

H. Lehmann, über Eigenschaften von $h_{\mu\nu}(1)$
 Ausbreitungsges. und Renormierungs-
 konstanten quantisierter Felder
 (Nuovo Cimento XI (1954),
 (1^o Aprile), 342)

a) Skalare Felder:

$$\left. \begin{aligned} \frac{\partial A}{\partial x_\mu} &= i[A(x), P_\mu] \\ [P_\mu, P_\nu] &= 0 \end{aligned} \right\}$$

$$P_\mu \Phi_k = k_\mu \Phi_k \quad (k_0 \geq 0)$$

$$\langle \Phi_0, A(x) A(x') \Phi_0 \rangle = \langle A(x) A(x') \rangle_0 \\ = i \Delta^{(+)\prime}(x-x')$$

$$\langle A(x') A(x) \rangle_0 = -i \Delta^{(-)\prime}(x-x')$$

$$\langle [A(x), A(x')] \rangle_0 = i \Delta'(x-x')$$

$$= -2i \varepsilon(x_0 - x'_0) \bar{\Delta}'(x-x')$$

$$\langle \{A(x), A(x')\} \rangle_0 = \Delta^{(0)\prime}(x-x')$$

$$\langle T A(x) A(x') \rangle_0 = \frac{1}{2} \Delta_F'(x-x')$$

$$\Delta_F'(x) = 2i [\theta(x_0) \Delta^{(+)\prime}(x) - \theta(-x_0) \Delta^{(-)\prime}(x)]$$

$$= \Delta^{(0)\prime}(x) - 2i \bar{\Delta}'(x)$$

$$\langle A(x) A(x') \rangle_0 = \sum_k (\Phi_0, A(x) \Phi_k) (\Phi_k, A(x') \Phi_0)$$

$$= \sum_k A_{0k}(x) A_{k0}(x')$$

$$= \sum a_{0k} a_{0k}^* \exp[ik(x-x')] \quad \text{No. 5 1953. 12. 5.000}$$

L. (2)

$$(\Phi_0, A(x)\Phi_k) = A_0k(x) \\ = a_0k \exp(ikx)$$

$$\left(\because \frac{\partial A_0k(x)}{\partial x^\mu} = i (\Phi_0, [A(x), \gamma_\mu] \Phi_k) \right. \\ \left. = i k_\mu A_0k(x) \right)$$

$$P(-k^2) = (2\pi)^3 \sum_{\mathbf{k}} a_0k a_0k^* \quad (\text{isotropy of vacuum})$$

$$\Delta^{(0)'}(x) = \int_0^\infty \Delta^{(0)}(x; \kappa^2) \rho(\kappa^2) d(\kappa^2)$$

$$\rho(-k^2) = \int_0^\infty \rho(\kappa^2) \delta(-k^2 + \kappa^2) d(\kappa^2)$$

$$\rho(\kappa^2) \geq 0$$

$$\Delta'(x) = \int_0^\infty \Delta(x; \kappa^2) \rho(\kappa^2) d(\kappa^2)$$

$$\langle [A(x, t), A(x', t)] \rangle_0 = 0$$

$$\langle [A(x, t), \dot{A}(x', t)] \rangle_0 = -i \delta(x - x') \\ \times \int_0^\infty \rho(\kappa^2) d(\kappa^2)$$

b) Spinor Fields:

$$\langle [\psi_\alpha(x, t), \bar{\psi}_\beta(x', t)] \rangle_0 = \gamma_{\alpha\beta}^4 \delta(x - x') \\ \times \int_0^\infty \rho(\kappa^2) d(\kappa^2)$$

$$\begin{aligned}
 (\Phi_{k_j} A(x) A(x') \mathbb{I}_L) &= \sum_j A_{k_j}(x) A_{j_l}(x') \quad L, (3) \\
 &= \sum_j a_{k_j} a_{j_l}^* \exp i(k_{n_j} x - k_{e_j} x') \\
 k_{n_j} &= k_{j_l} - k_{k_l}
 \end{aligned}$$

$$\begin{aligned}
 (\Phi_{k_l} A(x) A(x') \mathbb{I}_k) &= \sum_j a_{k_j} a_{k_j}^* \exp i k_{k_j} (x-x') \\
 &= \exp -i(k_{k_l} - k_0)(x-x') \\
 &\quad \times \sum_j a_{k_j} a_{k_j}^* \exp i k_{k_j} (x-x')
 \end{aligned}$$

(possible anisotropy of the state \mathbb{I}_k)

Linear theory:

$$a_{k_0} = a_{k_0}^* \neq 0 \quad \text{only for } k: \text{ one particle state}$$

$$a_{k_l} = a_{k_l}^* \neq 0 \quad \text{only for the case } n_{k_l} + 1 = n_l$$

$$\text{ant } a_{k_l} a_{k_l}^* = n_l$$

$$a_{k_l}^* a_{k_l} = n_{k_l}$$

$$a_{k_l} a_{k_j}^* = a_{k_j}^* a_{k_l}$$

$$\{k_\mu k_\nu + \kappa^2\} a_{\mu\nu} \Phi(\{n_{k_\mu}\}) = 0$$

$$\Phi = P_\kappa \Phi$$

$$\Phi = \begin{pmatrix} \Phi(0, \kappa, n_{\mu\nu}, 0) \\ 0 \\ 0 \\ 0 \end{pmatrix}$$

$$\int_{\mathcal{D}^4} d^4x e^{ik_\mu x_\mu} \{k_\mu k_\nu + \kappa^2\} a_{\mu\nu} \Phi = 0 \quad \downarrow$$

$$\{-\frac{\partial^2}{\partial x_\mu \partial x_\mu} + \kappa^2\} \varphi \cdot \Phi = 0$$

$$\left. \begin{aligned} \Phi' &= a_k^* \Phi \\ \Phi'' &= \varphi^* \Phi \end{aligned} \right\} \text{は } \Phi\text{-空間の外にあり.}$$

$$\int_{\mathcal{D}^4} d^4x e^{-ik_\mu x_\mu} \{-\frac{\partial^2}{\partial x_\mu \partial x_\mu} + \kappa^2\} \varphi(x) \Phi = 0$$