



Zs. 73. (3-4)

Rosenfeld: Bemerkung zur korrespondenz-  
mäßigen Behandlung des rel. Mehrkörper-  
problems. S. 253

Lanczos: zsl. Mg. als natürliche Eigen-  
schaft der Riemannsche Geometrie  
S. 147

25.74 78. S. 433

E. Regener

Über das Spektrum der Ultrastrahlung

I. Die Messungen im Herbst 1928.

25. 74. 7-8. S. 503

R. Schachenmeier, Wellenmechanische Vorstudien  
zu einer Theorie der Supraleitung

25. 74. 3-4 S. 143

F. London: Zur Theorie nicht adiabatisch  
verlaufender chemischer Prozesse.

# HET NEUTRON

door A. J. Rutgers

summary:

A survey is given of the literature about the neutron. At the end some remarks are made about nuclear constitution.

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4. Bothe. Phys. Z. 32, 692, 1931
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38, 1399, 1931
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1932,

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1929.

34. Barvart, Proc. R. Soc. 128, 345, 1930

35. Chau, Proc. Nat. Acad. 16, 431, 1930

36. Meitner and Hupfeld, Z. f. Ph. 67, 147,  
1932,

37. Jacobson, , 120, 145, 1932,

38. Rutgers, Nature 129, 361, 1932.

Weizsäcker? Ortmessung eines Elektrons  
durch ein Mikroskop.

25. 70 5,114 (eV)

grauher. el. dyn. 2073 Ortmessung,  
Genauigkeit  $\sim \frac{1}{2} \lambda$  (Heisenberg, 1927)

$$d \sim \frac{\lambda}{\sin \epsilon}$$

an  $\lambda$   $\sim \frac{h}{mc}$   $\sim 2,4 \times 10^{-12}$  m

$\lambda \sim \frac{h}{mc}$  ; order  $\sim 2,4 \times 10^{-12}$  m  
in  $\lambda$   $\sim \frac{h}{mc}$   $\sim 2,4 \times 10^{-12}$  m



# Jamou: mechanism of $\gamma$ -excitation by $\beta$ disintegration Nature 131, p. 57 (1935) (Jan 14)

neutron が proton へ  $\beta$  崩壊して electron を emit する  
過程では、upper level の neutron からの energy が  
electron へ  $\beta$  崩壊して proton  
へ excite されて  $\beta$  崩壊後の lower level へ遷移  
して hard  $\gamma$ -ray が emit される。

この  $\beta$  崩壊で proton number が 1 増加し、 $\beta$  崩壊で emit される  $\alpha$ -particle  
が emit される。emit される  $\gamma$ -ray は、その  $\beta$  崩壊で  
hard  $\gamma$  である。

また hard  $\gamma$ -ray が nucleus へ  $\beta$  崩壊。nuclear  
electron を excite して、その proton へ transition  
して characteristic  $\gamma$ -ray が emit される。  
(Gray, Parvack)

J. D. Cockcroft; E. T. S. Walton

Disintegration of light elements by Fast Protons

Jan. 7.

Nature, p. 23 (Vol 131) (1932)

Lithium:  $\alpha$ -particle range 8.4 cm, 2 cm  
( $\beta$ -ray  $\alpha$ -end  $\gamma$   $\alpha$   $\gamma$   $\alpha$ )

Boron: (one particle per two million protons  
at 500 Kilovolts)  $\beta$  to  $\gamma$ ,  $\alpha$   $\gamma$  intensive.

$\alpha$ -particle range 3 cm.

$\beta$ ''  $\alpha$   $\gamma$   $\alpha$ -particle 2 or 3 cm.

5 cm, range 1, 2, 2 + 2 cm.

For matrices  $\hat{H}_E$  electron & proton  $\hat{H}_P$   
 states.

→ 'sol'  $\hat{H}_P \hat{H}_E \psi = 0$ , sol "

$\hat{H}_P \psi = 0$ ,  $\hat{H}_E \psi = 0$  }  $\psi = \frac{1}{\sqrt{2}} (\phi_1 + \phi_2)$

-for }  $\hat{H}_E \psi = \phi$ ,  $(\hat{H}_P \phi = 0)$   
 $\hat{H}_P \psi = \phi$ ,  $(\hat{H}_E \phi = 0)$

,  $\phi_1 \neq \phi_2$  is not 0.

