
A discussion on the effects of QGP's electric conductivity on observables in high-energy heavy-ion collisions

Nicholas J. Benoit (PostDoc at Hiroshima Univ.)

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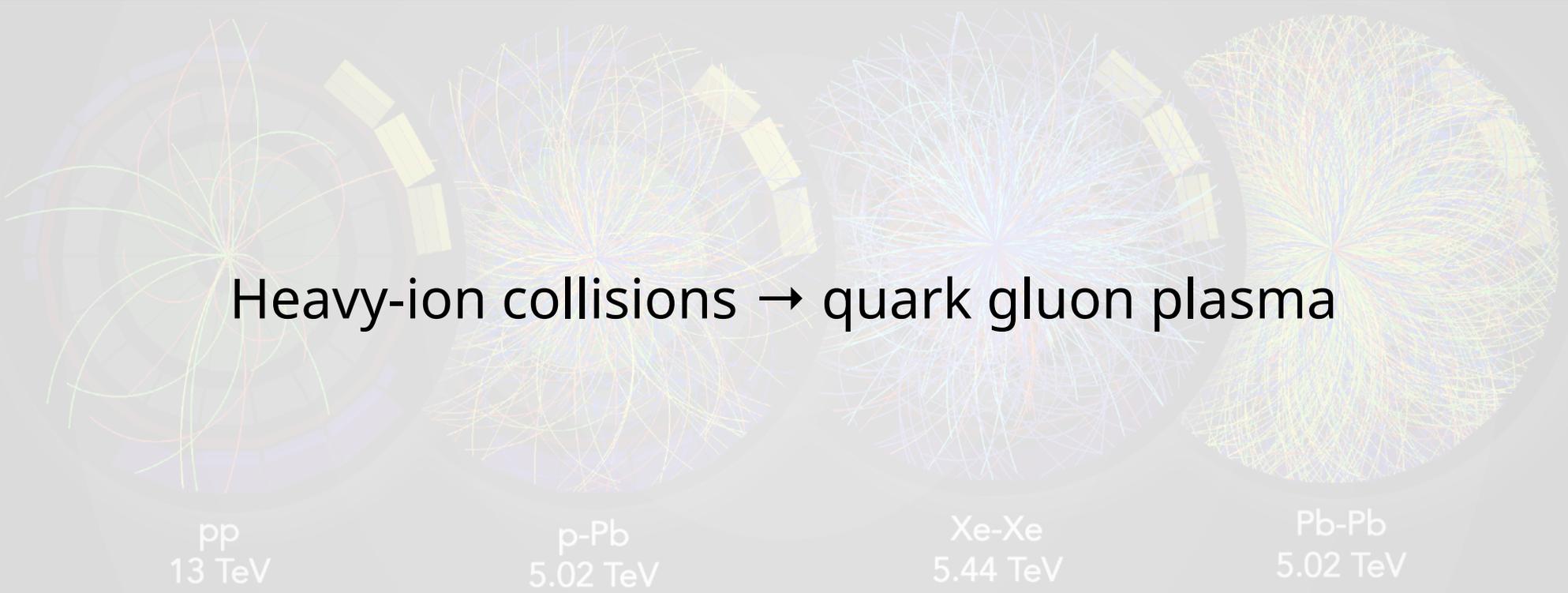
広島大学



Collaborators [paper soon to be released]:

T. Miyoshi (HU), C. Nonaka (HU, SKCM2, KMI), A. Sakai (HU), and H. R. Takahashi (KU)

Heavy-ion collisions → quark gluon plasma



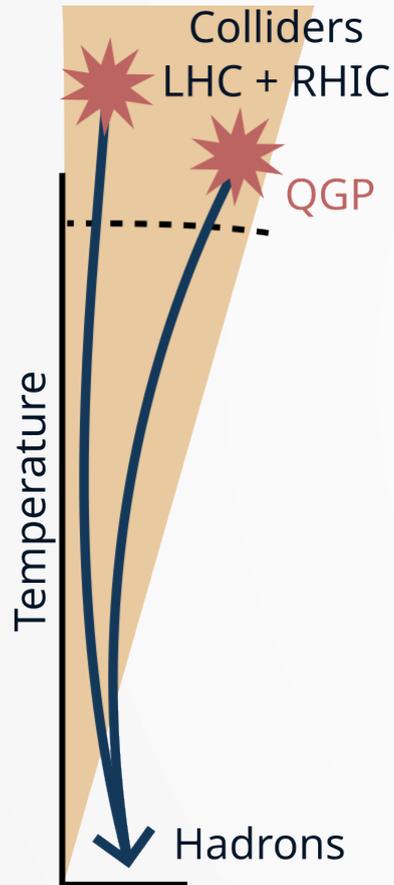
pp
13 TeV

p-Pb
5.02 TeV

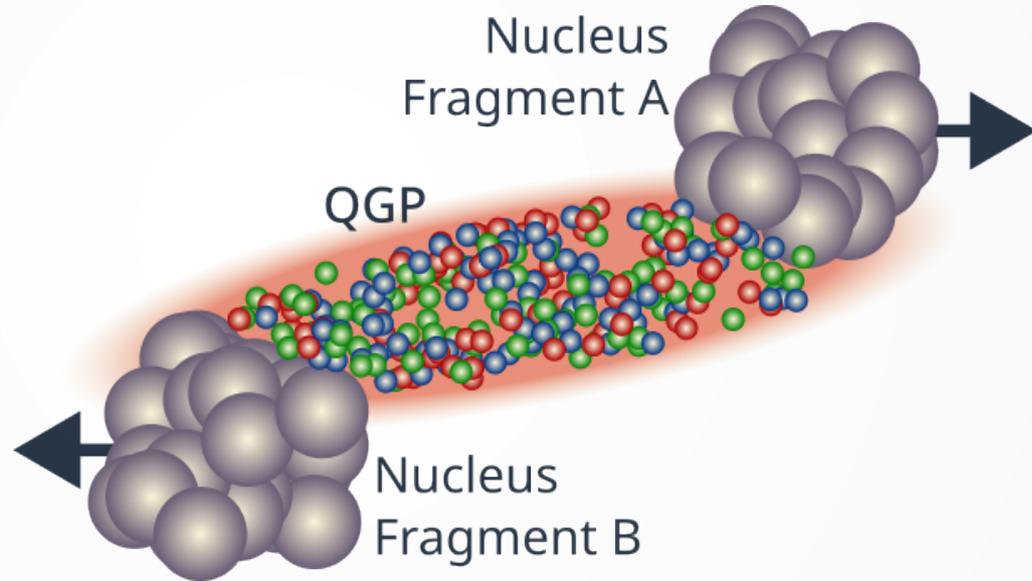
Xe-Xe
5.44 TeV

Pb-Pb
5.02 TeV

How to make quark gluon plasma



- How is quark gluon plasma (QGP) created in heavy-ion collisions?



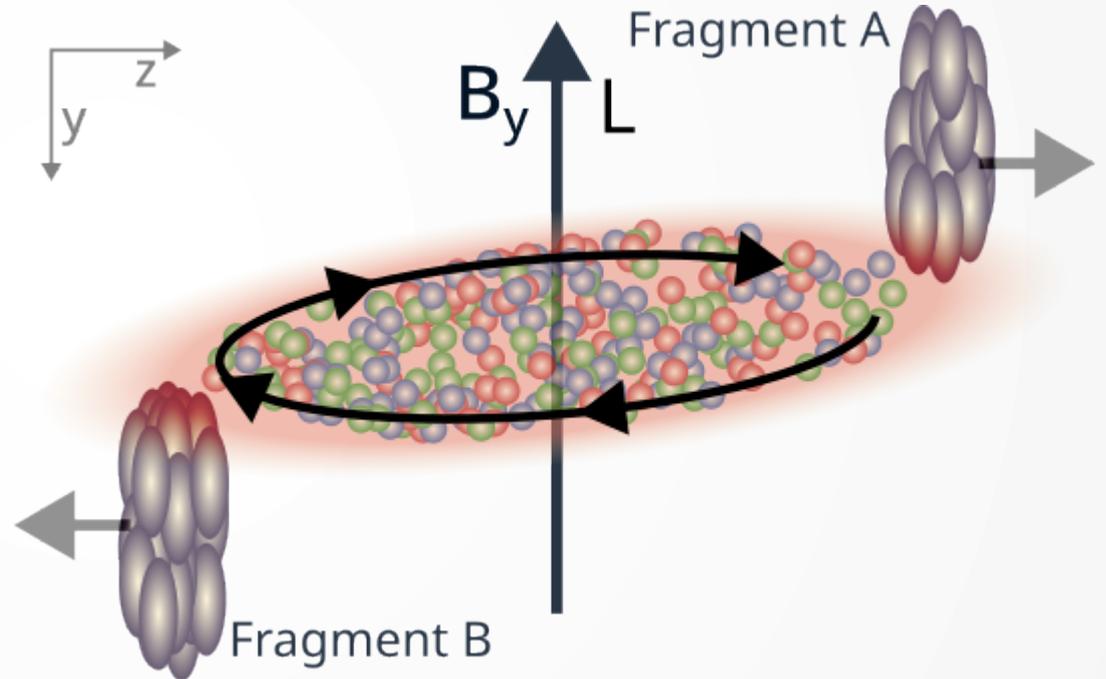
- Collide heavy nuclei like Au, Cu, Xe, Pb, etc. at relativistic energies

Non-central collisions and global polarization

Non-central collisions have many interesting properties,

- large orbital angular momentum
- strong magnetic field
- etc.

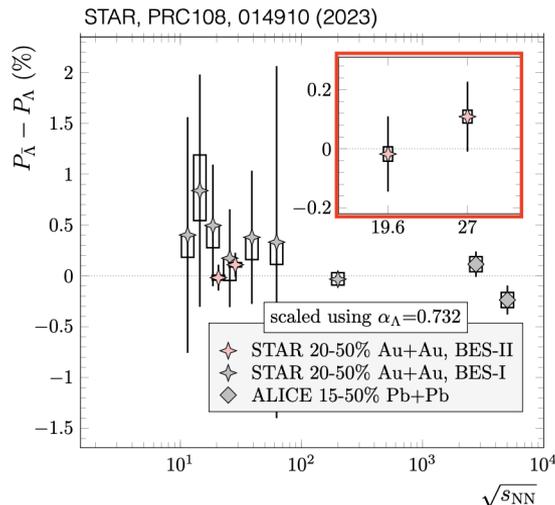
These result in a “global” polarization for particles



Non-central collisions and global polarization

$P_{\bar{\Lambda}} - P_{\Lambda}$ difference?

- B-field lifetime is likely very short (<0.5 fm/c) but could be sustained by QGP electric conductivity (unknown)
- Can we put constraint on the late-stage B-field by the polarization difference?



$$P_{\Lambda} \simeq \frac{1}{2} \frac{\omega}{T} + \frac{\mu_{\Lambda} B}{T}$$

$$P_{\bar{\Lambda}} \simeq \frac{1}{2} \frac{\omega}{T} - \frac{\mu_{\Lambda} B}{T}$$

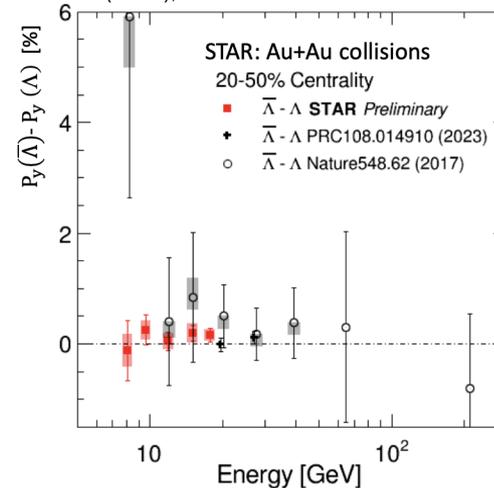
μ_{Λ} : Λ magnetic moment

$$B = (P_{\bar{\Lambda}} - P_{\Lambda})T / (-2\mu_{\Lambda})$$

$$\simeq 30\Delta P_{\Lambda} \times 10^{14} [T]$$

$$eB/m_{\pi}^2 \sim 10\Delta P_{\Lambda} \quad (T=120 \text{ MeV})$$

Q. Hu (STAR), SQM2024

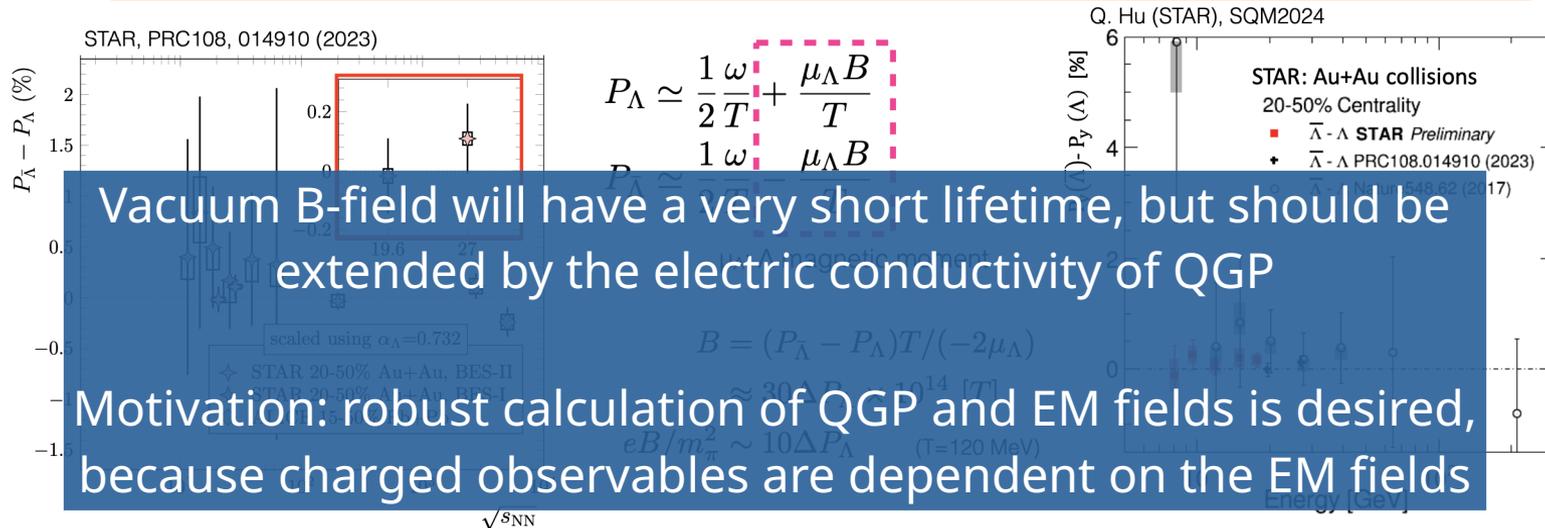


New BES-II results with better precision do not show significant difference.
An upper limit of the late-stage B-field: $B \lesssim 10^{13} T$

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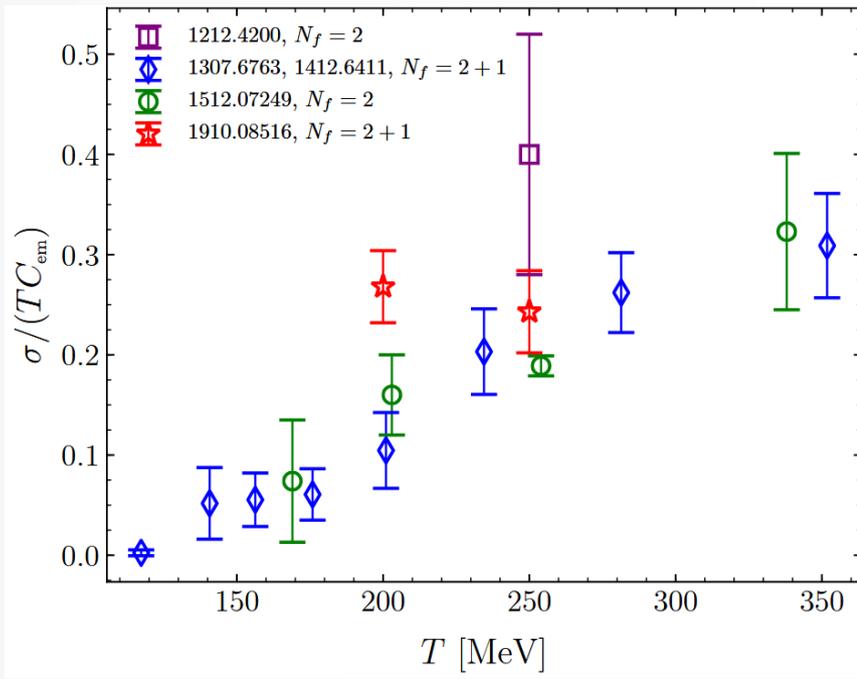


