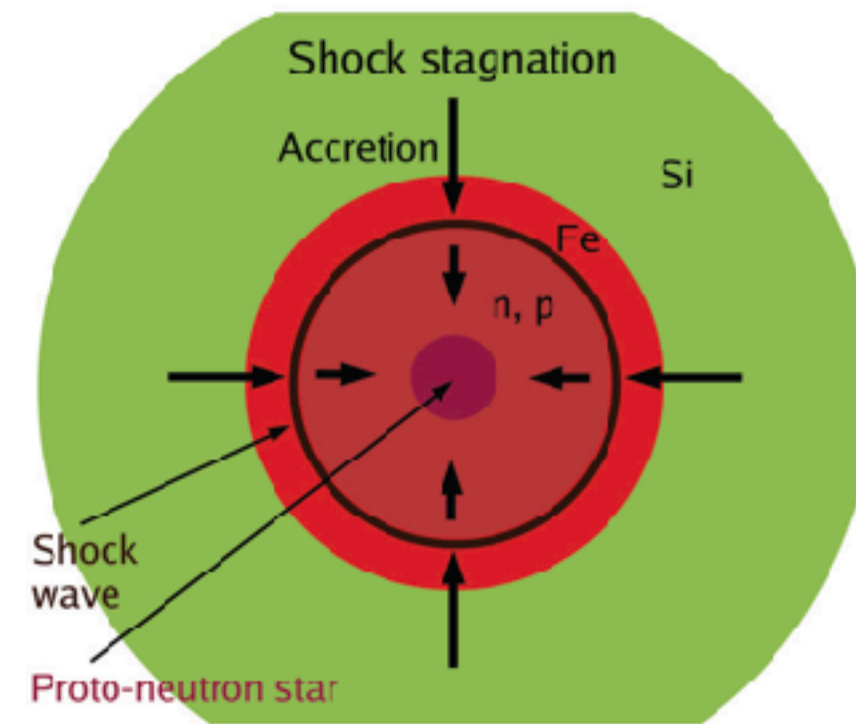
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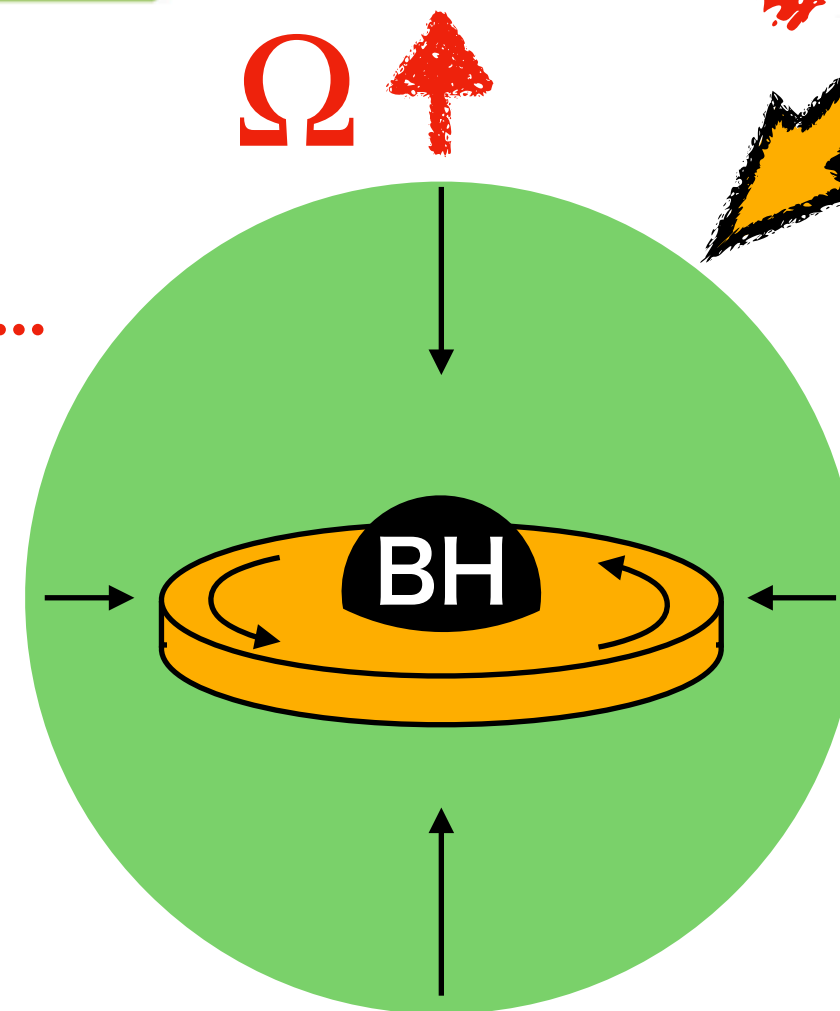
Shock developed → stall



Fail to re-energize the shock

If the star has sufficient rotation...