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京都大学基礎物理学研究所 湯川記念館史料室

J. Solomon. — Sur quelques difficultés de la  
théorie des quanta

J. de Physique <sup>Série VII</sup> <sup>(No 10)</sup>  
Tome II p 3 21

① L'analogie entre les équations de Maxwell et les  
équations de Dirac

1. Rosenfeld et Solomon (J. Phys. t. 2 (1931), p. 129) =  $\text{Zur}$

(électromagnétique) (le champ), Hamiltonien  $\mathcal{H}$ :

$$\begin{aligned} \vec{F} &= \sqrt{\Delta} \vec{F} \\ \vec{F}^\dagger &= -\sqrt{\Delta} \vec{F}^\dagger \end{aligned} \quad (1)$$

2. Équation de Hamilton avec courant de convection.  
hamiltonien du rayonnement pur

$$\mathcal{H}_R = \frac{1}{2} \vec{F}_k \vec{F}_k^\dagger$$

$$\left. \begin{aligned} \vec{F}_k + \text{rot} \left( \frac{\text{rot} \vec{F}}{\sqrt{\Delta}} \right) &= S \\ \vec{F}_k^\dagger - \text{rot} \left( \frac{\text{rot} \vec{F}^\dagger}{\sqrt{\Delta}} \right) &= S \end{aligned} \right\} \quad (4)$$

$$\text{div} \vec{F} = \text{div} \vec{F}^\dagger = \rho \quad (5)$$

3. Application de la théorie des (spinors).

(I) L'énergie électrostatique propre de l'électron.

1. généralisés

2. L'énergie électrostatique propre de l'électron dans  
la théorie quantique relativiste des champs.

3. L'énergie électrostatique propre et la nouvelle  
théorie du champ

$\approx \frac{1}{2} \int (\mathcal{H})$  infinie

(III) L'équation de Dirac et les états d'énergie  
négative

Nature No 3237 Supplement, p 825, 1931  
 Jeans: Beyond the Milky Way. vol 128

The galactic system  
 intra-galactic systems.  
 The Breaking-up of the Universe  
 Three agencies: (1) rotation (2) tidal action  
 (3) gravitation instability  
 The scattering of the Universe

1. négative et états d'énergie négative ~ passage à  
 potentiel  $\sim \frac{h}{4\pi m_0 c}$ , ordre, intervalle  $\sim \frac{m_0 c^2}{eV}$   $\sim$   
 $\sim$   $\frac{h}{4\pi m_0 c}$   $\sim$   $\frac{h}{4\pi m_0 c}$  (Nikolsky, ZS, 62, (1930) 677)  
 (Sauter, ZS, 69, (1931), 742)

2. première orbite de l'hydrogène  $\sim \dots$   
 $\int_0^{\frac{h}{4\pi m_0 c}} |\psi|^2 dr \sim 4\pi^2 \left(\frac{h}{4\pi m_0 c}\right)^3 |\psi(0)|^2 = \pi^2 \left(\frac{2\pi e^2}{hc}\right)^3$   
 $= \pi^2 \alpha^3$ .

particle  $\sim \frac{h}{4\pi m_0 c}$   $\sim$   $\frac{h}{4\pi m_0 c}$   $\sim$   $\frac{h}{4\pi m_0 c}$   $\sim$   $\frac{h}{4\pi m_0 c}$   
 P330  
 $\sim$  longueur minimale  $\sim \frac{h}{m_0 c}$   $\sim$  introduire  $\sim$   
 $\sim$   $\frac{h}{m_0 c}$   $\sim$   $\frac{h}{m_0 c}$   $\sim$   $\frac{h}{m_0 c}$   $\sim$   $\frac{h}{m_0 c}$   
 Lorentz transf  $\rightarrow$

Born u. Kummer: Ansätze zur Q.E.D.