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Quark-gluon plasma (QGP) electric conductivity

- Studied by **lattice calculations (~10 papers)**, pQCD, and kinetic transport theories

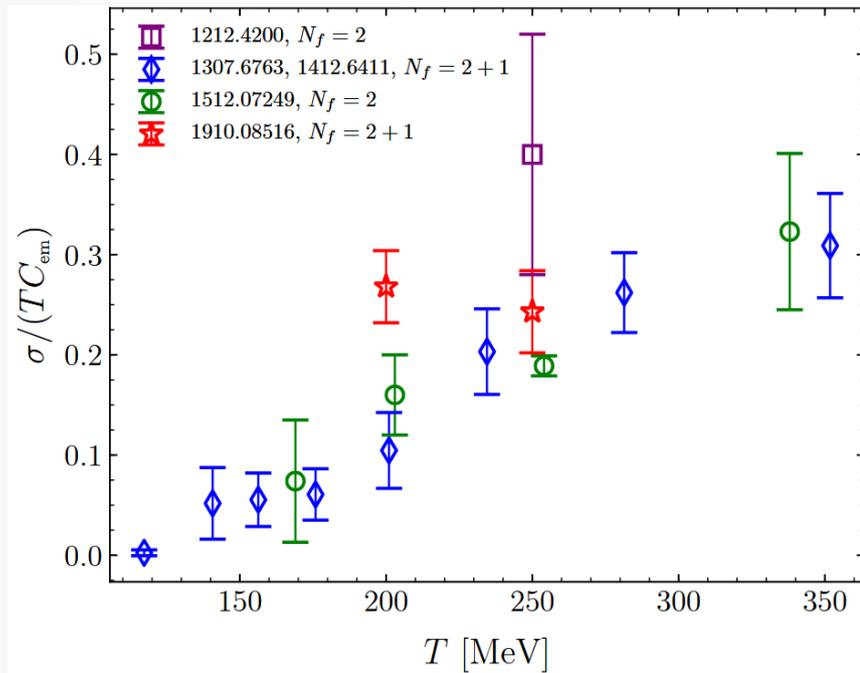
Review: G. Aarts and A. Nikolaev; EPJ. A 57, 118 (2021) arXiv:2008.12326 [hep-lat]



Quark-gluon plasma (QGP) electric conductivity

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- Uses linear-response theory (Kubo formula)
- Low energy limit - electromagnetic spectral function

$$\sigma = \frac{1}{6} \frac{\partial}{\partial \omega} \left(\int d^4x e^{i\omega t} \langle [j_\mu^{\text{em}}(t, x), j_\mu^{\text{em}}(0, 0)] \rangle \right) \Big|_{\omega=0}$$

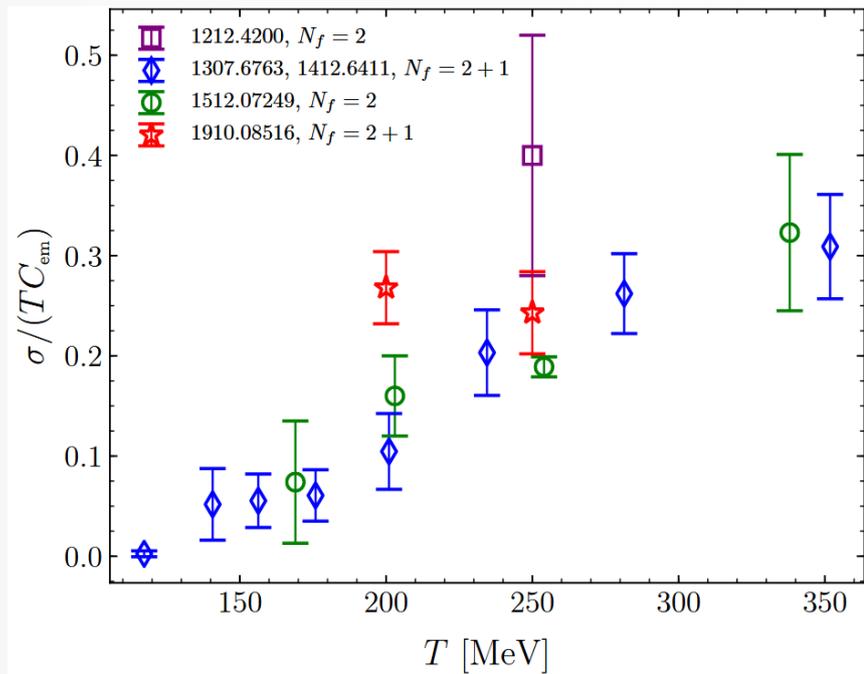
where the EM current depends on # of quark flavors

$$j_\mu^{\text{em}}(x) = \sum_{f=1}^{N_f} (eq_f) \bar{\psi}^f(x) \gamma_\mu \psi^f(x)$$

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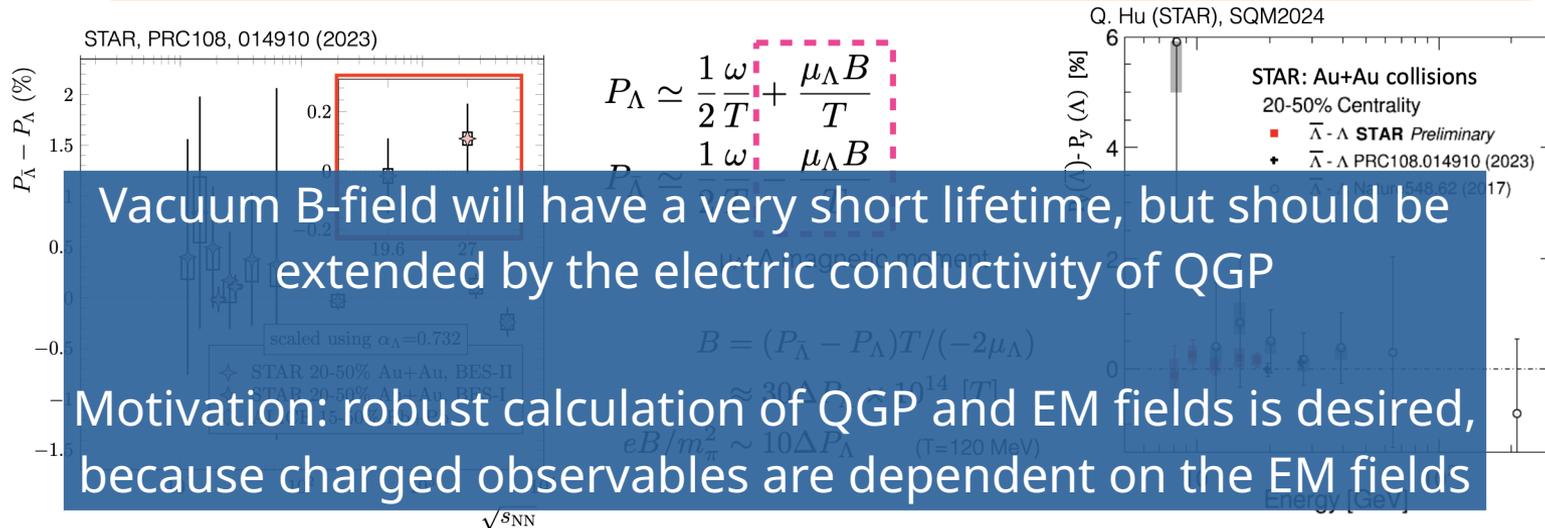
$$j_\mu^{\text{em}}(x) = \sum_{f=1}^{N_f} (eq_f) \bar{\psi}^f(x) \gamma_\mu \psi^f(x)$$

- Does not include external magnetic field effects
- Uses approximately realistic pion mass
- General agreement among results using a variety of methods and parameters (see paper)

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What observables can we use?

Charged leptons and thermal photons

Dileptons → Y. Akamatsu, H. Hamagaki, T. Hatsuda, and T. Hirano, PRC.85.054903 (2012).

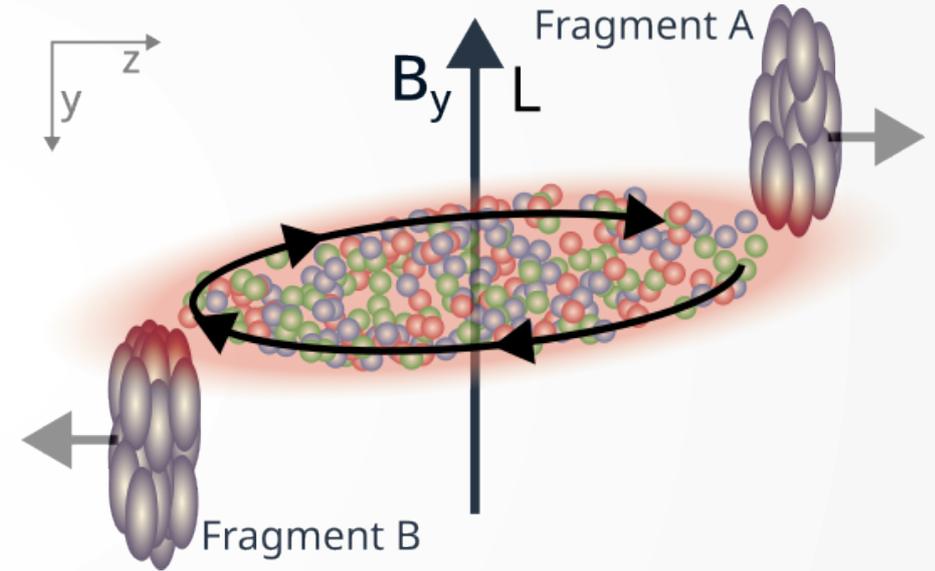
Photons → J.-A. Sun and L. Yan, PRC.109.034917 (2024)

Charge dependent directed flow (v_1)

Asymmetric collisions → Y. Hirono, M. Hongo, and T. Hirano, PRC.90.021903 (2014)

Symmetric collisions → U. Gürsoy, D. Kharzeev, and K. Rajagopal PRC.89.054905 (2014)

As some example papers.



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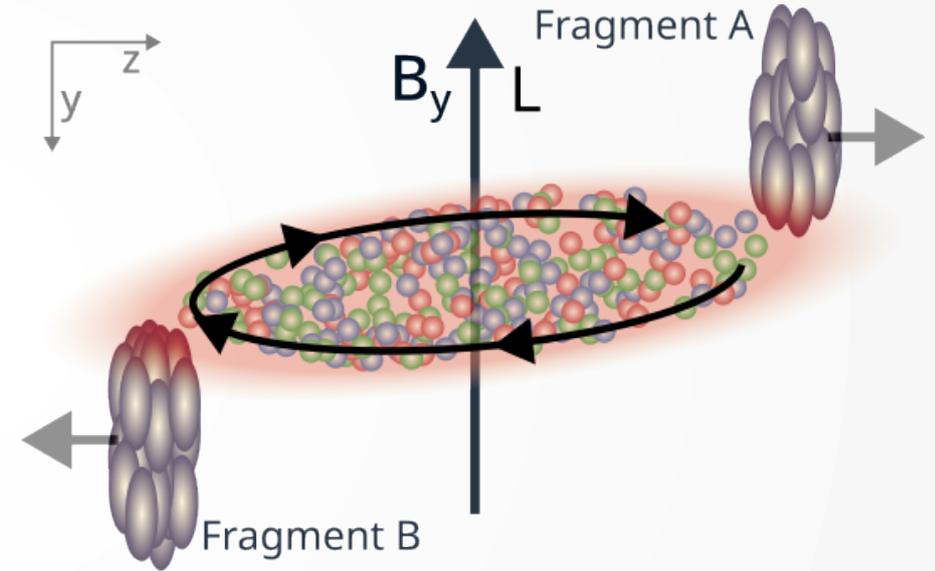
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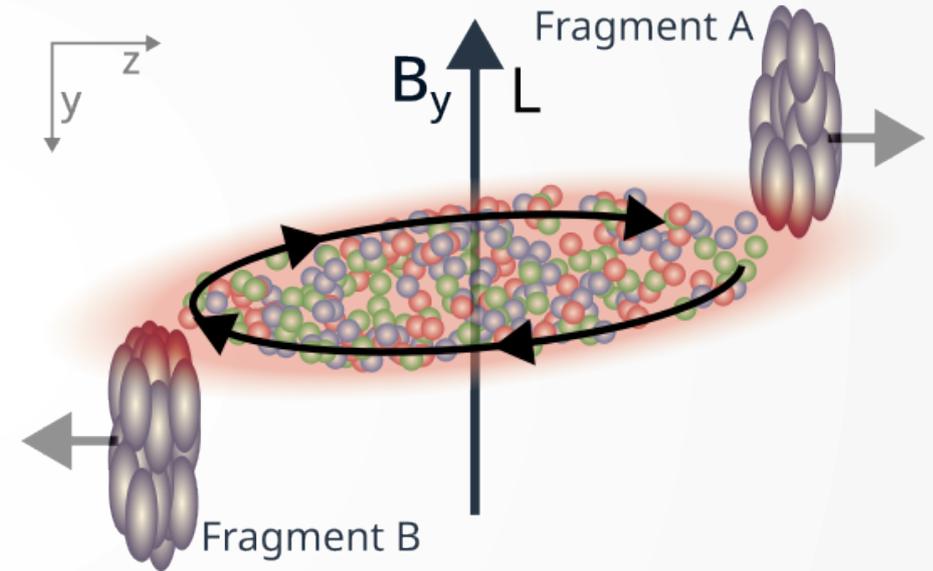
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This works focuses on these two

How can we model such a system?

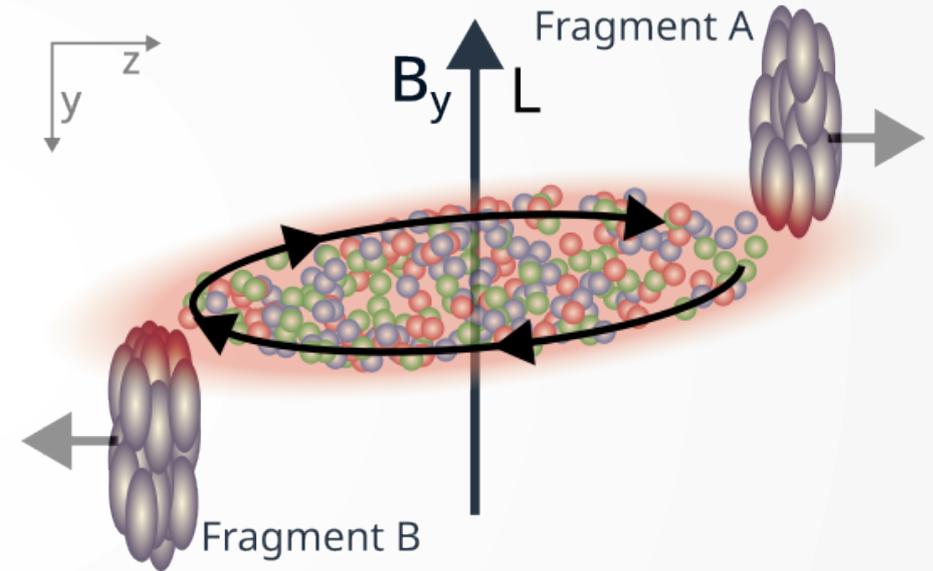
- In medium EM fields evolve at a similar rate as the QGP expansion (QGP is a conductive medium)
- EM fields could modify the collective behavior of charged particles
- Charged currents in the plasma should influence the EM fields



Use relativistic resistive magneto-hydrodynamics

How can we model such a system?