Accretion disk formed around the BH.
Further activities possible.



*B-field can also help explosion
e.g. Obergaulinger & Aloy 2021

BH-disk activities

- Gamma-ray bursts (GRBs)

BH-disk is one of the promising central engines

(e.g., Woosley et al. 1993...)

Energy liberated by viscous accretion:

$$GM_{\text{BH}}M_{\text{disk}}/r_{\text{disk}} \approx 3 \times 10^{52} \text{ erg} \left(\frac{M_{\text{BH}}}{10M_{\odot}} \right) \left(\frac{M_{\text{disk}}}{0.1M_{\odot}} \right) \left(\frac{r_{\text{disk}}}{10^7 \text{ cm}} \right)^{-1}$$

It can drive an energetic outflow from disk.

(MacFadyen & Woosley 1999)

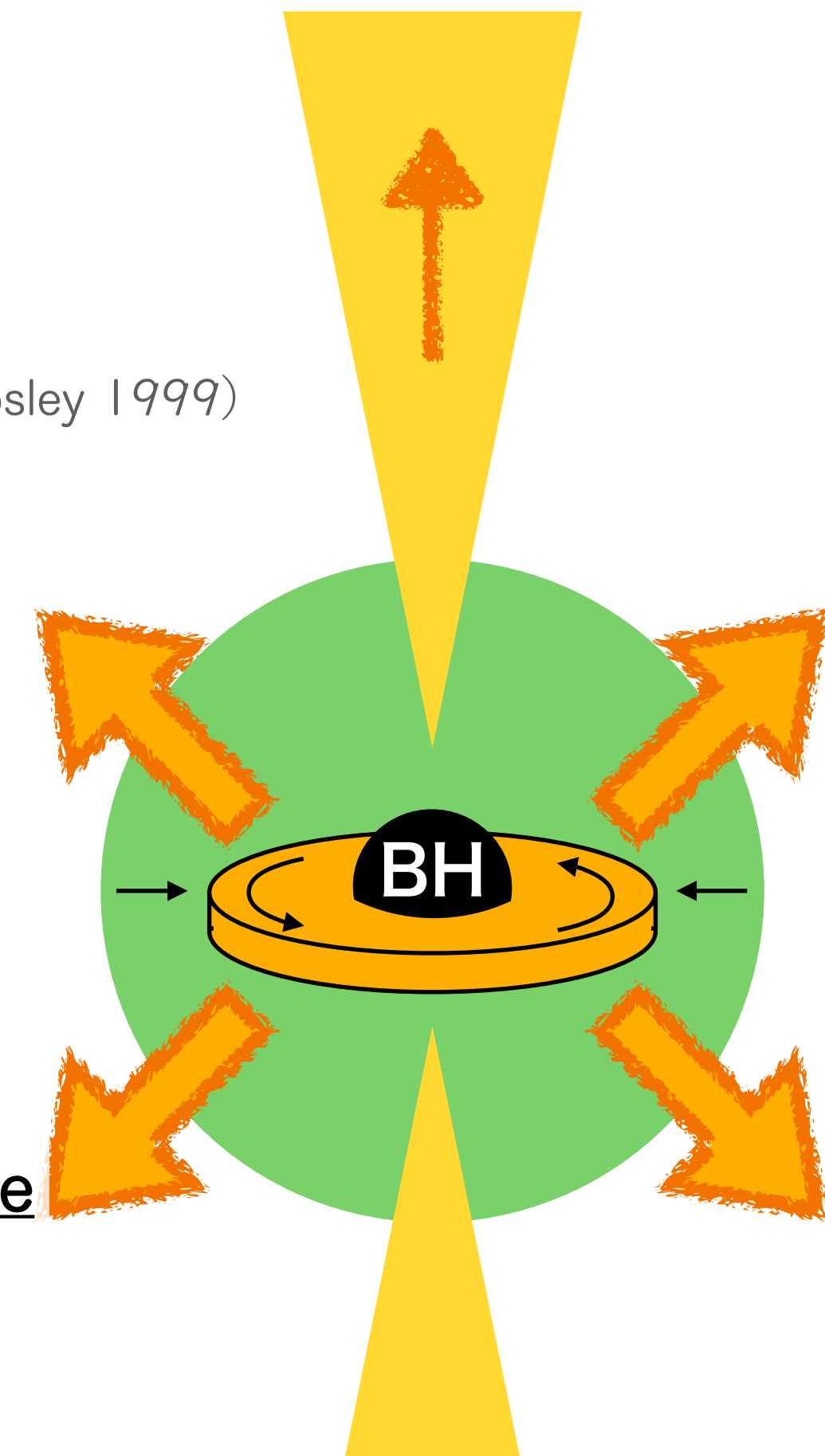
(at least some) long GRBs are accompanied by energetic SNe

