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16 + 2B

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(1)  $e^{ikz}$  / 場合



$\vec{k}$  方向  $\sim$  plane wave  $\sim e^{i(kz)}$

$$R = (R_{in} \cos \beta, R_{in} \sin \beta, R_{out})$$

$$R = (r \sin \theta \cos \phi, r \sin \theta \sin \phi, r \cos \theta)$$

$$R_{out} = Rr (\sin \theta \sin \alpha \cos(\phi - \beta) + \cos \theta \cos \alpha)$$

$\phi = A(\alpha, \beta) + \pi$  for  $\alpha, \beta = \pi/2$  / 積分  $\int_0^{\pi/2} \dots$   $Z=0$  ( $\phi = \theta = \pi/2$ )

$\pi$  / 原点を中心した半径  $a$ , circular slit  $\gamma$  ( $\theta = \pi/2$ )

bundle  $= \pi/2$ .

$$y(r, \theta) = \int_0^{\pi/2} A(\alpha, \beta) e^{iRr (\sin \theta \sin \alpha \cos \beta + \cos \theta \cos \alpha)} d\alpha d\beta \quad \text{--- ①}$$

$\phi = (\phi - \beta)$  /  $\theta = \pi/2$   $\rightarrow$   $\phi = \pi/2 + \pi - \beta = \pi - \beta$

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~~Symmetry~~ Symmetry in  $\theta, \phi$  is  $A(\alpha, \beta)$  in  $\alpha, \beta$  function  $r, r'$

274  $A(\alpha)$   $r, r'$

D  $\beta = 2\pi$   $\int_0^{2\pi} \cos \theta \sin \theta d\theta$

$$\int_0^{2\pi} e^{i k r \sin \alpha \cos \beta} d\beta = \int_0^{2\pi} e^{i k r \sin \alpha \cos \beta} d\beta \quad \left[ \int_0^{2\pi} k r \sin \alpha \cos \beta d\beta \right]$$

$$= 2 \int_0^{\pi} e^{i k r \sin \alpha \cos \beta} d\beta$$

$$= 2\pi J_0(kr) = 2\pi J_0(kr \sin \alpha)$$

[  $\therefore$  Hansen's formula  $J_n(z) = \frac{1}{\pi} \int_0^{\pi} e^{iz \cos \theta} \cos n\theta d\theta$  ]

277  $r, r'$

$$\psi(r, \theta) = 2\pi \int_0^{\frac{\pi}{2}} A(\alpha) e^{i k r \sin \alpha \cos \theta} J_0(kr \sin \alpha) d\alpha$$

$\theta = \frac{\pi}{2}$   $r, r'$

$$\psi(r, \frac{\pi}{2}) = \begin{cases} 1 & r < a \\ 0 & r > a \end{cases}$$

278  $r, r'$

$$\psi(r, \frac{\pi}{2}) = 2\pi \int_0^{\frac{\pi}{2}} A(\alpha) J_0(kr \sin \alpha) d\alpha = \begin{cases} 1 & r < a \\ 0 & r > a \end{cases} \quad (2)$$





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$z = -d$ ,  $\rho = \frac{1}{2} \frac{d}{z}$ ,  $\rho = \frac{1}{2} \frac{d}{z}$ ,  $\rho = \frac{1}{2} \frac{d}{z}$

$$F(\theta) = \frac{R_0}{2a} \int \int f(\sin \theta \sin \alpha \cos(\phi - \beta) + \cos \theta \cos \alpha) T_1(R_0 \sin \alpha) \cos \alpha \, d\alpha \, d\beta$$

$\frac{1}{2} \frac{d}{z}$

$$I = |F(\theta)|^2$$

" scattering intensity  $\propto \frac{1}{z^2}$

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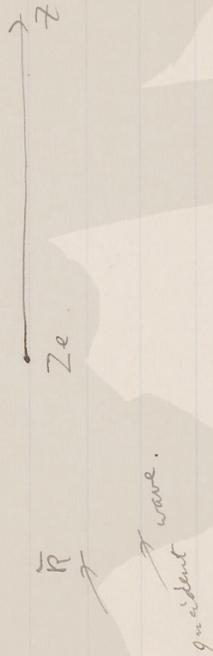
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(ii) Coulomb Field の場合