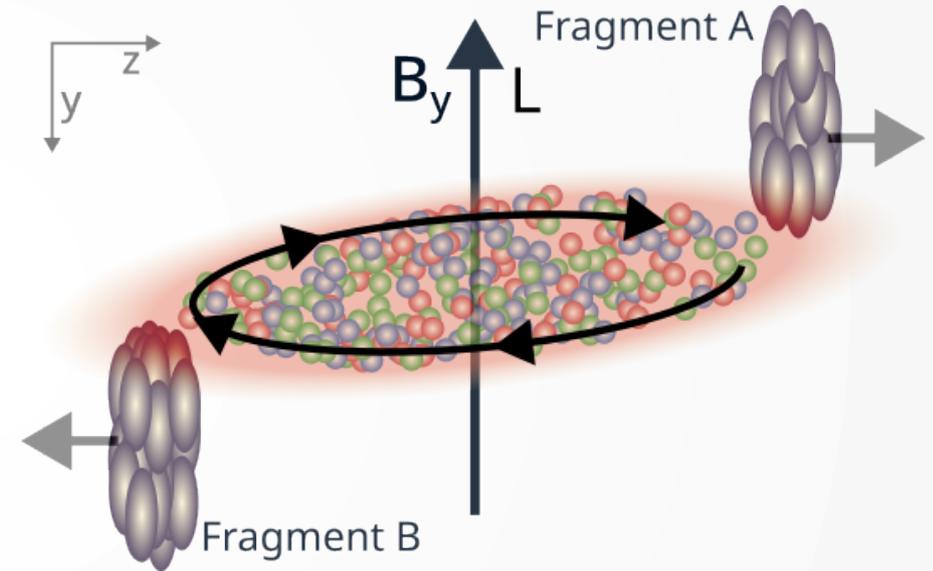
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Use relativistic **resistive** magneto-hydrodynamics

How can we model such a system?

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Use relativistic resistive **magneto-hydrodynamics**

Details of our RRMHD model

- Start from ideal-hydrodynamics + Maxwell equations

$$\nabla_{\mu} N^{\mu} = 0 \quad \text{Continuity equation (i.e., net-baryon current)}$$

$$\nabla_{\mu} T^{\mu\nu} = 0 \quad \text{Total energy-momentum}$$

$$\nabla_{\mu} S^{\mu} \geq 0 \quad \text{2nd law of thermodynamics}$$

Maxwell's equations

$$\nabla_{\mu} F^{\mu\nu} = -J^{\nu}$$

$$\nabla_{\mu} {}^*F^{\mu\nu} = 0$$

$$T^{\mu\nu} = T_{\text{m}}^{\mu\nu} + T_{\text{f}}^{\mu\nu}$$

$$T_{\text{m}}^{\mu\nu} = (\epsilon + p_{\text{gas}}) u^{\mu} u^{\nu} + p_{\text{gas}} g^{\mu\nu}$$

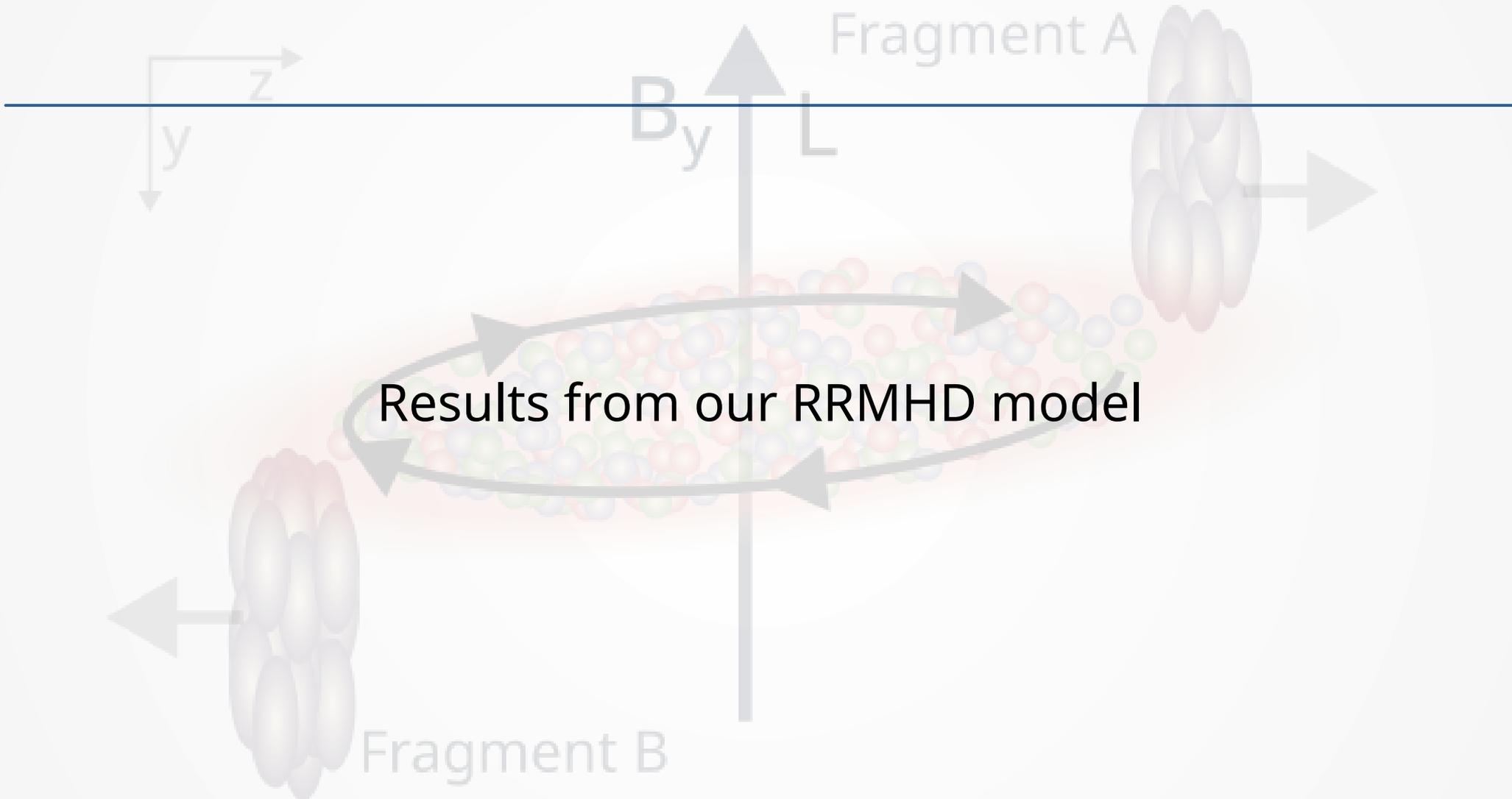
fluid matter part

$$T_{\text{f}}^{\mu\nu} = F^{\mu\lambda} F_{\lambda}^{\nu} - \frac{1}{4} g^{\mu\nu} F^{\lambda\delta} F_{\lambda\delta}$$

EM fields part

$$\nabla_{\mu} T_{\text{m}}^{\mu\nu} = -J_{\mu} F^{\mu\nu}$$

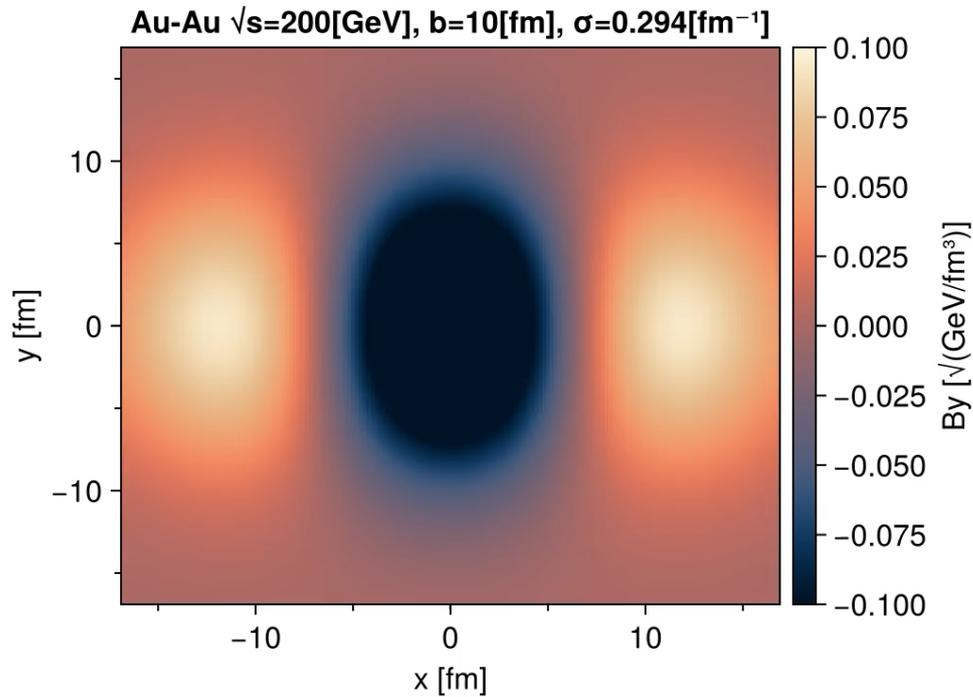
EM fields acting on the plasma and vice-versa



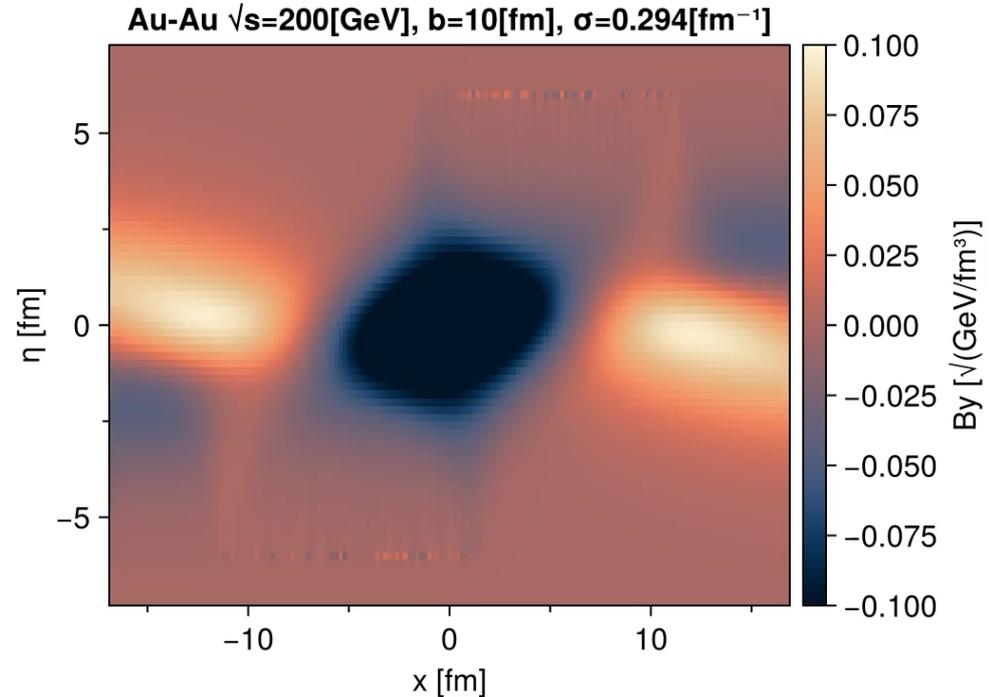
Results from our RRMHD model

Time evolution of the magnetic field

Transverse Plane ($\eta=0$)



Collision Plane ($y=0$)



Time evolution of the magnetic field

Our RRMHD model vs. an analytic estimation

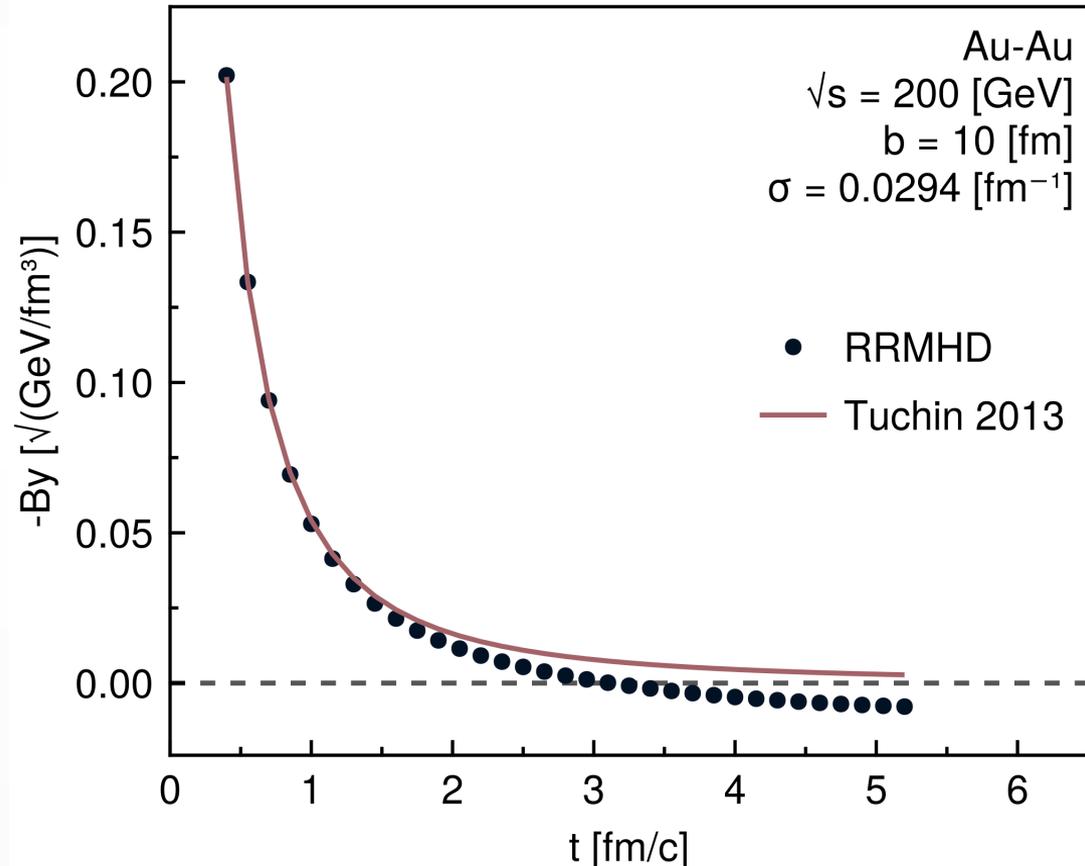
Tuchin; PRC.88.024911 (2013), arXiv:1305.5806 [hep-ph]

