

Angular Momenta in the Theory of Hyperfine Structure  
 The theory of the hyperfine structure of one electron spectra has recently been treated by means of the relativistic equation of Dirac. The writer wishes to show that in this case the angular momenta of electron alone are no more constant, but the total angular momenta including the nuclear spin angular momenta still conserve. The proof is very simple. According to the ordinary notation, the Hamiltonian is written in the form

$$H = -eV(r) + c\alpha(p + \frac{e}{c}A) + \alpha_4 mc^2,$$

where  $A = \frac{[M \times r]}{r^3}$

represents the field caused by the nuclear magnetic moment  $\mu$ . The components of  $M$  are assumed to take eigenvalues  $\mu, \mu \frac{k-1}{k}, \mu \frac{k-2}{k}, \dots, -\mu \frac{k-1}{k}, -\mu$

and to satisfy the relations etc.

$$\mu_y \mu_z - \mu_z \mu_y = i \frac{\mu}{k} \mu_x$$

$$\text{or } [M \times M] = i \frac{\mu}{k} M,$$

but, commute with other variables, and so take eigenvalues

$$\mu_1, \mu \frac{k-1}{k}, \mu \frac{k-2}{k}, \dots, -\mu \frac{k-1}{k}, -\mu. \quad \alpha = \beta, \alpha_4 = \beta_3.$$

$\alpha$  and  $\alpha_4$  can also be expressed by Dirac's matrices  $\beta, \alpha, \alpha_4$  as follows. In this case, the angular momenta of the electron

$$[r \times p] + \frac{\hbar}{4\pi c} \sigma$$

commute with the terms of the Hamiltonian, except the nuclear spin term, which does not commute, for which the relation

$$\left\{ [r \times p] + \frac{\hbar}{4\pi c} \sigma, \frac{\sigma [M \times r]}{r^3} \right\} = \left\{ -e\beta, \frac{\sigma [M \times r]}{r^3} \right\} \left\{ [r \times p] + \frac{\hbar}{4\pi c} \sigma \right\}$$

$$= -e\beta \cdot \frac{\hbar}{2\pi i} \left\{ [r \times (\sigma \times M)] + \sigma \times [M \times r] \right\}$$

is established through the following calculation:

$$\sigma [r \times p] (r \times M) - (r \times (\sigma \times M)) [r \times p] = \frac{\hbar}{2\pi i} [r \times (\sigma \times M)]$$

$$\text{or } \{ y p_z - z p_y \} i x (\sigma_y \mu_z - \sigma_z \mu_y) + y (\sigma_z \mu_x - \sigma_x \mu_z) + z (\sigma_x \mu_y - \sigma_y \mu_x) + \{ y p_z - z p_y \}$$

$$- \{ x (\sigma_y \mu_z - \sigma_z \mu_y) + y (\sigma_z \mu_x - \sigma_x \mu_z) + z (\sigma_x \mu_y - \sigma_y \mu_x) \} = \frac{\hbar}{2\pi i} [r \times (\sigma \times M)]_x$$

$$= \frac{\hbar}{2\pi i} \{ y (\sigma_x \mu_y - \sigma_y \mu_x) - z (\sigma_z \mu_x - \sigma_x \mu_z) \}$$

$$\text{by using the principle } p_x x - x p_x = \frac{\hbar}{i} \text{ etc. } = \frac{\hbar}{2\pi i} [ \sigma \times (M \times r) ]$$

$$\text{or } \frac{\hbar}{4\pi c} \sigma_x \{ \sigma_x (\mu_y z - \mu_z y) + \sigma_y (\mu_z x - \mu_x z) + \sigma_z (\mu_x y - \mu_y x) \}$$

$$- \{ \sigma_x (\mu_y z - \mu_z y) + \sigma_y (\mu_z x - \mu_x z) + \sigma_z (\mu_x y - \mu_y x) \} \frac{\hbar}{4\pi c} \sigma_x$$

$$= \frac{\hbar}{2\pi i} \{ \sigma_y (\dot{m}_x y - \dot{m}_y x) - \sigma_z (\dot{m}_y x - \dot{m}_x y) \} = \frac{\hbar}{2\pi i} [ \dot{\Phi} \times (\mu \times r) ]_x$$

by using the formulae relations

$$\sigma_x \sigma_y - \sigma_y \sigma_x = 2i \sigma_z \quad \text{etc.}$$

Accordingly,  $\frac{\hbar}{4\pi} \dot{\Phi} \{ r \times p \} + \frac{\hbar}{4\pi} \dot{\Phi}$  is not constant and  $\frac{\hbar}{2\pi i} \{ \dot{\Phi} \times (\mu \times r) \}$  is of motion.