(lc-BL)

- Explosion energy $E_{\text{K}} = (0.8 - 4.4) \times 10^{52} \text{ erg}$

- ^{56}Ni mass $M_{\text{Ni}} = (0.2 - 0.5) M_{\odot}$ (Cano et al. 17)

Viscosity-driven outflow from disk would naturally explain such SNe



Neutrino cooling vs viscous heating

For high temperatures, ($kT \gtrsim 1 \text{ MeV}$) neutrino emission cools down the disk.

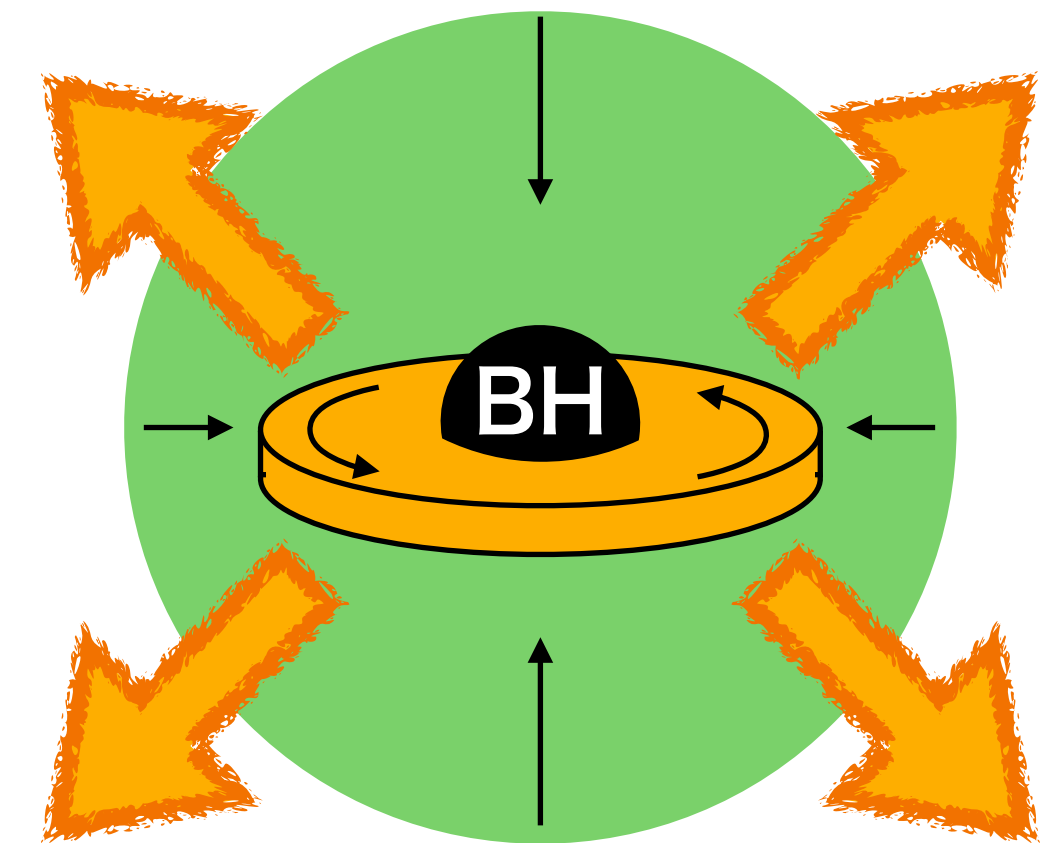
Neutrino emission timescale : $t_{\text{weak}} \sim 1 \text{ s} \left(\frac{kT}{1 \text{ MeV}} \right)^{-5}$