- Captured @ center of grid (collision)
- At early times they agree because the initial conditions are the same
- Late time is different because of **different source term** in calculation

$$\nabla_{\mu} F^{\mu\nu} = -J^{\nu}$$

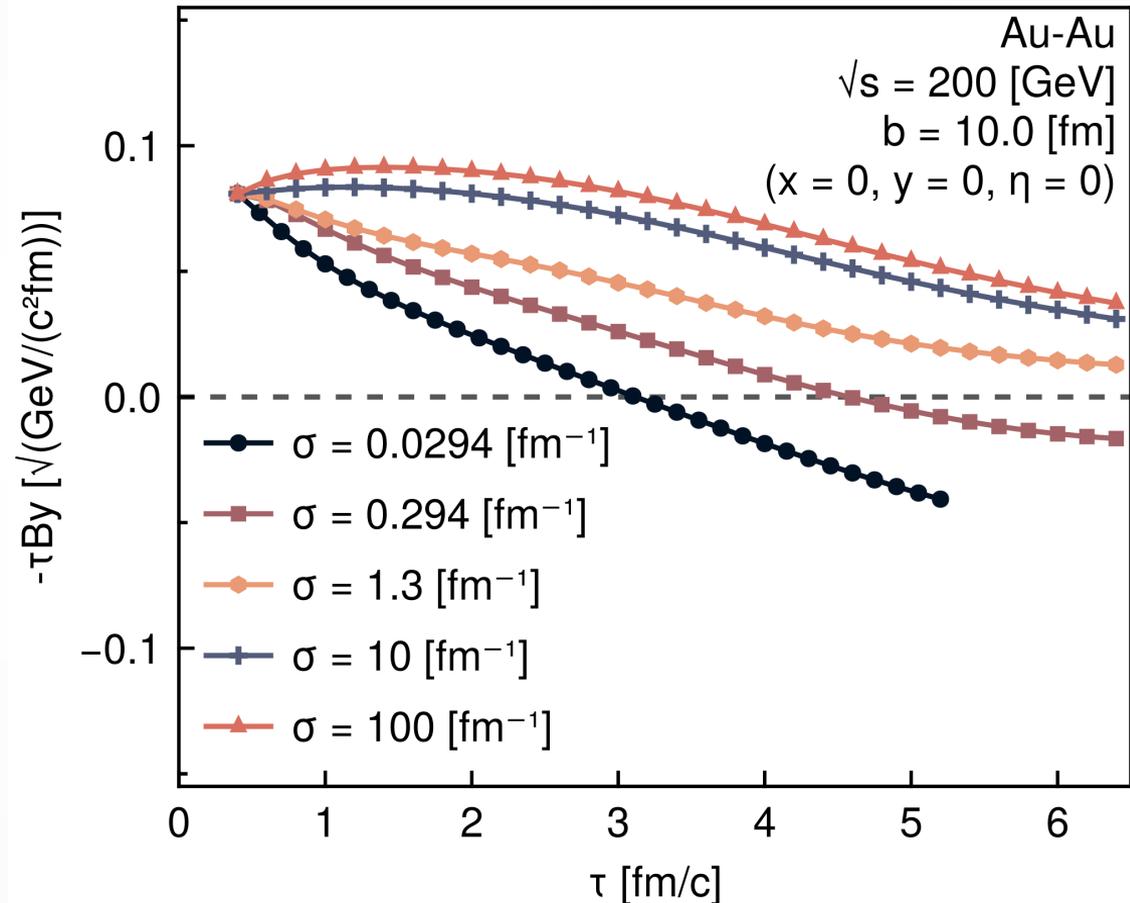
RRMHD = only QGP sources
Analytic = only collision spectators

$$\nabla_{\mu} {}^*F^{\mu\nu} = 0$$

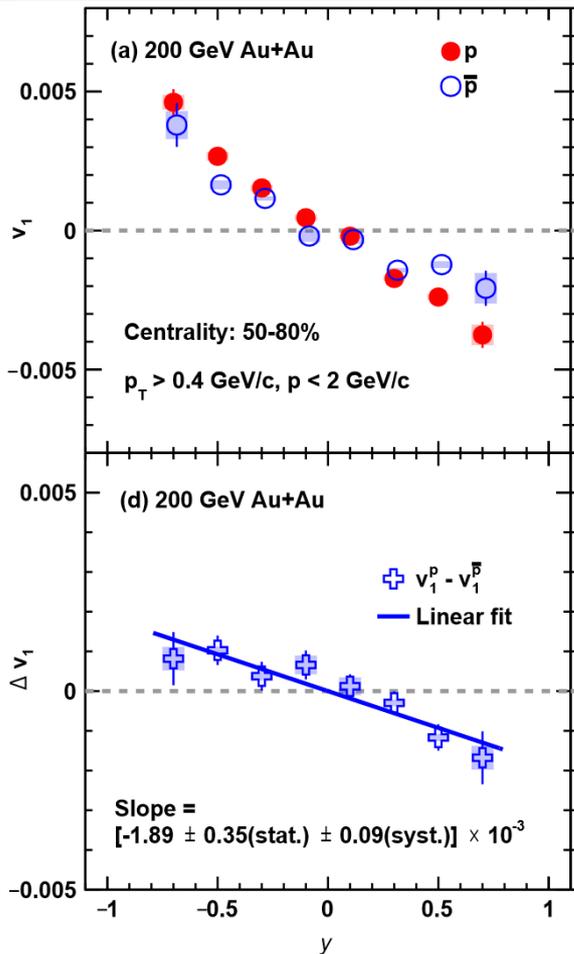


Time evolution of the magnetic field

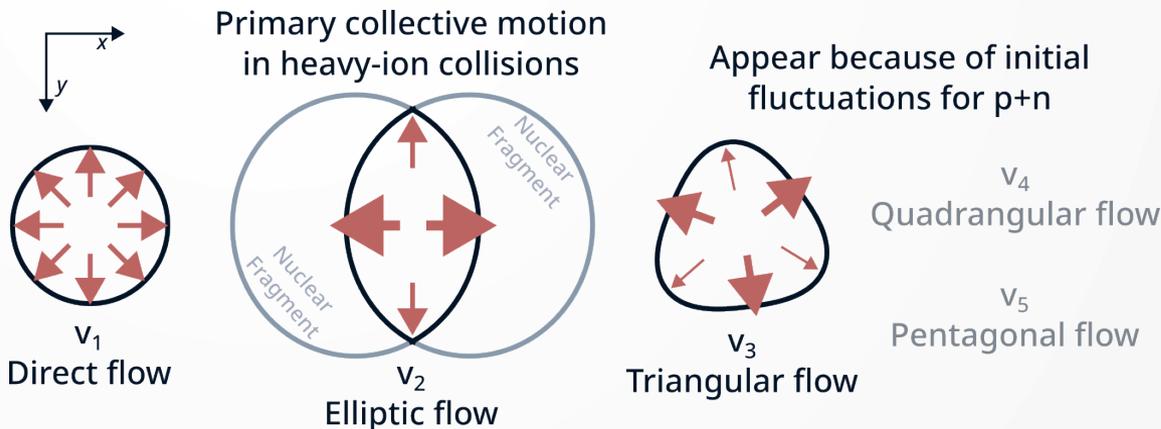
- Captured @ center of grid (collision) starting from the same initial condition
- A larger electric conductivity means longer field lifetime



Recent STAR experimental result

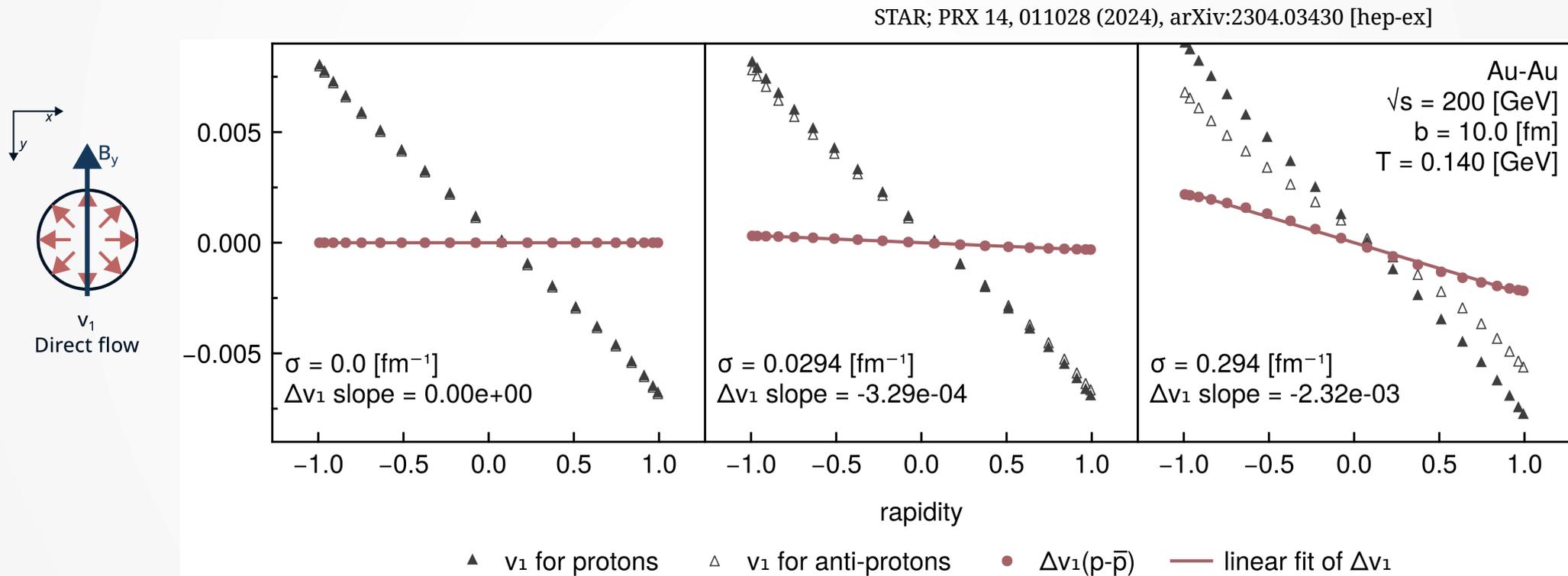


- New 2024 results from STAR experiment for RHIC @ Phys. Rev X 14, 011028 (2024) arXiv:2304.03430 [hep-ex]
- They measure direct flow of charged hadrons like protons, kaons, and pions
- And they calculate the difference in flow between negatively – positively charged hadrons



RRMHD Δv_1 results

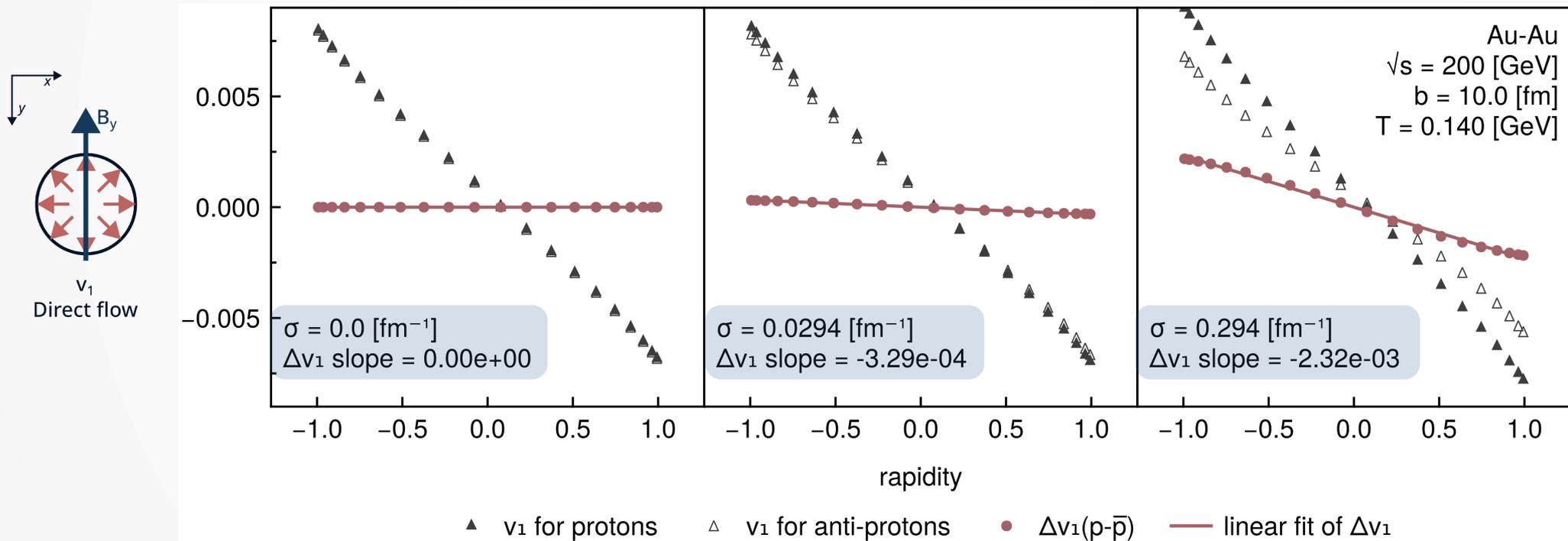
- Our RRMHD model can reproduce the negative slopes from STAR results



RRMHD Δv_1 results

- Our RRMHD model can reproduce the negative slopes from STAR results
- Value of the negative slope depends on the electric conductivity of QGP

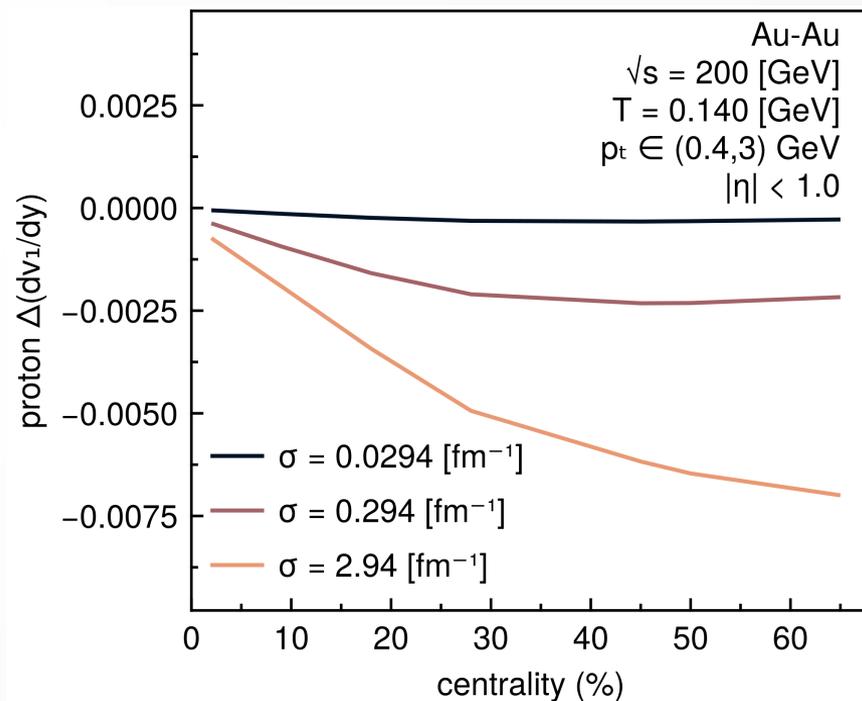
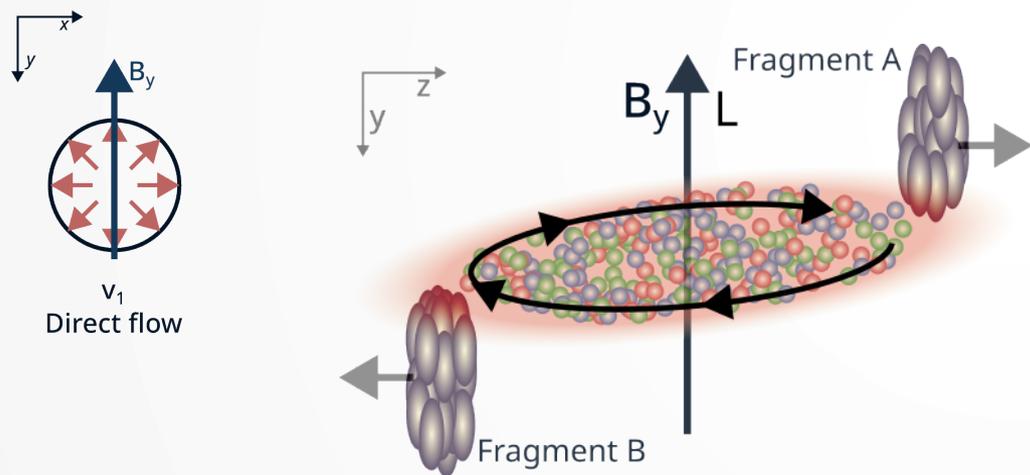
STAR; PRX 14, 011028 (2024), arXiv:2304.03430 [hep-ex]