DS. 69, 141, 1951,

§ Einleitung

§ 1. Vermeidung der Singularität beim Coulombschen Potential.

$$\frac{e \cdot \kappa r}{r} = \int \frac{4\pi e^{2\pi i(p_0 r)}}{4\pi^2 p^2 - \kappa^2} dp, \quad \kappa > 0 \quad (1)$$

$$\frac{1}{r} \rightarrow \frac{1}{\pi p^2} \quad (2)$$

$$\varphi(r) = \frac{e}{\pi} \int_{p < p_0} \frac{e^{2\pi i(p_0 r)}}{p^2} dp = \frac{e}{\pi} \frac{2}{\pi} \int_0^{2\pi p_0 r} \frac{\sin x}{x} dx, \quad (4)$$

$$r \rightarrow \infty: \varphi(r) \rightarrow \frac{e}{r}, \quad (5)$$

$$r \ll \frac{1}{p_0}: \varphi(r) = 4e p_0 \left(1 - \frac{2}{\pi} \alpha^2 p_0^2 r^2 + \dots\right) \quad (6)$$

Die Anzahl der ein Elektron aufbauenden Wellen ist endlich; Wellen mit einer Wellenlänge kürzer als eine bestimmte Länge  $r_0 \approx 1/p_0$  sind selten, und zwar gemäß dem Gaußschen Fehlergesetz\*:

$$(7) \quad g(p_0) = g_0 e^{-\alpha \left(\frac{p_0}{p_0}\right)^2} \int g(p_0) dp_0 = g_0 p_0^3,$$

$$(8) \quad \varphi(r) = \frac{g_0}{\alpha} \int \frac{e^{-\alpha \left(\frac{p_0}{p_0}\right)^2}}{p^2} e^{2\pi i(p_0 r)} dp,$$

$$\Delta \varphi = -4\pi e \int g(p_0) e^{2\pi i(p_0 r)} dp,$$

$$\varphi(r) = \frac{\epsilon}{r} \frac{2}{\sqrt{\pi}} \int_0^{\frac{r}{r_0}} e^{-y^2} dy \quad (14)$$

$$r \rightarrow \infty: \varphi(r) \rightarrow \frac{\epsilon}{r} \quad (15)$$

$$r \ll r_0: \varphi(r) = \frac{2\epsilon}{r_0} \left(1 - 2\left(\frac{r}{r_0}\right)^2 + \dots\right) \quad (16)$$

$$r_0 = \frac{1}{\beta_0}$$

elektrostat. Feldenergie

$$U_0 = \frac{1}{8\pi} \int (\text{grad } \varphi)^2 dS = \frac{1}{2} \int \rho \varphi dS$$

$$U_0 = \frac{1}{\sqrt{2}} \frac{\epsilon^2}{r_0}$$

## § 2. Die Nullpunktsenergie des Hohlraums

Anzahl der Eigenschaften

$$dz = \frac{8\pi}{c^3} \nu^2 d\nu$$

1 kan  $\nu =$

$$(23) \quad dz = 2 g(\nu) dP = g_0 \frac{8\pi}{c^3} \nu^2 e^{-\frac{\pi(\nu)^2}{\lambda_0^2}} d\nu,$$

statt Null.  $N, P, E, \dots$

$$E_0 = \frac{h}{2} 2 \int_0^{\infty} \nu g(\nu) dP = h\nu_0 \frac{2\pi^3}{\lambda_0^3}$$

$$= h\nu_0 \pi^2 z. \quad (25)$$

$$z = 2 \int g(\nu) dP = \frac{2}{\lambda_0^3} \quad (24)$$

Anzahl der Lichtquanten  $h\nu_0$  im Grundzustand

$$n_0 = \frac{E_0}{h\nu_0} = \pi^3 z.$$

§3. Beziehung zur Q.M.

Dichte  $f(x) = \int g(p_0) e^{i\pi(p_0 x)} dP$  (19)  $\frac{h\nu_0}{2} \sim \frac{z}{2}$

$$f(x-x') = \int \psi(x, p_0) \psi(x', p_0) dP, \quad (28)$$

$$\psi(x, p_0) = e^{-\frac{\pi}{2} \left(\frac{p_0}{p_0}\right)^2} e^{i\pi(p_0 x)} \quad (27)$$

(28) " (27), Freisystem  $\rightarrow \hat{H}^2 \sim$  Einheit operator  $\rightarrow \hat{H}^2 \sim$

$\rightarrow$  orthogonal  $\rightarrow$  "normieren"  $\rightarrow$   $\int \psi(x, p_0) \psi(x', p_0) dx = \delta(p_0 - p_0')$

$$g(p_0) = e^{-\pi \left(\frac{p_0}{p_0}\right)^2} \int g(p_0) dP = p_0^3 \quad (29)$$

mangeregter Hohlraum  $\rightarrow$  ruhenden Teilchen  $\rightarrow$   $\frac{h\nu_0}{2} \sim z$

welle  $\rightarrow$  Nullwelle  $\rightarrow$   $\frac{h\nu_0}{2} \sim z$ , Verteilung: Experiment  $\rightarrow$  無關係.