$$(1) \left( \left\{ \left[ r \times p \right] + \frac{\hbar}{4\pi} \dot{\Phi} \right\} H - H \left\{ \left[ r \times p \right] + \frac{\hbar}{4\pi} \dot{\Phi} \right\} \right) = -e \phi \frac{\hbar}{2\pi i} \frac{\{ \dot{\Phi} \times (\mu \times r) \}}{r^3} \neq 0$$

Now let us consider the quantity

$$\frac{\hbar}{2\pi} \mu$$

which commutes with the terms of the Hamiltonian other than the nuclear spin term, for which

$$\frac{\hbar}{2\pi} \mu \left\{ -e \phi \frac{\Phi \{ \mu \times r \}}{r^3} \right\} - \left\{ -e \phi \frac{\Phi \{ \mu \times r \}}{r^3} \right\} \frac{\hbar}{2\pi} \mu = -e \phi \frac{\hbar}{2\pi i} \frac{\mu \{ \mu \times (r \times \mu) \}}{r^3}$$

is established by means of the formula

$$\begin{aligned} & \frac{\hbar}{2\pi} \mu_x \{ \mu_x (y \sigma_z - z \sigma_y) + \mu_y (z \sigma_x - x \sigma_z) + \mu_z (x \sigma_y - y \sigma_x) \} \\ & - \left\{ \mu_x (y \sigma_z - z \sigma_y) + \mu_y (z \sigma_x - x \sigma_z) + \mu_z (x \sigma_y - y \sigma_x) \right\} \frac{\hbar}{2\pi} \mu_x \\ & = \frac{\hbar}{2\pi} (-i) \frac{\mu}{r} \{ \mu_y [ r \times \Phi ]_z - \mu_z [ r \times \Phi ]_y \} \\ & = \frac{\hbar}{2\pi i} \frac{\mu}{r} [ \mu \times (r \times \Phi) ]_x, \end{aligned}$$

so that

$$(2) \quad \frac{\hbar}{2\pi} \frac{\mu}{r} \mu H - H \frac{\hbar}{2\pi} \frac{\mu}{r} \mu = -e \phi \frac{\hbar}{2\pi i} \frac{[ \mu \times (r \times \Phi) ]}{r^3}$$

From these two formulae (1) and (2), using the identities

$$\begin{aligned} & [ r \times [ \dot{\Phi} \times \mu ] ] + [ \dot{\Phi} \times [ \mu \times r ] ] + [ \mu \times [ r \times \dot{\Phi} ] ] \\ & = \dot{\Phi} (r \mu) + \mu (r \dot{\Phi}) + \mu (\dot{\Phi} r) - r (\dot{\Phi} \mu) + r (\mu \dot{\Phi}) - \dot{\Phi} (\mu r) = 0 \end{aligned}$$

The required formula

$$\frac{\hbar}{4\pi} \dot{\Phi} [ r \times p ] + \frac{\hbar}{4\pi} \dot{\Phi} + \frac{\hbar}{2\pi} \frac{\mu}{r} \mu H - H \left( \mu + \frac{\hbar}{4\pi} \dot{\Phi} + \frac{\hbar}{2\pi} \frac{\mu}{r} \mu \right) = 0$$

is deduced. So the sum of the angular momenta of the electron and the nuclear spin angular momenta are constants of motion, of course, the orbital angular momentum of the nucleus does not appear in this case as it is fixed in the centre, but still the spin angular momenta play their rôle in the expression of the total angular momenta.