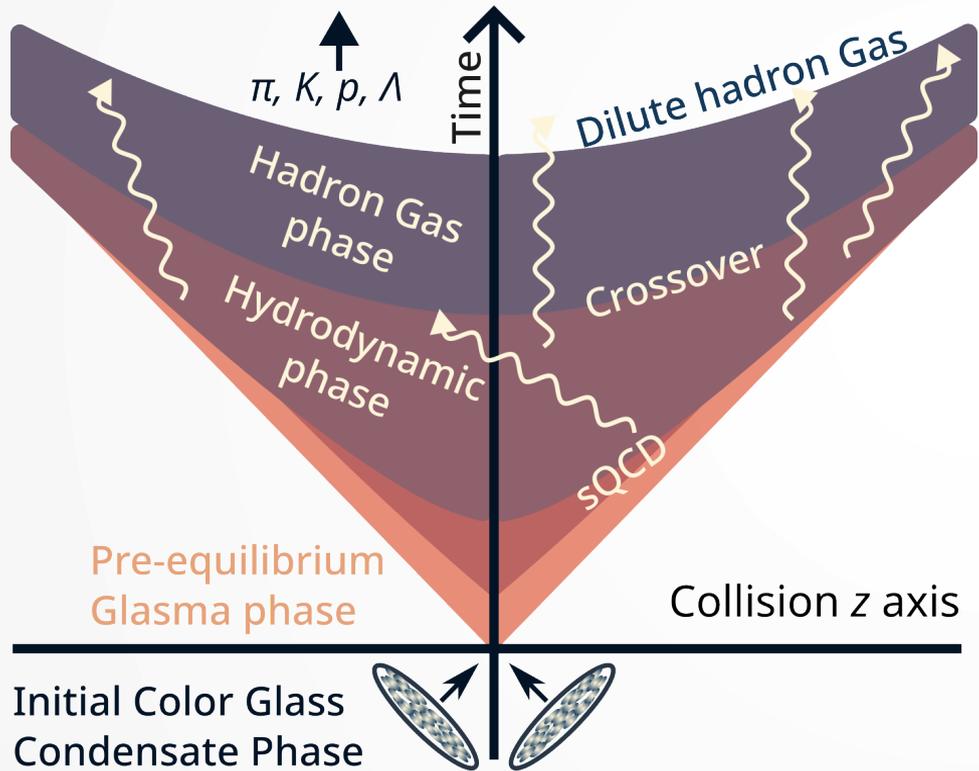
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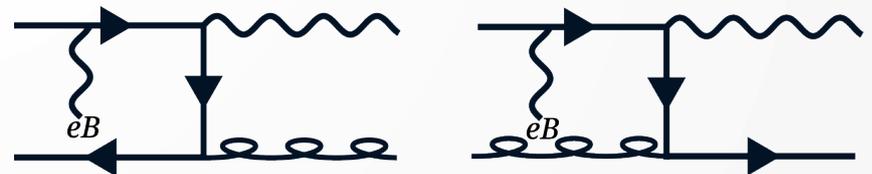
Thermal Photons from quark-gluon plasma



- Quark-gluon plasma will radiate photons as it expands and cools



- EM fields can increase the thermal photon production



Thermal Photons from quark-gluon plasma

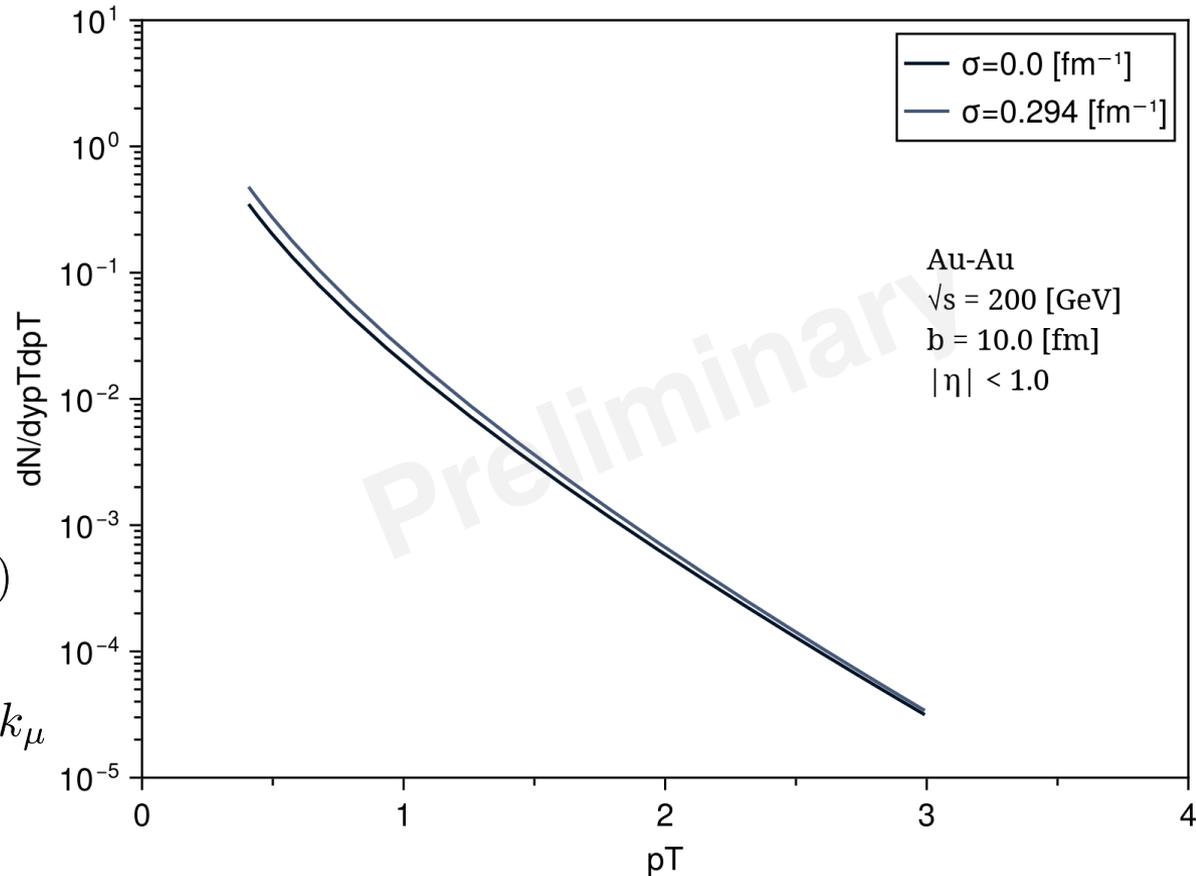
- Small number of photons produced, because EM fields are “weak” compared to QGP

$$E_k \frac{d\mathcal{R}}{d^3\vec{k}} = E_k \frac{d\mathcal{R}^{\text{QGP}}}{d^3\vec{k}} + E_k \frac{d\mathcal{R}^{\text{EM}}}{d^3\vec{k}}$$

$$E_k \frac{d\mathcal{R}^{\text{EM}}}{d^3\vec{k}} \simeq C \alpha_s \alpha_{\text{EM}} \mathcal{I} \mathcal{L}_c \sum_a \delta f_{a,\text{EM}}^{(1)}(X, k)$$

$$\delta f_{a,\text{EM}}^{(1)}(X, k) = -\frac{-f_{a,\text{eq}}(1 - f_{a,\text{eq}})}{T \chi_{ei} k^\mu u_\mu} e \sigma Q_a e^\mu k_\mu$$

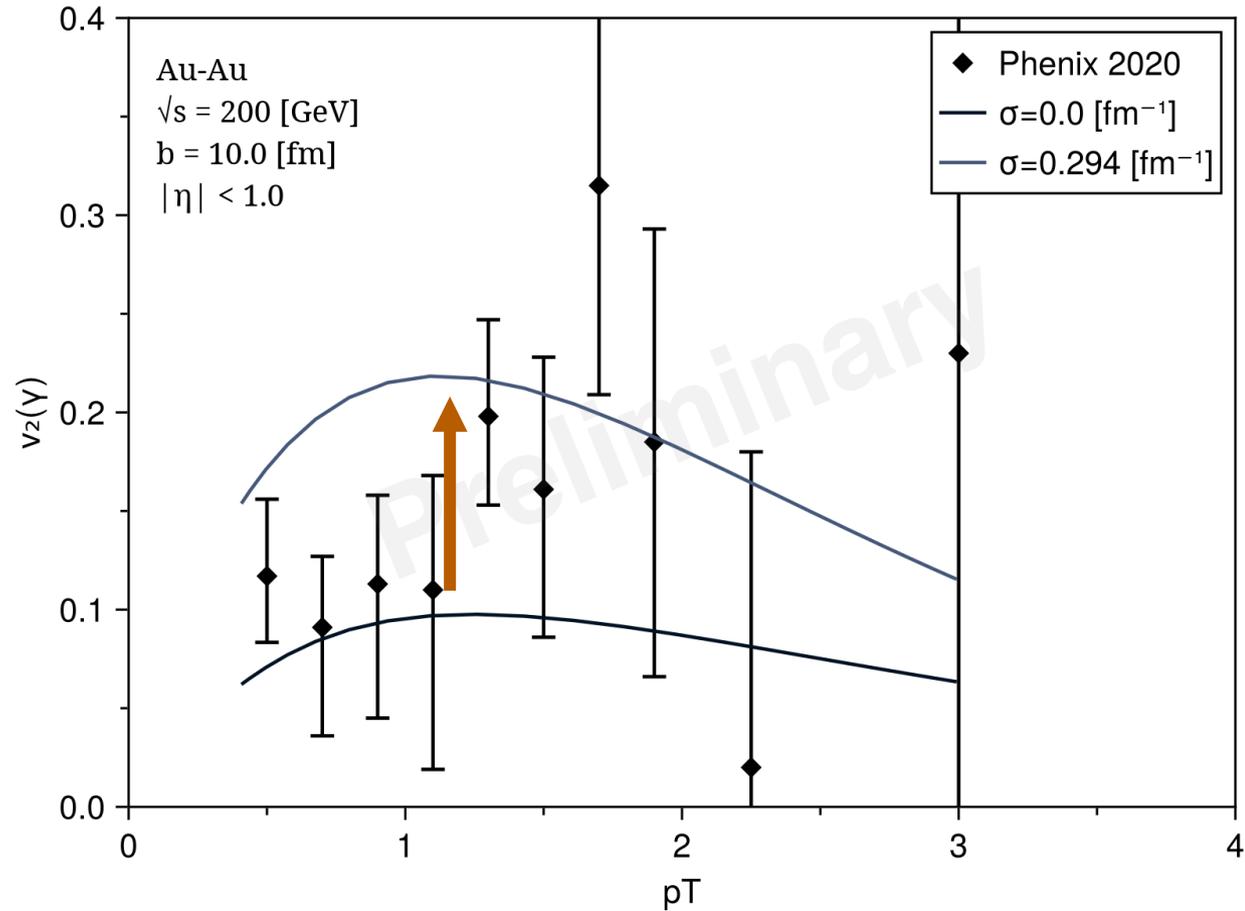
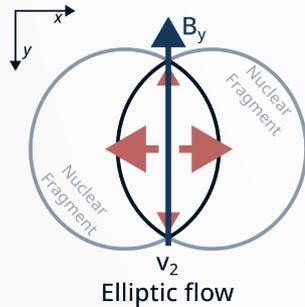
J.A. Sun and L. Yan; PRC.109.034917 arXiv:2311.03929 [hep-ph]



Thermal Photons from quark-gluon plasma

- Enhancement of the photon elliptic flow when the EM fields are included

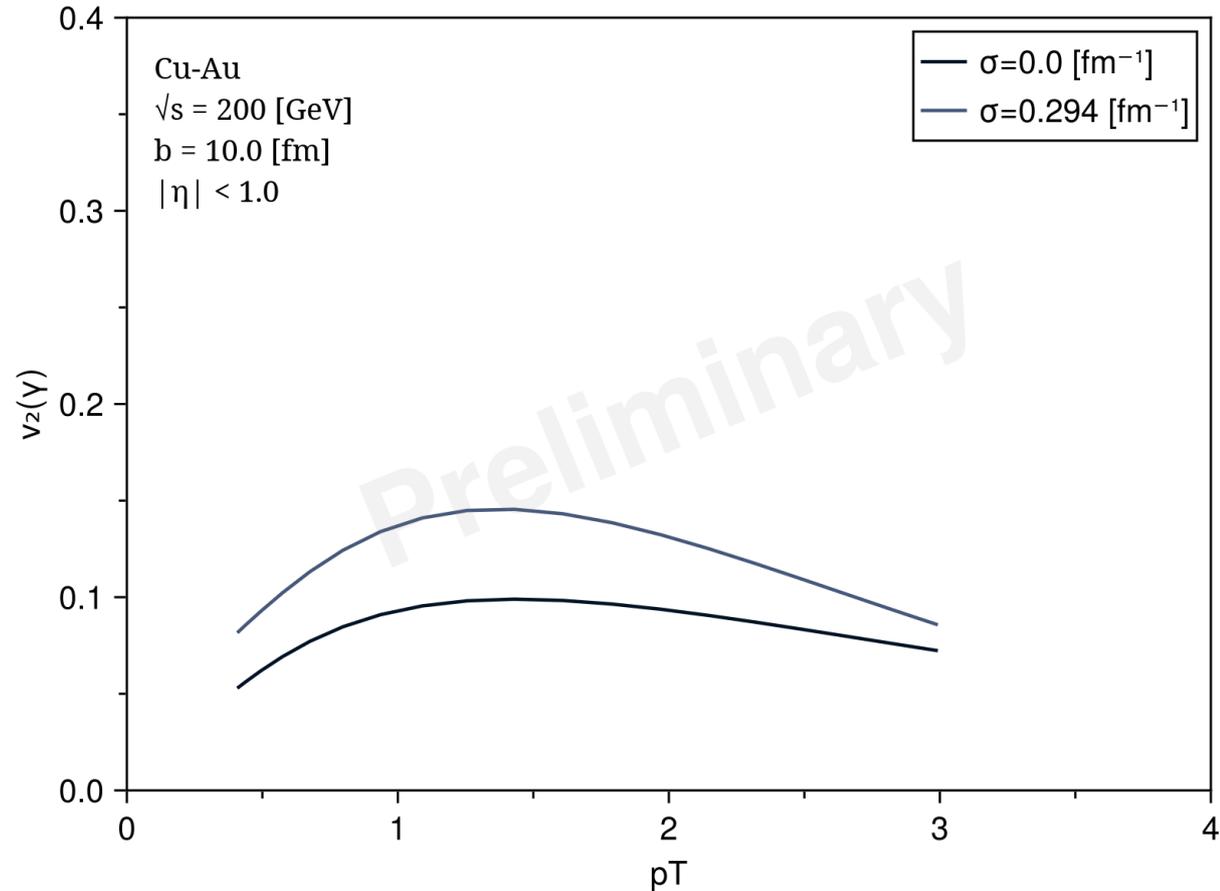
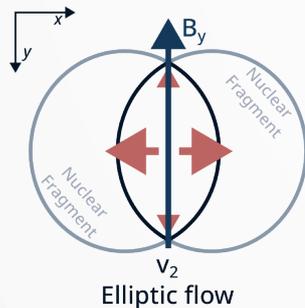
$$v_2(\gamma) \equiv \frac{v_0 v_2 + v_0^{\text{EM}} v_2^{\text{EM}}}{v_0 + v_0^{\text{EM}}}$$



Thermal Photons from quark-gluon plasma

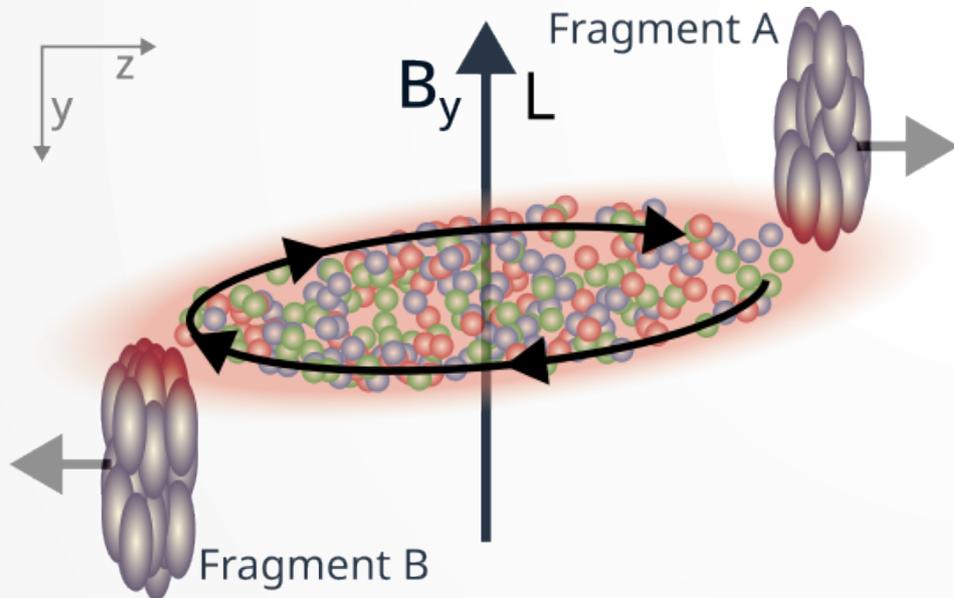
- Asymmetric collision the QGP thermal photon is similar to symmetric collision
- But, enhancement from the EM fields is much smaller

$$v_2(\gamma) \equiv \frac{v_0 v_2 + v_0^{\text{EM}} v_2^{\text{EM}}}{v_0 + v_0^{\text{EM}}}$$