~~任意の有向波~~

Incident wave, 有向波  $R = |k|$ , wave function  $\psi$ ,  
 Nucleus の位置  $z = 0$  座標系  $R = |k|$  - parallel +  
 $|k| \neq |k'|$  散乱  $r = |R| \neq |k|$



第一項近似  $1 \ll r$

$$\psi \sim \left( 1 - \frac{Z^2}{i(kr - R^2)} \right) e^{i(kr - R^2)} + \frac{Z}{kr - R^2} e^{i(kr - R^2) \log(kr - R^2)} + i\alpha + 2i\eta_0 \quad (1)$$

$$\alpha = \frac{2\pi Z Z' e^2}{k v} = \frac{m Z Z' e^2}{\hbar^2 k}$$

$$e^{2i\eta_0} = \Gamma(1 + i\alpha) \Gamma(1 - i\alpha)$$

$m, Z, Z', k$ : incident particle, reduced mass  $B \ll \dots$  change  $\eta_0$



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$$\psi(R, \frac{R}{d}) = \int_0^{\frac{R}{d}} A(\alpha) e^{\alpha} \left[ (R \text{ and } \cos(\alpha \beta) - R \alpha \cos \alpha) + \log \left( \frac{R}{R \sqrt{1+\alpha^2}} - R \sin \alpha \cos(\alpha \beta) + R \alpha \cos \alpha \right) \right] d\alpha d\beta$$

$$\psi(R, \frac{R}{d}) = \begin{cases} 1 & R < a \\ 0 & R > a \end{cases}$$

$\alpha \log(R \sqrt{1+\alpha^2} - R \alpha \cos \alpha + R \alpha \cos \alpha)$  / term  $\sim \alpha$   
 $\alpha$  変化  $\rightarrow$  振盪内, 変化  $\rightarrow$  高  $\rightarrow$   $R \alpha$ , order  $\sim \alpha$   
 振盪内  $\rightarrow$   $R \alpha$  の変化以外  $\rightarrow$  生じる変化,  $R \ll d + \dots$   
 大体

$$\frac{\alpha R R}{2 R d} = \frac{\alpha R}{2 d}$$

↑  $\alpha \sim \alpha$

$$\alpha = \frac{m \lambda^2 e^2}{k^2 R} = \frac{\sqrt{m} \lambda^2 e^2}{k \sqrt{2} E}$$

$\lambda \rightarrow$  大  $\rightarrow$   $\lambda \sim 50$   $\rightarrow$   $\lambda \sim 1 + \dots$  (Particle, energy  $\sim$   
 増える  $\alpha \sim \dots$ )

$$R \ll \frac{d}{10}$$

↑  $\alpha \sim \log[\dots]$ , term  $\sim$

$$\sim \log R d (1 + \cos \alpha)$$

↑  $\alpha \sim \alpha$

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故  $\beta$  による積分の 容易 = 遂行出来

$$\psi(R, \frac{\pi}{2}) = 2i \int_0^{\frac{\pi}{2}} A(x) J_0(RR \sin \alpha) e^{-iRa \cos \alpha + i\alpha \log R a} Ra (H \cos \alpha) d\alpha \quad (5)$$

$RR = \lambda$ , ~~###~~  $\sin \alpha = x$ ,  $t = \frac{\pi}{2} - \alpha$

$$\psi(R, \frac{\pi}{2}) = \begin{cases} \int_0^{\lambda} & \lambda < Ra \\ \int_0^{\pi/2} & \lambda > Ra \end{cases}$$

$$= 2i \int_0^{\lambda} \frac{A(\sin^{-1} x)}{x \sqrt{1-x^2}} e^{-iRa \sqrt{1-x^2} + i\alpha \log R a} Ra (H \sqrt{1-x^2}) J_0(\lambda x) x dx$$