$t_{\text{weak}} \lesssim t_{\text{vis}}$ (NDAF) phase: weak/no outflow

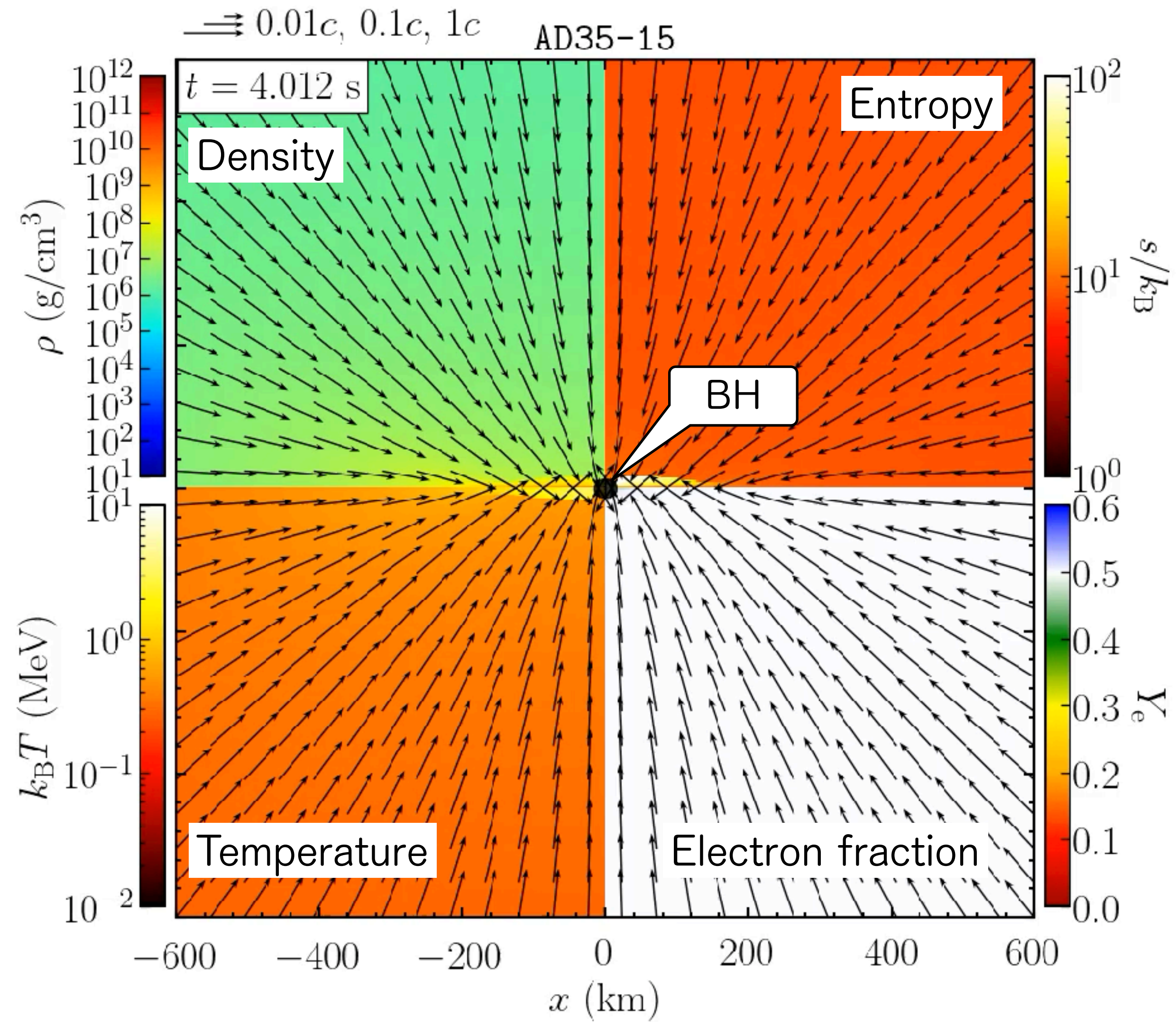
$t_{\text{weak}} \gg t_{\text{vis}}$ phase: strong outflow would be driven