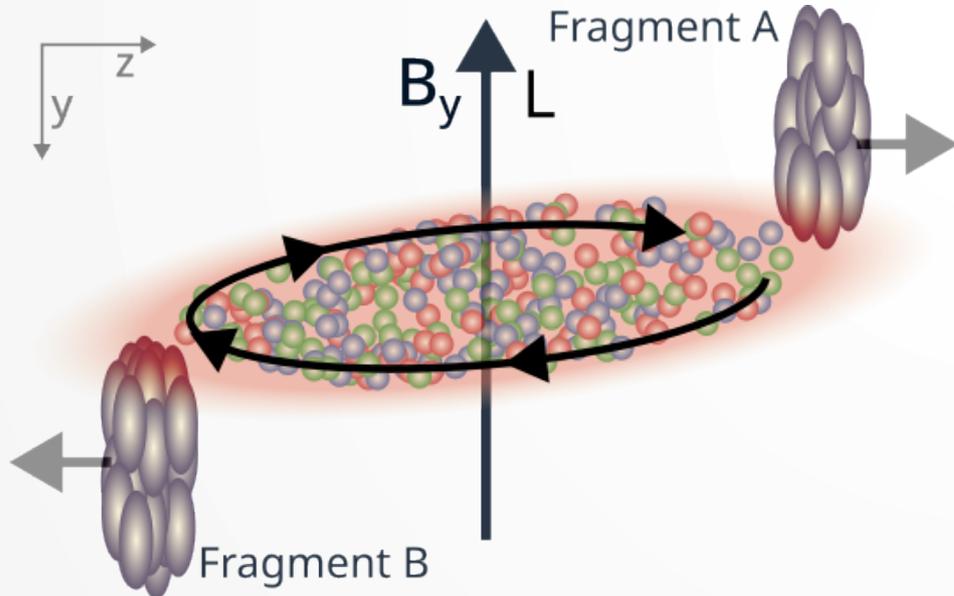
Summary and Outlook

- We have applied relativistic resistive magneto-hydrodynamics (RRMHD) to studying charged observables in heavy-ion collisions
- Our results show how important EM fields are to those charged observables
- Progress toward a robust calculation of QGP+EM fields



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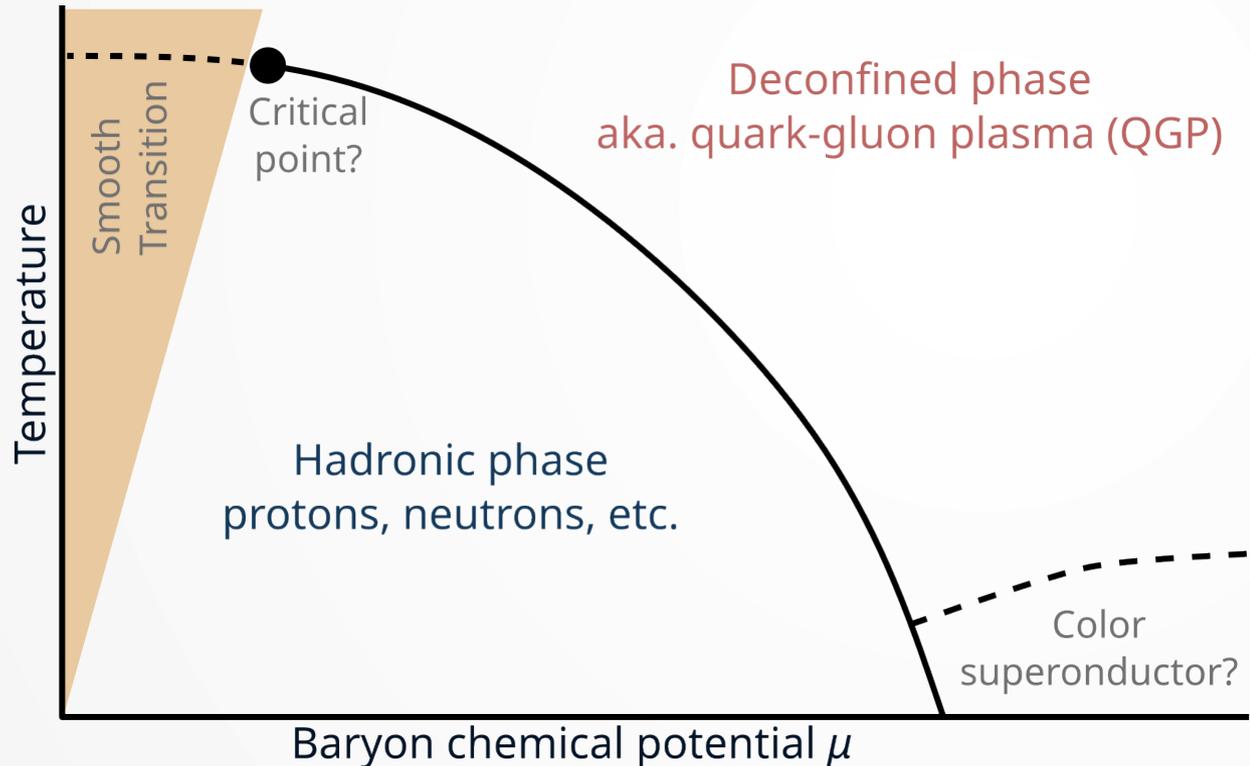
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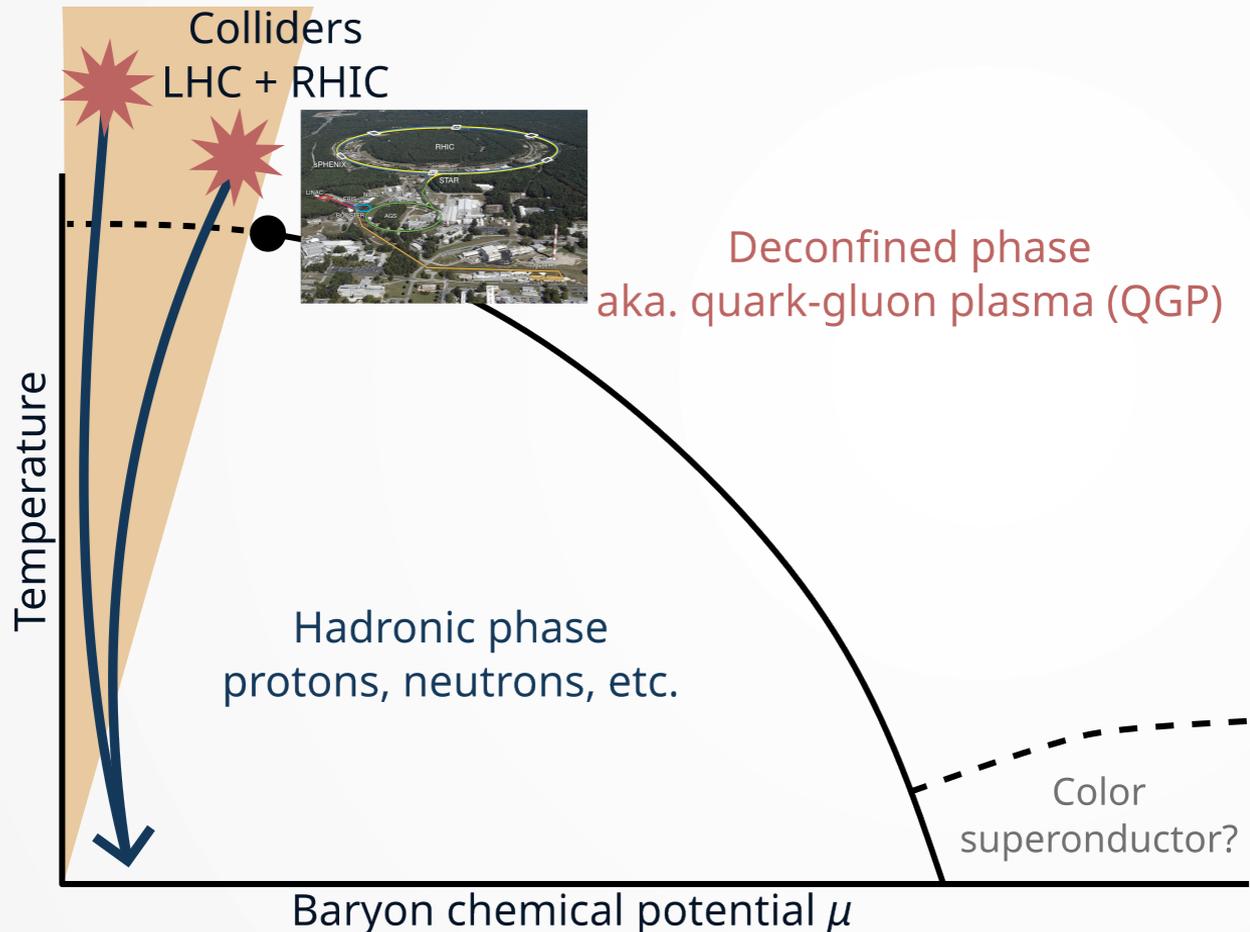
- In the future,
 - Non-equilibrium initial state
 - Hadron Gas + EM fields phase
 - Hydrodynamic fluctuations
 - Vortexes + EM fields
 - etc.
- Lots of work and results to come!

Heavy-ion collisions

- Heavy-ion collisions study ordinary matter under extreme conditions



Heavy-ion collisions

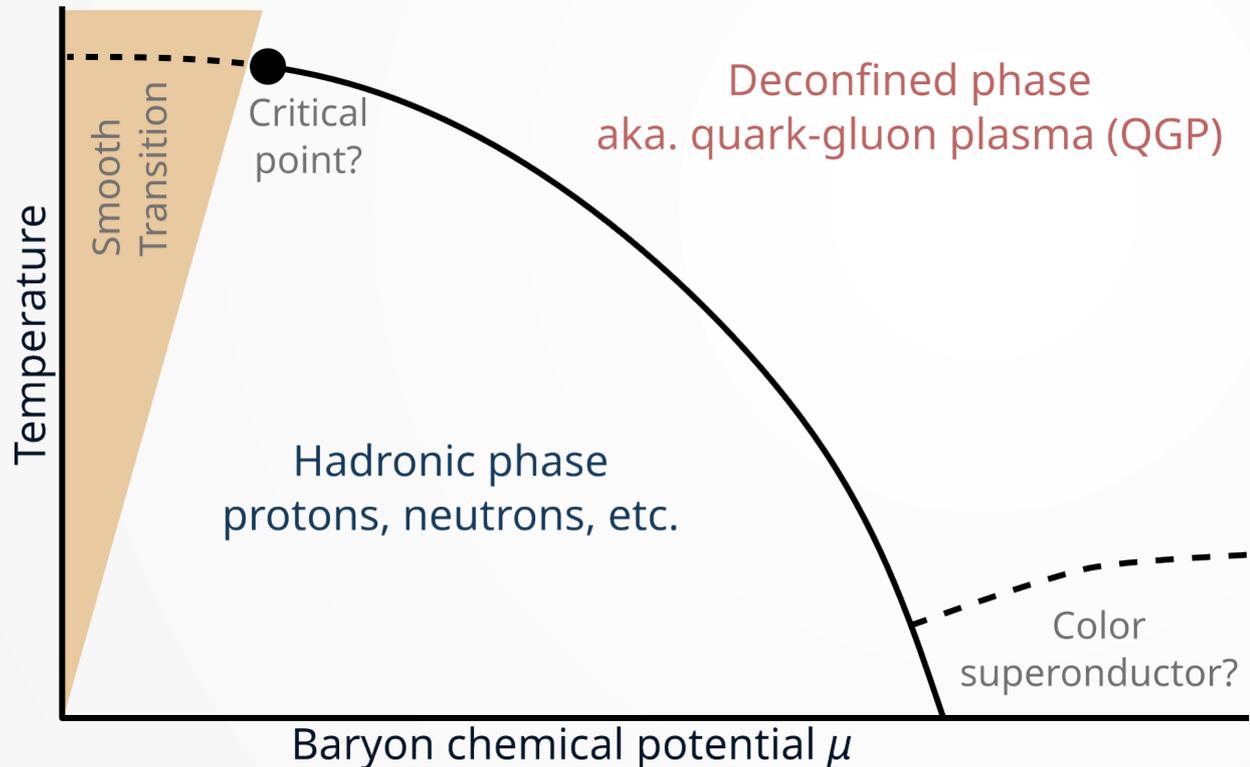


We use high-energy colliders to study the deconfined phase of QCD matter called quark-gluon plasma (QGP)

SPS @ CERN (1982~2001?)
RHIC @ BNL (2001~2025)
LHC @ CERN (2010~)

Why study quark-gluon plasma?

- Better understanding of the Strong force and confinement
- Possible similar conditions as the early universe (high temperature)



- Medium energy collisions can study the QCD phase transition and critical point
- Low energy collisions can be connected to studies of neutron stars (large baryon chemical potential)

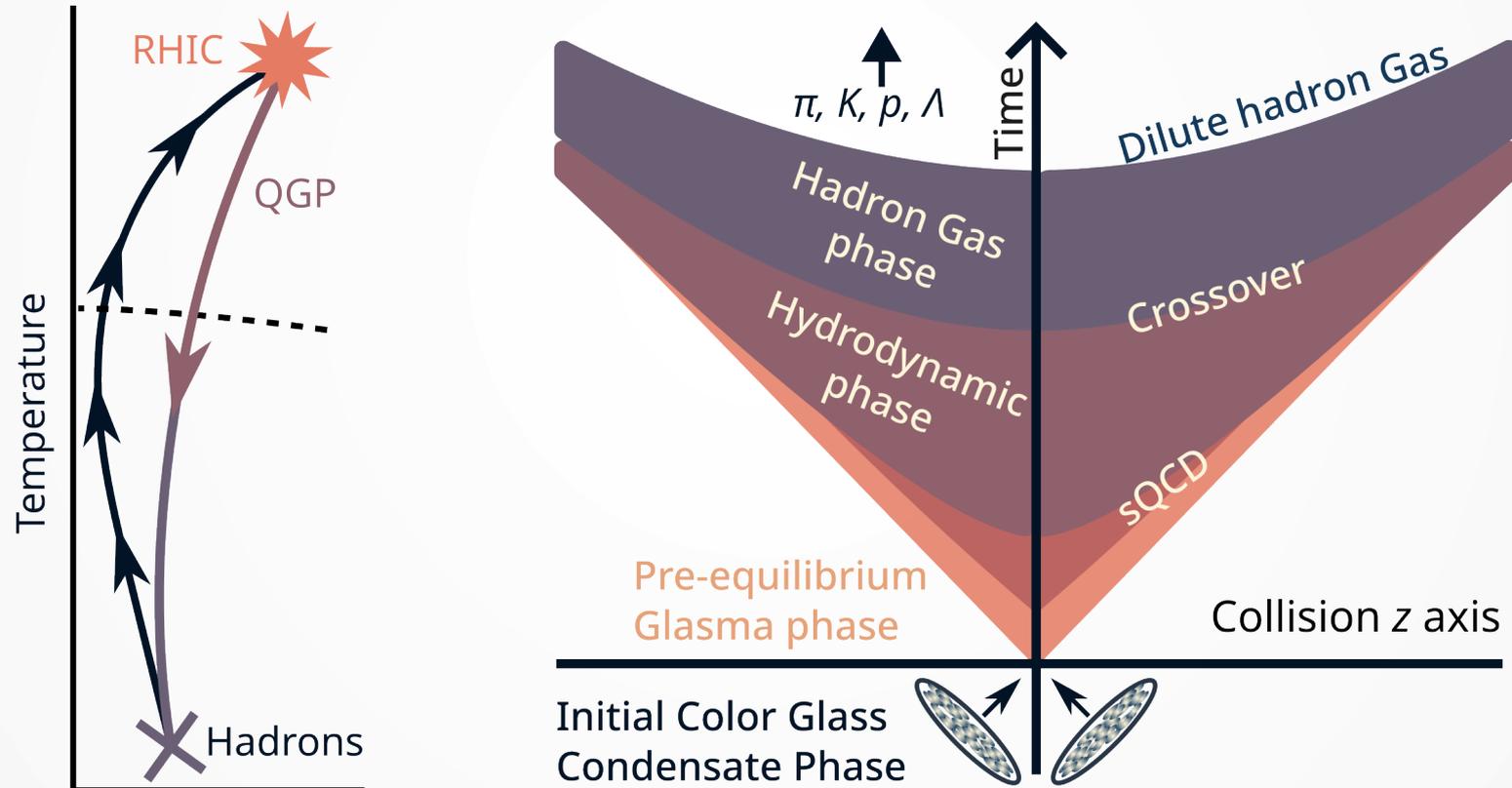
What kind of collective motion appears?