§4. Vereinigung von Teilchen

$$F(x-x') = A e^{-\alpha(x-x')^2} \quad G(x-x') = B e^{-\beta(x-x')^2}$$

$\rightarrow$  Matrixprodukt  $\rightarrow$   $\rightarrow$   $\rightarrow$

$$H(x-x') = \int F(x-x') G(x''-x') ds''$$

$$= AB \int e^{-\alpha(x-x'')^2 - \beta(x''-x')^2} ds''$$

$$= C e^{-\gamma(x-x')^2}$$

$$\frac{1}{\gamma} = \frac{1}{\alpha} + \frac{1}{\beta}$$

$$A = \left(\frac{\alpha}{\pi}\right)^{1/2}, \quad B = \left(\frac{\beta}{\pi}\right)^{1/2} \rightarrow C = \left(\frac{\gamma}{\pi}\right)^{1/2}$$

Vereinigte Teilchen, Radius  $\rightarrow$

$$\lambda_0^2 = r_1^2 + r_2^2$$

$$\lambda_0 = \frac{c}{\nu_0} r_2 \quad \lambda = 2\lambda_0 \quad \text{ Bohr}$$

$$h\nu = h \cdot 2 \frac{c}{\lambda_0} = \frac{e^2}{\sqrt{2} r_1}$$

$$\lambda_0 = r_2$$

$$2) \quad \frac{r_1}{r_2} = 2 \frac{e^2}{\sqrt{2} h c} = 2 \frac{\alpha}{2\pi\sqrt{2}}$$

$$\alpha = \frac{2\pi e^2}{h c} = \frac{1}{136}$$

$$\frac{m_1}{m_2} = \frac{r_2}{r_1} = 1210 \approx$$

$\approx 21,52$   $\times \times \times \times$   $\approx 21,52$   $\times \times \times \times$   $1858$   $\times \times \times \times$



(Proc. Roy. Soc. 117, 1928)

p. 625 Flint & Fischer: The fundamental eq. of wave mechanics & the metrics of space

p. 630 Flint: Relativity and the Quantum Theory

p. 637. Flint & Richardson: On a minimum proper time and its applications: (1) to the number of the chemical elements (2) to some uncertainty relations.

rest mass  $m$  +  $x$  particle, momentum,

$$p_\alpha = m g_{\alpha\beta} \frac{dx^\beta}{ds}, \quad (1)$$

$$p^\alpha (= m \frac{dx^\alpha}{ds})$$

$$p_\alpha p^\alpha = m^2$$

momentum = associate  $\Rightarrow$  action  $W$

$$p_\alpha = \frac{\partial W}{\partial x^\alpha}$$

$$\therefore g_{\alpha\beta} \frac{\partial W}{\partial x^\alpha} \frac{\partial W}{\partial x^\beta} = m^2,$$

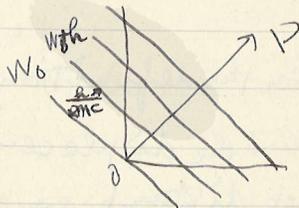
$$|\text{grad } W|^2 = m^2.$$

$$\frac{\partial W}{\partial x^\alpha} \frac{dx^\alpha}{ds} = m,$$

$\nabla_\alpha \frac{dx^\alpha}{ds} \approx \text{div. on } x^\alpha \sim \text{world line}$ ,  $\frac{\partial W}{\partial x^\alpha} = \text{Lagrange } p_\alpha$

$$\frac{dW}{ds} = m$$

$$W = ms + W_0$$



vol 183.