I performed simulations of the post-merger remnant with

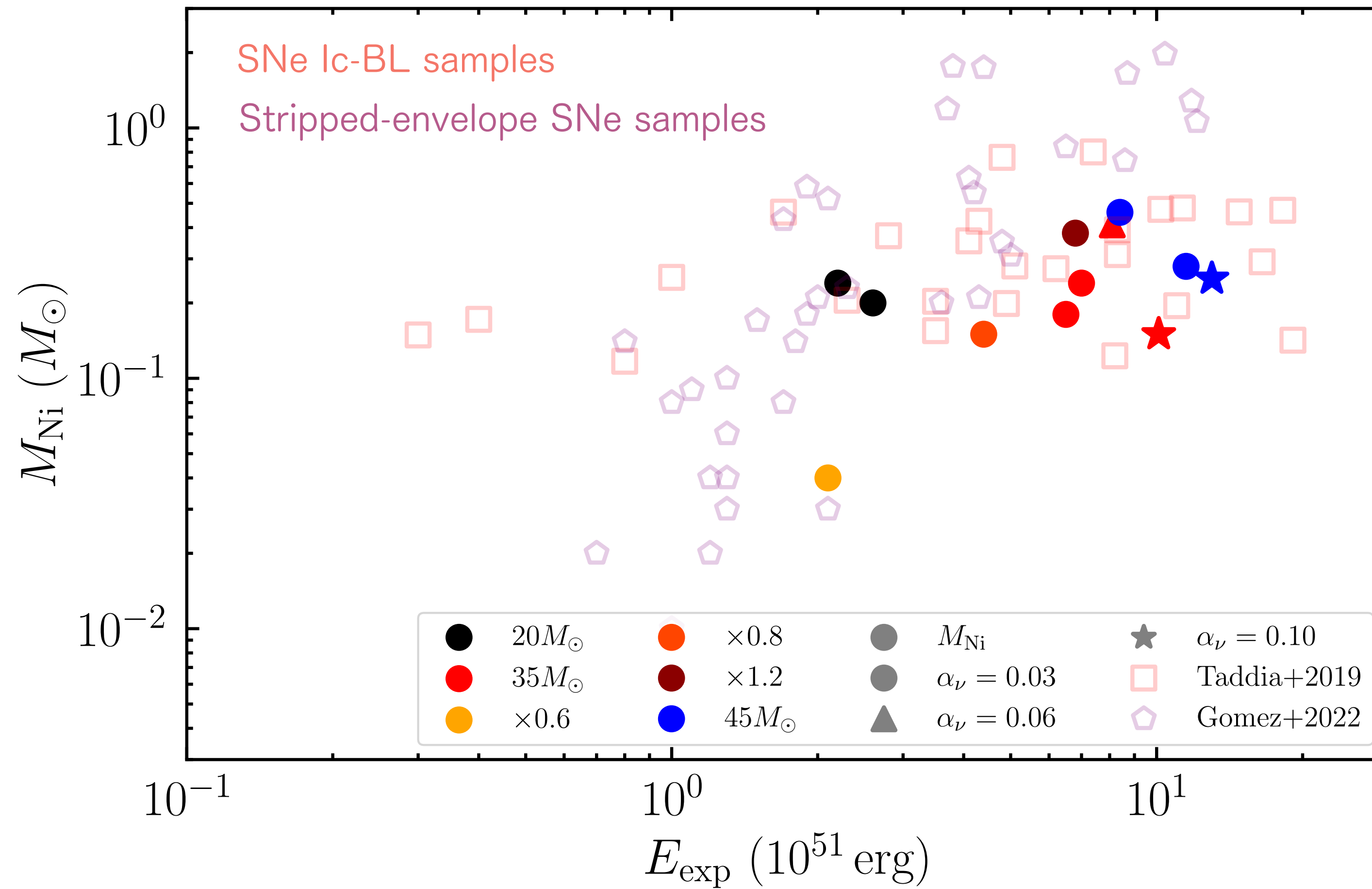
- GR gravity
- Neutrino cooling
- Viscous angular momentum transport