- Details of the collective motion are captured in transport parameters
- For QGP we have,

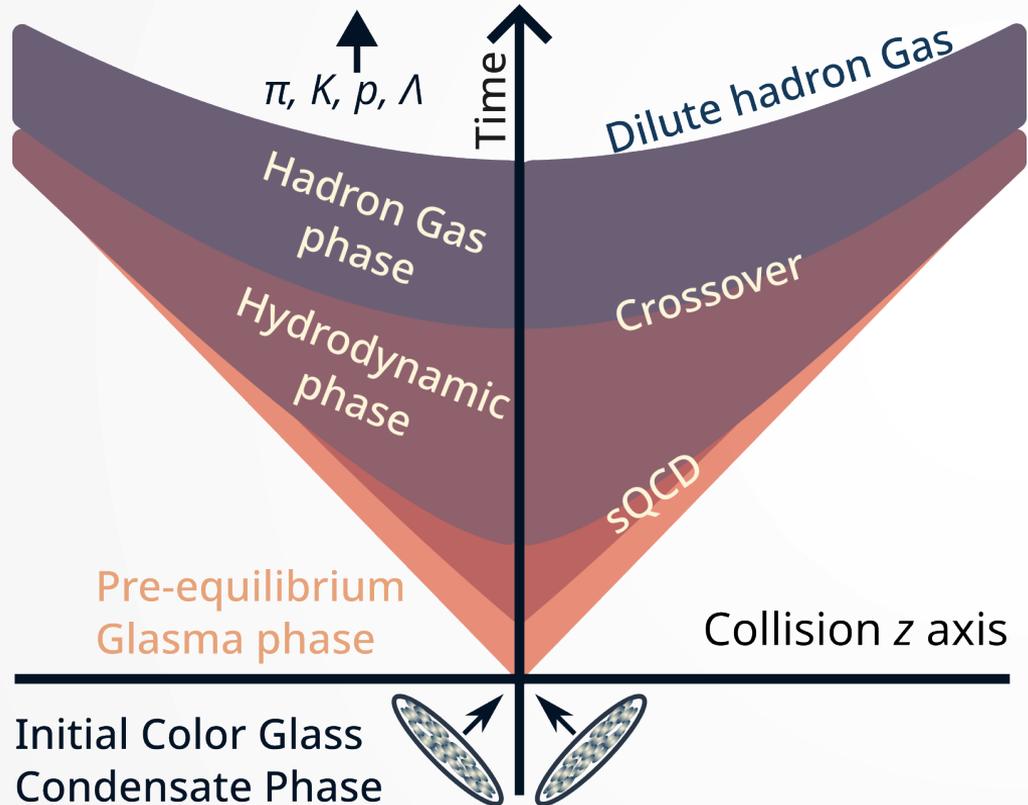
Shear viscosity	η	Energy-momentum
Bulk viscosity	ζ	
Charm-diffusion coefficient	D	Heavy-flavor quantum numbers
Thermal conductivity	κ	Heat via baryon current
Electric conductivity	σ	Electrical charges via the electric current

Stages in relativistic heavy-ion collisions

- Studying QGP is not simple, many phases before experimental measurement



Why model using hydrodynamics?



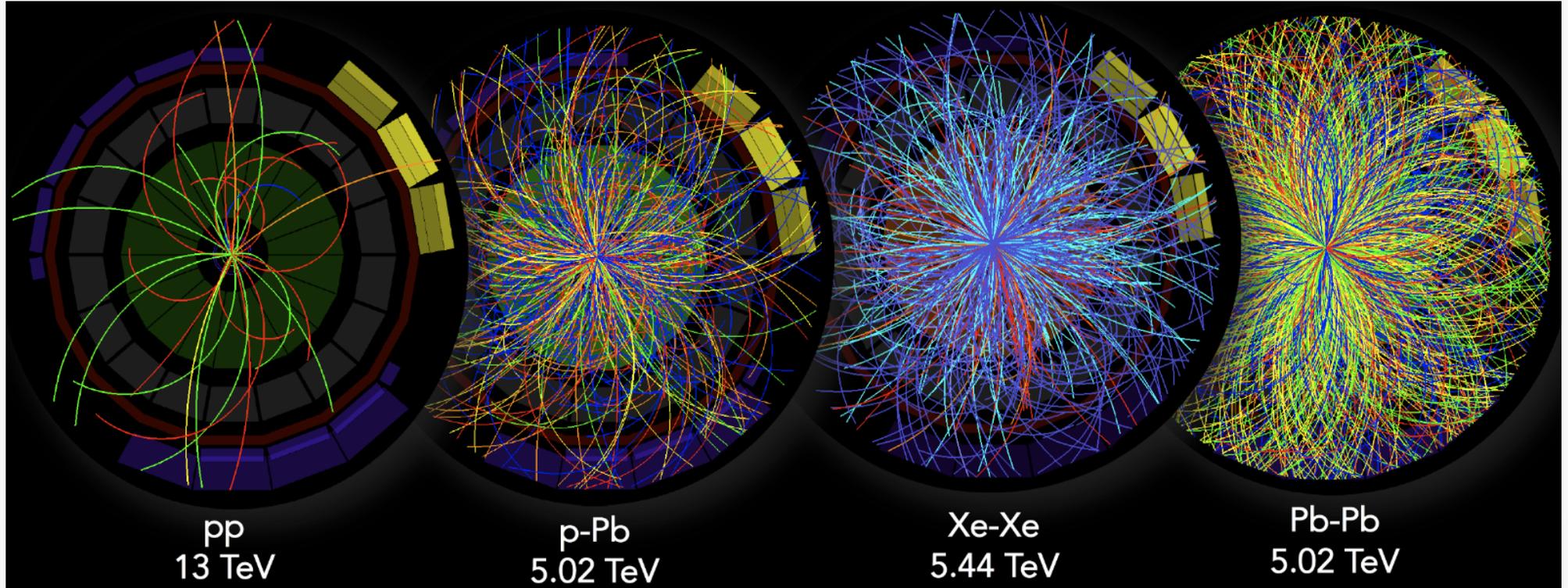
- Why the hydrodynamics phase?
- Many liberated quarks and gluons, so it becomes a many-body problem

What about kinetic transport models like the Boltzmann equation?

- Let's look to experiments for an answer

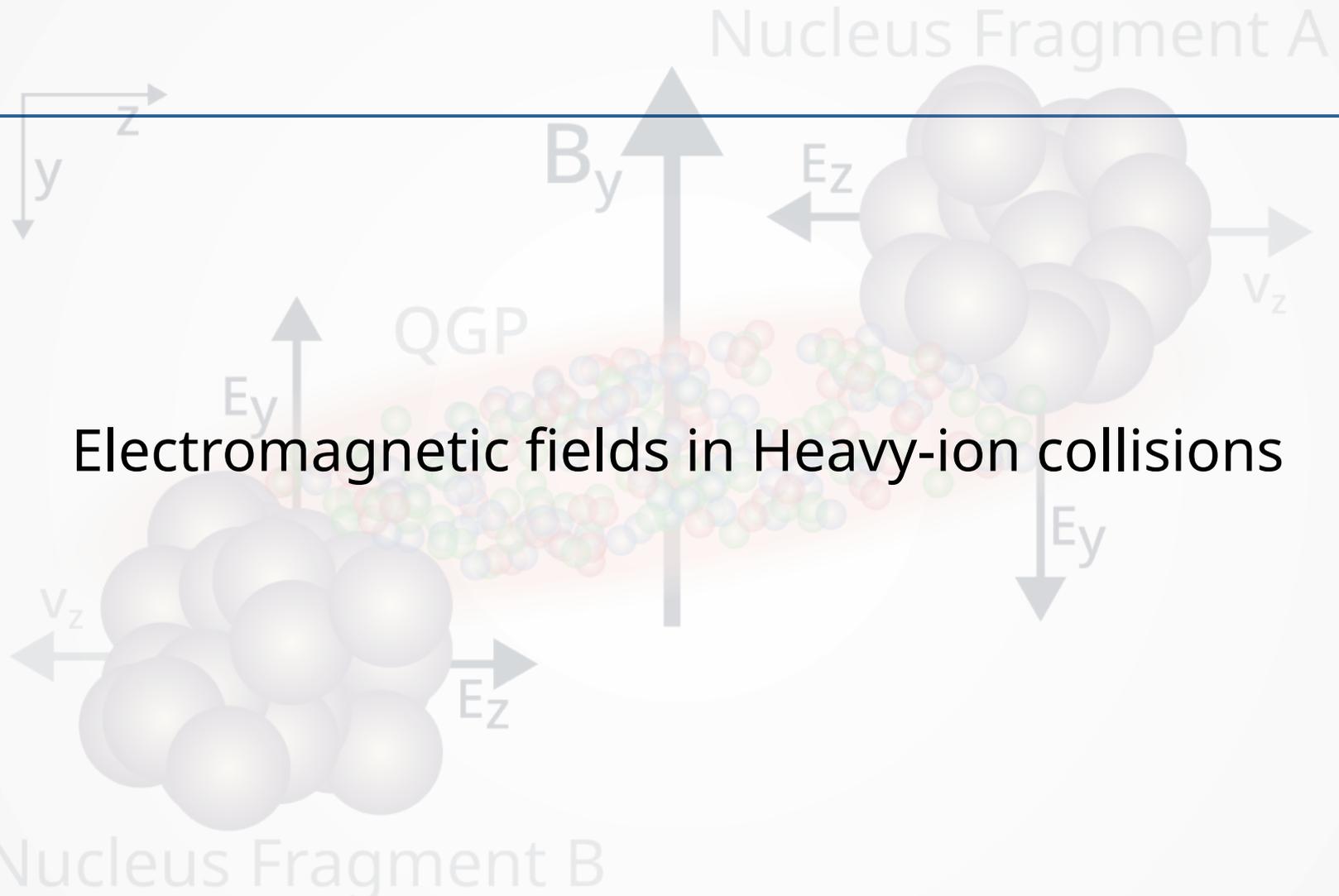
Why model using hydrodynamics?

- Heavy-ion collisions produce a large amount of particles



From: A.P. KALWEIT, PoS (EPS-HEP2023) 027

Electromagnetic fields in Heavy-ion collisions



What about the protons?

Resources: Jackson (1975), Feynman (2010)

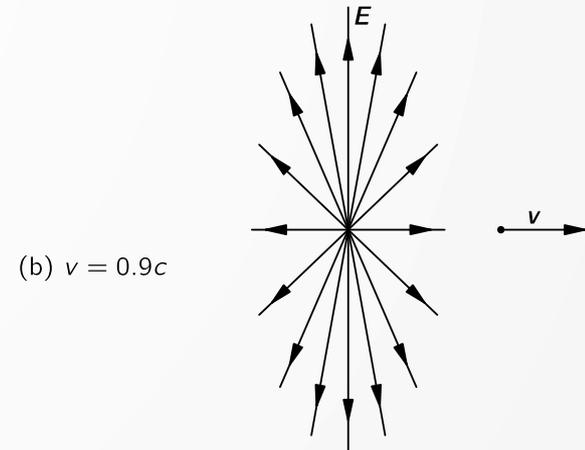
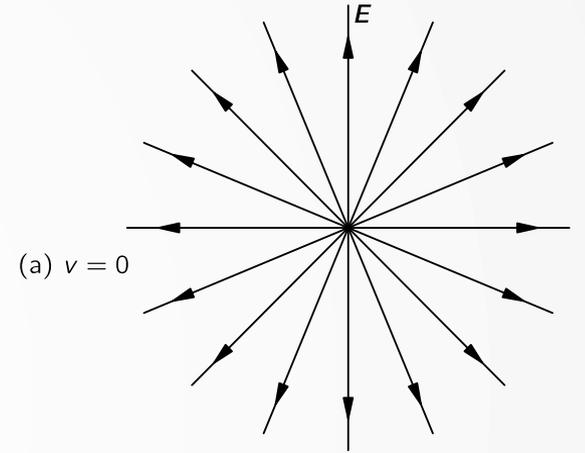
- Using a simple picture, of a point charge moving at a constant velocity in the z-direction

$$E_{x,y} = q\gamma \frac{\vec{x}}{[\vec{x}^2 + \gamma^2(z - vt)^2]^{3/2}},$$

$$E_z = q\gamma \frac{z - vt}{[\vec{x}^2 + \gamma^2(z - vt)^2]^{3/2}}$$

aka: Lienard-Wiechert Fields for a constant velocity

$$\vec{B} = \vec{v} \times \vec{E}$$

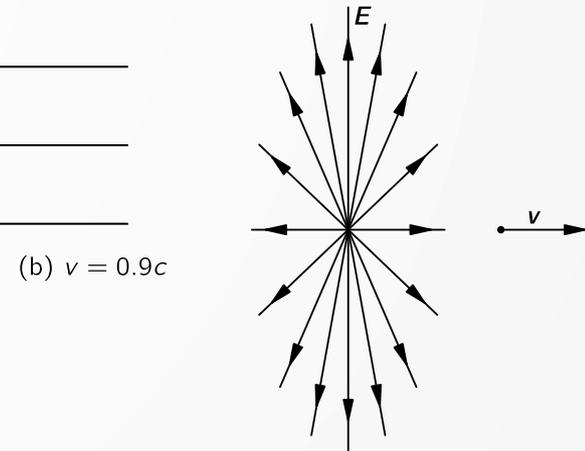
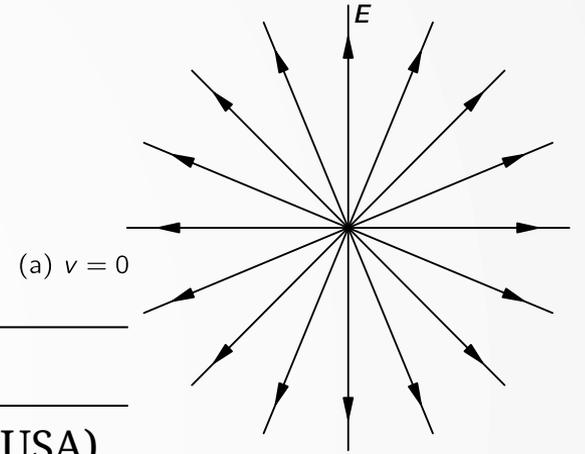