G. P. Thomson: Diffraction of Electrons by  
Single crystals p. 2.

Mott: The Theory of the effect of Resonance  
Levels on Artificial Disintegration p. 228

Chandrasekhar: Stellar Coeffs. of Absorption  
and Opacity p. 241

Kronig: Q. T. of Dispersion of in Metallic  
Conductors. II. p. 255

Eddington: The Property of wave Functions  
p. 311

Hulme: Photoelectric effect for  $\gamma$  rays.  
p. 381

Wilson: Theory of Electronic Semi-Conductors

Eddington: On the value of Cosmical Const.  
p. 605

Arnott: Diffraction of Electrons in Gases.  
p. 615

On the definition of distance in curved space,  
and the displacement of the spherical lines  
of distant sources. by E. T. Whittaker. p. 106.

$$\begin{aligned} m \cdot c \Delta X' &= \frac{m \cdot g \cdot x}{c} - p_x \cot \theta + \Delta E \frac{\Delta x}{c} \\ &\approx + \Delta E \frac{\Delta x}{c} - p_x \cot \theta \approx p_x \cot \theta \\ &= E \frac{h}{c} \end{aligned}$$

L. Rosenfeld  
Bemerkung zur Korresp. Behand. des rel. Mehrkörper-  
problems 25. 73, S 253. (~~3~~<sup>4</sup>-6)

Lanczos  
El. Mg. als natürlicher Eigenschaft der Riemann-  
schen Geometrie 25 73, S 147 (3-4)

Karl Jaspers: Philosophie  
I, Philosophische Weltorientierung 10.60  
II, Existenzzerhellung 13.20  
III, Metaphysik 8.40  
(Springer)

Plexner, Die Universität in Amerika,  
England, Deutschland, 1982.  
RM 19.60

25. 73 7-8 }  
Ann. 13 1. } Jan 27th  
Proc. Roy. Soc. 823 }  
Arch.

Saitoh:  
kleinsche Paradoxien

Proc. Roy. Soc. vol 134,

No 823. p. 154. Riezler: The Scattering of  $\alpha$ -Particles  
 by Light Elements.

p. 103. Taylor: Interact. Energy of 2  $\alpha$ -Particle:  
 at close distances, determined from the  
 anomalous scattering in Helium,

~~No 823~~ p. 137. Gurney: Quantum Mechanics of Electrolysis

No. 824. p. 524 Eddington: On the Mass of the Proton

$$\left\{ 10 \left( i E_s \frac{\partial}{\partial \theta_s} \right)^2 + 136 \left( i E_s \frac{\partial}{\partial \theta_s} \right) + 1 \right\} \psi = 0$$

or  $(135.9264 i E_s \frac{\partial}{\partial \theta_s} + 1) (0.0735692 i E_s \frac{\partial}{\partial \theta_s} + 1) \psi = 0$

$$\frac{135.9264}{0.0735692} = 1847.68 \quad (\text{obs value } 1846.6 \pm 0.5)$$

Boss; Phil. Mag., 12 p. 632 (1911)

$$\frac{2\pi m c \alpha}{h} = \frac{\sqrt{N}}{R}$$

$$d^2 \left( i E_s \frac{\partial}{\partial s} + \frac{2\pi m c}{h} \right) \psi = 0$$

$$ds = R d\chi_s = (R/\sqrt{N}) d\theta_s$$

$$\left\{ \alpha \left( i E_s \frac{\partial}{\partial \theta_s} \right) + 1 \right\} \psi = 0.$$

root sum ..  $1 \rightarrow 1 \uparrow \uparrow + 1 \uparrow \uparrow \uparrow \uparrow \uparrow \uparrow$

209 field  $\rightarrow$  ..  $\rightarrow$  energy  $e\phi$   $\rightarrow$   $\rightarrow$   $\rightarrow$   
 .. energy ..  $e\phi$  ..

Rasetti : Über die Anregung von  
Neutronen in Beryllium

78, 165  
(1932)

### § Die Eigenschaften des Neutrons

$$\text{neutron } \sigma_f \sim \frac{e^2}{mc^2}$$

$$\sigma_p \sim \frac{hc}{2\pi e^2} \cdot mc$$

Massendefekt

$$\bar{E} = c \cdot \sigma_p \approx 137 mc^2$$

Neutronen Kerns Massendefekt  $\approx \frac{1}{137} \cdot 100 \text{ MeV}$

Neutron = 4e- Elektron, Bindungsenergie

Massendefekt  $\approx 2 \text{ MeV} mc^2$ , order  $10^{-8}$

Neutron,  $h \cdot 5400 \text{ s}^{-1} \approx 137 mc^2$ , order  $10^{-7}$

3/4  $10^{-8}$

Neutronen  $\alpha$  Strahlung, Streuung

kl. Osz. der Frequenz  $\nu_0 + 2 \text{ MeV} \approx 4.8 \cdot 10^{11} \text{ s}^{-1}$