Disk formation \rightarrow NDAF \rightarrow Outflow



Comparison with observations

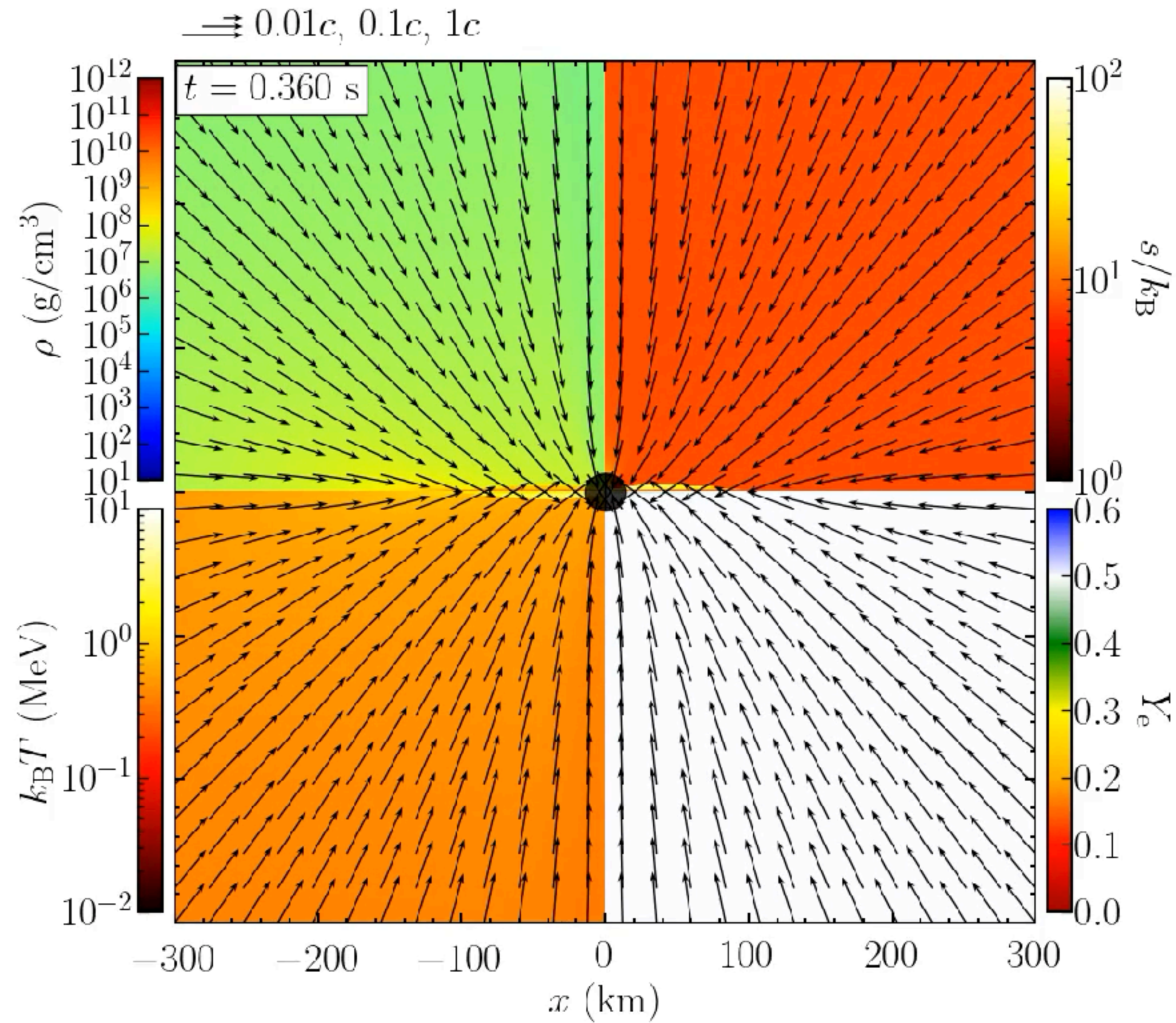


Nucleosynthesis calculation in the ejecta $\rightarrow M_{\text{Ni}} \gtrsim 0.1M_{\odot}$