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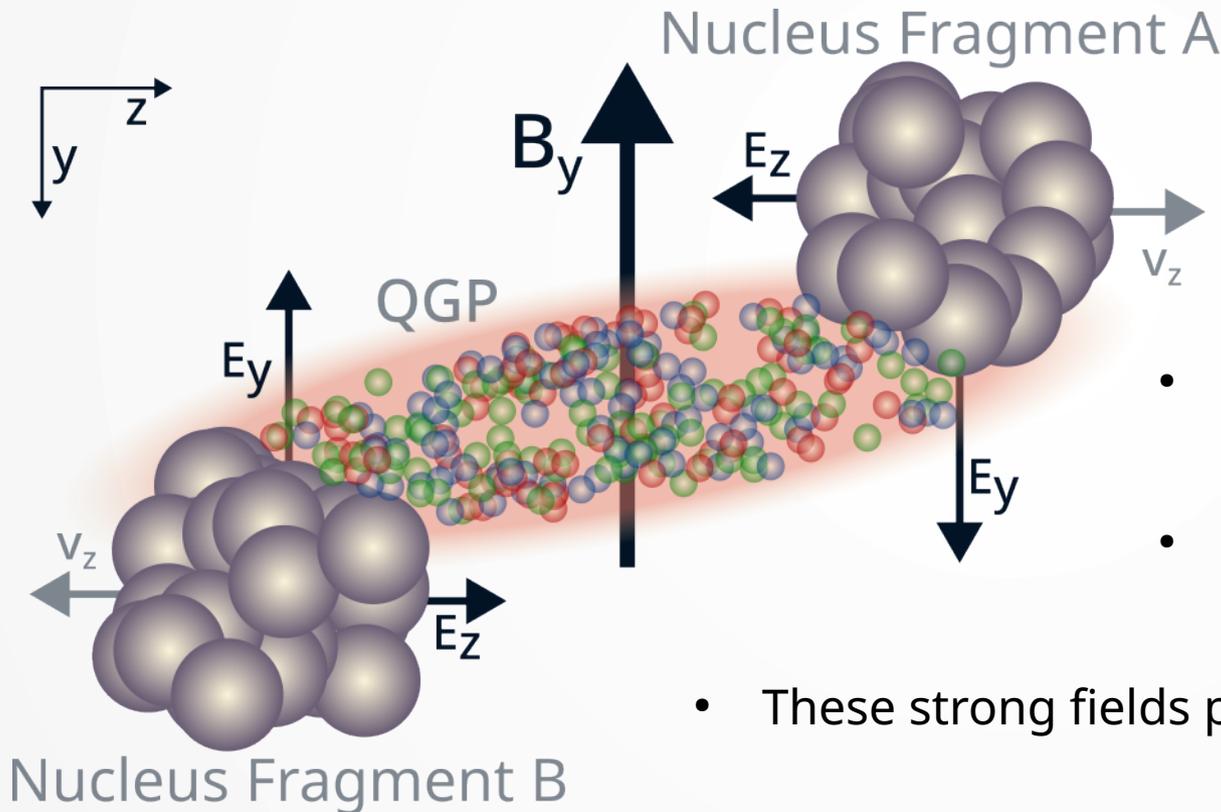
- Summing over all of the charges in a nuclei, produces the strongest EM fields ever created

$\sim 10 \text{ T}$	MRI
$\sim 45 \text{ T}$	Strongest continuous magnetic fields in labs (China & USA)
$\sim 10^9 \text{ T}$	EM fields become nonlinear (Schwinger limit)
$\sim 10^{13} \text{ T}$	Strongest magnetic pulsars
$\sim 10^{14} \text{ T}$	Possible field strength at 200 GeV RHIC collisions



Role of the electromagnetic (EM) fields

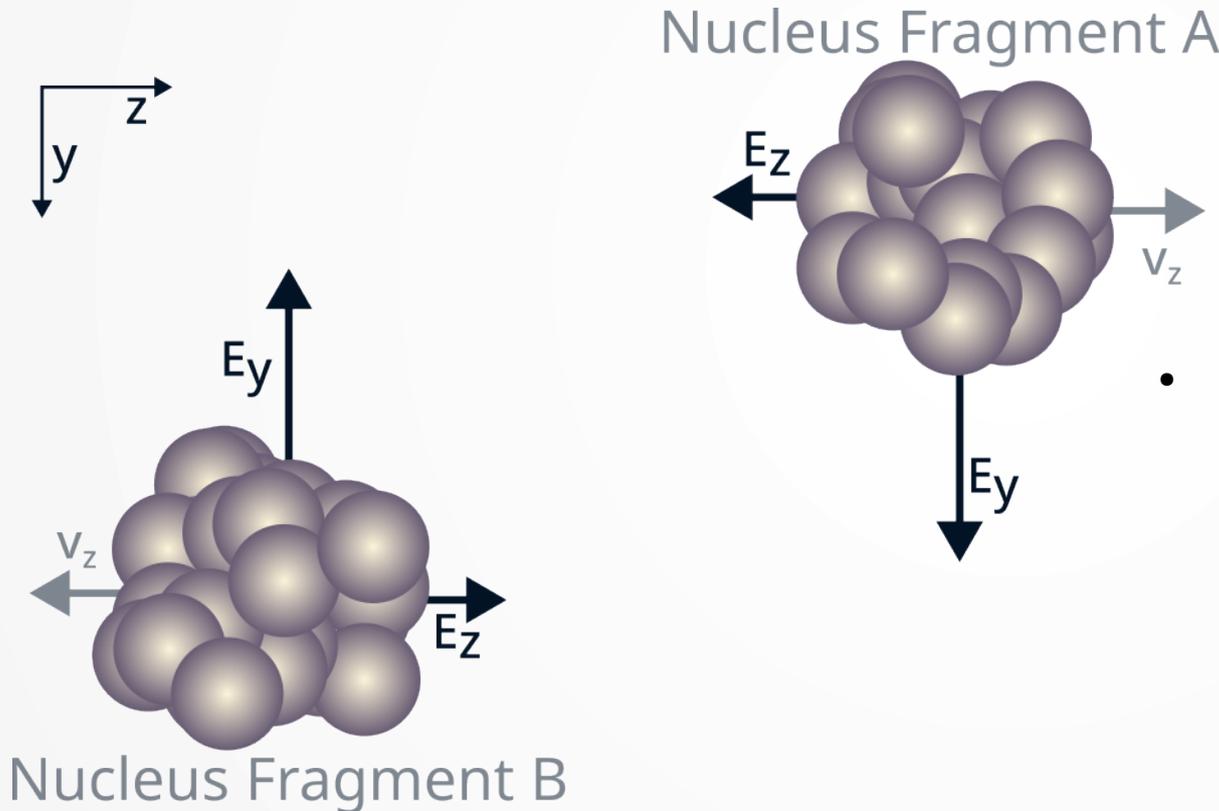
- Fragments of the collision play an important role as a source of the EM fields



- E-fields from fragmented protons are squished and stretched
- A strong B-field follows in the wakes of the moving nuclei
- These strong fields penetrate the quark-gluon plasma

Role of the electromagnetic (EM) fields

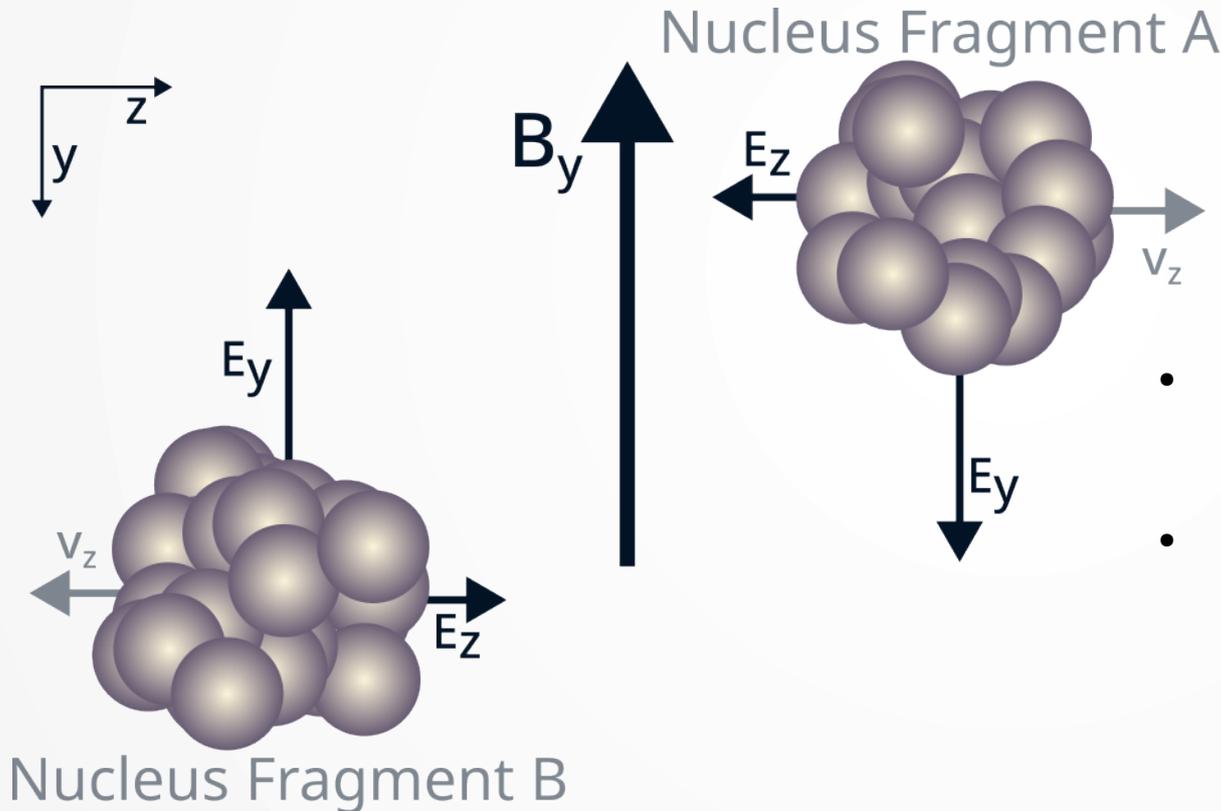
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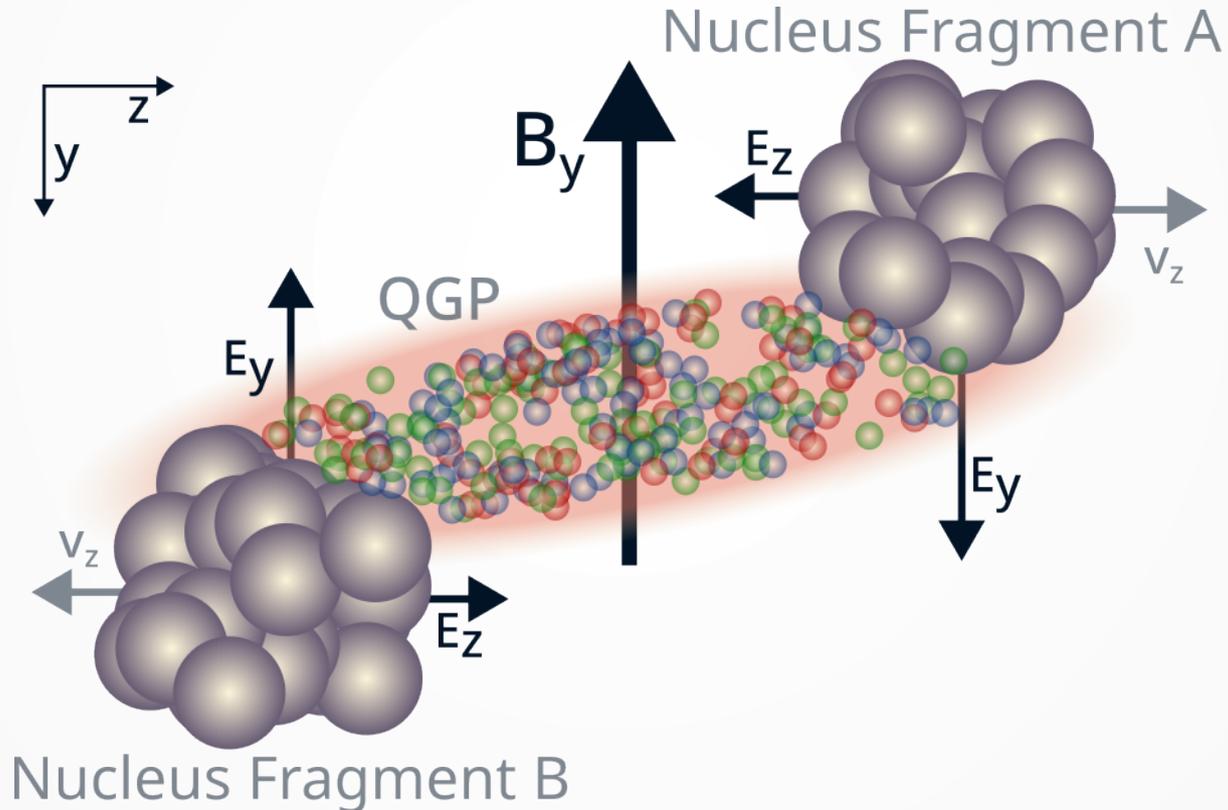
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History of relativistic resistive magneto-hydrodynamics

- We are the first group to apply RRMHD to the study of QGP



Electromagnetic dissipation for QGP photons

Electromagnetic fields inside QGP

- EM fields penetrating QGP drive charge carriers out-of-equilibrium

$$J^\mu = qu^\mu + \sigma F^{\mu\nu} u_\nu$$

EM current in the QGP medium

First order dissipation from the EM fields

- Taking the Boltzmann equation in the relaxation time approximation focus on 2→2 processes,

$$k^\mu \partial_\mu f_a + eQ_a F^{\mu\nu} k_\mu \frac{\partial f_a}{\partial k^\nu} = -\frac{k^\mu u_\mu}{\tau_R} \delta f_{a,EM}^{(n)}$$

Vlasov term for the external EM fields

Order “n” corrections to the quark distribution function because of dissipation from the EM fields

J. A. Sun and L. Yan
Phys. Rev. C 109, 034917 (2024)

Electromagnetic dissipation for QGP photons

Electromagnetic fields inside QGP

- For this calculation we focus on 1st order corrections,

$$k^\mu \partial_\mu f_a + eQ_a F^{\mu\nu} k_\mu \frac{\partial f_a}{\partial k^\nu} = -\frac{k^\mu u_\mu}{\tau_R} \delta f_{a,EM}^{(n)}$$

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$$f_a = f_{a,eq} + \delta f_{a,EM}^{(1)} + \delta f_{a,EM}^{(2)} + \delta f_{a,EM}^{(3)} + \dots$$

Ordered by the EM field strength

$$\delta f_{a,EM}^{(1)}(X, k) = -\frac{-f_{a,eq}(1 - f_{a,eq})}{T\chi_{el}k^\mu u_\mu} e\sigma Q_a e^\mu k_\mu$$

Electric conductivity of QGP from Landau matching with the current

$$J^\mu = qu^\mu + \sigma F^{\mu\nu} u_\nu$$

EM fields in the fluid rest frame

$$e^\mu = (\gamma v_k E^k, \gamma E^i + \gamma \epsilon^{ijk} v_j B_k)$$

Electromagnetic dissipation for QGP photons

Electromagnetic fields inside QGP

- What we do is calculate the fluid + EM field contributions using hydrodynamics

Temperature and four-velocity

Electric conductivity value

Values that come from a hydrodynamic calculation

$$\delta f_{a,EM}^{(1)}(X, k) = - \frac{-f_{a,eq}(1 - f_{a,eq})}{T \chi_{el} k^\mu u_\mu} e \sigma Q_a e^\mu k_\mu$$

Electric susceptibility of QGP

$e^\mu = (\gamma v_k E^k, \gamma E^i + \gamma \epsilon^{ijk} v_j B_k)$
Spacetime dependent EM fields in QGP medium

$$\chi_{a,el} = - \frac{1}{3} \int \frac{d\vec{p}}{(2\pi)^3 E_p} (p^\sigma p^\nu \Delta_{\sigma\nu}) \frac{-f_{a,eq}(1 - f_{a,eq})}{p^\mu u_\mu}$$

Electromagnetic dissipation for QGP photons

Electromagnetic fields inside QGP

- What we do is calculate the fluid + EM field contributions using hydrodynamics
- All of those values can be calculated self-consistently using relativistic resistive magneto-hydrodynamics (RRMHD)

Temperature and four-velocity

Electric conductivity value

Values that come from a hydrodynamic calculation

$$\delta f_{a,EM}^{(1)}(X, k) = - \frac{-f_{a,eq}(1 - f_{a,eq})}{T \chi_{el} k^\mu u_\mu} e \sigma Q_a e^\mu k_\mu$$

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Thermal Photons from quark-gluon plasma

- Enhancement of the photon elliptic flow when the EM fields are included