Wirkungsquerschnitt

$$\sigma = \frac{8\pi}{3} \left(\frac{e^2}{mc^2}\right)^2 \left(\frac{\nu^2}{\nu_0^2 - \nu^2}\right)^2$$

z.B.

$$h\nu = 5.15 \text{ MeV} : \sigma = 1.5 \cdot 10^{-28} \text{ cm}^2$$

$$\therefore h\nu_0 = 42.6 \text{ MeV}$$

von Nr Elektronen, Bindungsenergie

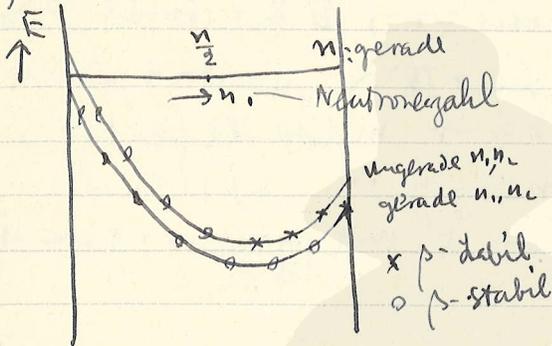
order  $10^{-7}$  zu Neutron,  $h \cdot 5400 \text{ s}^{-1}$

2  $10^{-8} / \nu - 3.8 \cdot 10^{-8}$

Über den Bau der Atomkerne, II.

Heisenberg Zs. 78, S. 156, 1932.

§ Stabilität der Kerne



Beck 47 / 407  
 50 578

Pauli Prinzip für N und P im Kern

Max und Min von  $n_1$  (für jede  $n$ ) ~~gerade~~

'' ''  $n_2$  ( ''  $n_2$  ) ''

(leichtes Element,  $1 \leq n \leq 10$ )

$n_1, n_2$ : Anzahl Kernprotonen und Neutronen

§ Streuung von  $\gamma$ -Strahlen am Atomkern

Rayleigh Streuung:  $\sigma_R \approx \sigma_N \cdot n_1^2$

Raman ''  $\sigma_R \approx \sigma_N \cdot n_1$

(Meitner, Huffeld 15, 705)









7カ>δ。

Long Range /  $\alpha, \beta$ , H particle reenergy + range +  
1. 倫理... definite + (2M) = 4...  
子 + 2 + ...

Williams: The Passage of  $\alpha$ - and  $\beta$ -Particles  
through Matter and Bohr's theory of Collisions  
(Proc. Roy. Soc. B35, 108, 1932)

= 2M

Bethe (Ann. Physik 5, 325, 1930)

, non-relativistic theory of collision //

Gaunt (Proc. Camb. Phil. Soc. 23, 732, 1927)

, 理論の基礎 = stopping power + ...

$\beta$ -particle, velocity,  $v \approx 0.98c$

39% ... systematic  
ment) + relativity correction (ca. 20%) \*  
velocity + ...

(total ionisation in non atomic gas,  
straggling of  $\alpha$ -particle in light element gas)

1. 倫理... Rel. Cor. // Williams: Proc. Roy. Soc. B30, 310,  
328

9用... order = ...



→ atomic & molecular spectra, hyperfine  
structure, nucleus, moment  $\rightarrow \vec{I} + \vec{S} + \vec{M}$   
核の構造  $\rightarrow$   $^{12}\text{O}$  nuclear structure, spin  $\rightarrow$   $\vec{I}$   
核の構造  $\rightarrow$   $^{12}\text{O}$

2. Note: 目的は最近の文献をとり Nuclear Physics  
= 同様の性質、この初期の核物理の出来事  
systematic + 規則性、又暗示的 + 2つの理論  
2つ + 2つ、理論 + 観測 = 核物理  
核子 =  $p, n$

理論 = 核子  
先験的核子 + 行成る核子 = nucleus  
核子  $\alpha$ -particle, proton の energy  
level 有 = 行成る核子 + 核子の核子  
 $p, n$

1.  $\gamma$ -ray spectra (連続的)

2. Spontaneous disintegration <sup>核子</sup>  $\beta$ -ray,  
continuous spectra, 連続的 +  $\beta$  核子の核子  
 $p, n$

3. Artificial Transformation = 核子 Proton  
"range" of exciting  $\alpha$ -particle, range of  
核子の nucleus, 核子, nucleus  $p$ , proton  
of  $\alpha$ -particle energy level + 核子の核子  
 $p, n$