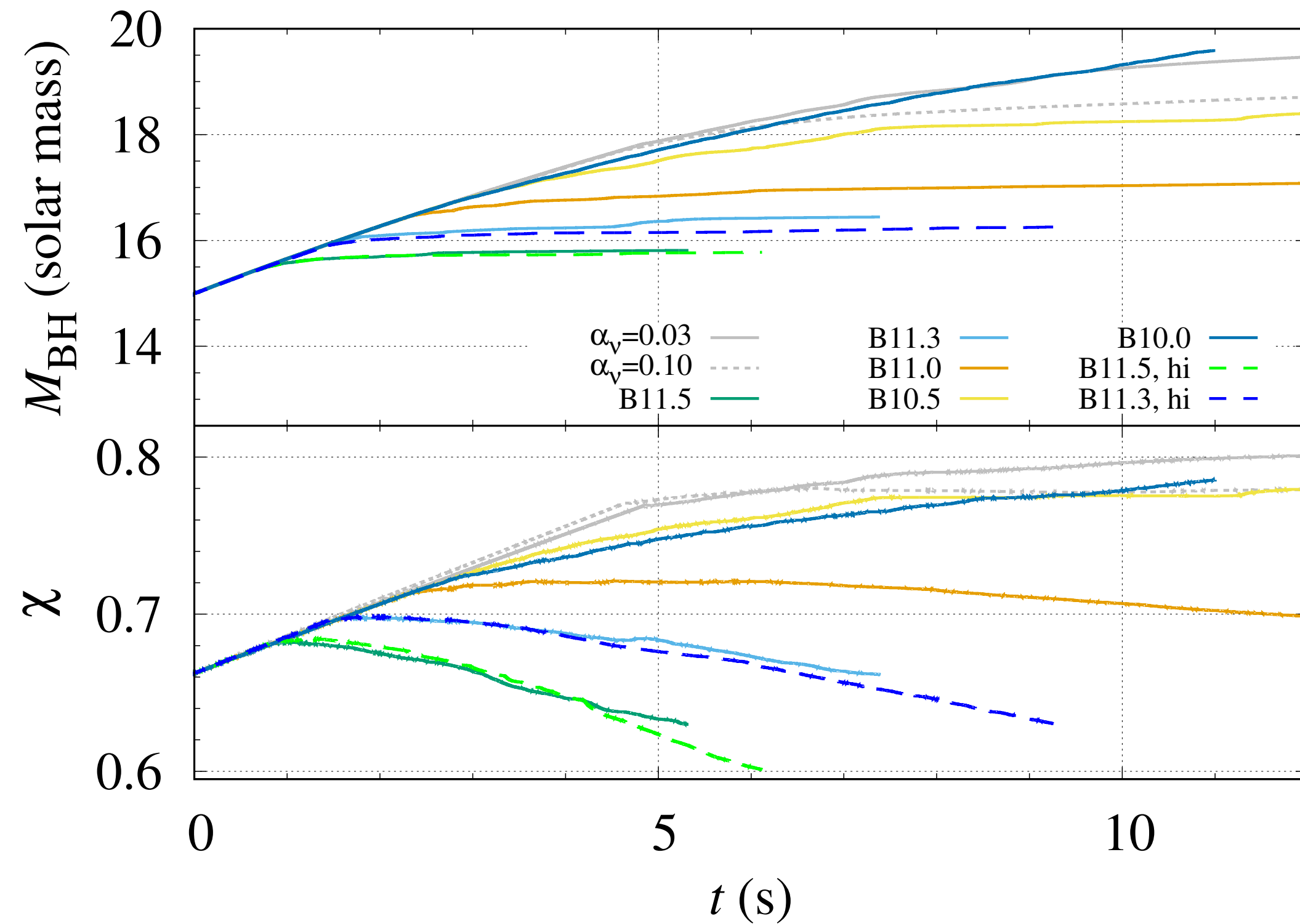
BH-disk can be an engine of the energetic explosion.

MHD models for GRB jets

Shibata, SF+24



MHD models for GRB jets



Spin down of BH by the feedback of energy extraction is numerically observed(!)

Summary

Numerical-relativity simulations of the collapses of rotating supermassive stars

Bounce-shock-induced ejecta with mass up to 1% of initial star, $v \sim 0.2c$

Kinetic energy $\sim 10^{55} - 10^{56}$ erg

(Mass is likely dominated by the swept-up cloud surrounding the star)

Bounce-shock-induced ejecta with mass up to 1% of initial star, $v \sim 0.2c$

Numerical relativity simulations of collapses of rotating massive stars:

- It can explode with $E \sim 10^{52}$ erg driven by disk outflow (\leftarrow viscous model)
- It can synthesize sufficient amount ($\gtrsim 0.1M_{\odot}$) of ^{56}Ni (\leftarrow viscous model)
- It can drive a jet (\leftarrow MHD model with an ideal config.)
- No significant r-process in the ejecta