$\frac{A(x)}{2\pi}$

$$= \int_0^{\lambda} \bar{A}(x) J_0(\lambda x) x dx \quad (6)$$

$\bar{A}(x)$  なる  $x$  の  $\lambda$  による積分の contribution については、  
 積分の上端  $\lambda$  による積分  
 上、同様

$$R \ll \frac{a}{10} \text{ かつ } \lambda = RR \ll \frac{Ra}{10}$$

1 能関  $\bar{A}(x)$  成立する  $\frac{a}{10} \approx 1$  なる  $t$  なる  $\lambda$ ,  $\lambda \ll 10 \sqrt{E}$   
 (E: volt unit) となる  $\lambda$  なる  $\lambda$ , 1. 非常  $\lambda$  なる  $\lambda$   
 なる  $\lambda$  なる  $\lambda$

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故 = Hamble, 逆演算. 近似的に = 応用出来.

$$\begin{aligned}
 \bar{A}(u) &= \frac{2\pi A(\beta_0 - i\epsilon)}{x\sqrt{1-x^2}} \int_0^{k_0} -i\epsilon d\sqrt{1-x^2} + i\alpha h_y k d \cdot (1+\sqrt{1-x^2}) \\
 &= \int_0^{k_0} 4(\beta_0 - \frac{\pi}{2}) J_0(u) \lambda d\lambda \\
 &= \int_0^{k_0} J_0(u) \lambda d\lambda \\
 &= \frac{k_0}{x} J_1(k_0 x)
 \end{aligned}$$

故 =

$$A(\alpha) = \frac{k_0}{2\alpha} \cos J_1(k_0 x) \cdot e^{i\alpha d \cos \alpha - i\alpha h_y k d (1 + \cos \alpha)} \quad \textcircled{7}$$

故 =

Let Coulomb field wave, scattered part.

$$\frac{1}{k_r - k_{i0}} e^{i\alpha h_y k d (k_r - k_{i0}) + i\alpha + 2i\theta_0} \quad \textcircled{8}$$

$$= A(\alpha) \gamma \theta \gamma_i \quad \alpha, \beta = \text{positive } i\alpha \text{ or } -i\alpha$$

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$f(\theta) = \int A(\alpha) \frac{\alpha}{R} \frac{e^{-i\alpha \log(1 - \sin\theta \sin\alpha \cos\beta - \cos\theta \cos\alpha)}}{1 - \sin\theta \sin\alpha \cos\beta - \cos\theta \cos\alpha} d\alpha d\beta$

$I(\theta) = |f(\theta)|^2$

Scattered intensity

SHI.  $f(\theta)$  / 奇点  $\beta = \pi/2$   $\alpha \rightarrow 2\pi$ ,  $\alpha = \pi$   $\theta = \pi$   
 $\theta = \delta$ ,  $\theta + \delta$  ( $\delta, \pi$ ) /  $\theta = \pi$   $\theta = 0$   $\theta = \pi$  奇点分る  
 $\delta \rightarrow 0$ , limiting case  $\pi$  奇点  $\alpha = \pi$   $\theta = \pi$   $\theta = \pi$

$A(\alpha) = \mathcal{D} \cdot H(\alpha)$

$|f(\theta)| = \frac{\alpha a}{2\pi} \int_{-\pi}^{\pi} d\alpha \frac{e^{iR\alpha \cos\alpha - i\alpha \log(1 - \sin\theta \sin\alpha \cos\beta - \cos\theta \cos\alpha)}}{1 - \sin\theta \sin\alpha \cos\beta - \cos\theta \cos\alpha} \alpha \beta$

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V

Coulomb field scattering, wave function, incidence (方向)  $k \rightarrow k'$

$$\psi \approx \left(1 - \frac{\beta^2}{i(kr - k'r')} e^{i(kr - k'r')} + \frac{\beta}{kr - k'r'} e^{i(kr - k'r')} + i\beta \log(kr - k'r') + i\alpha + 2i\eta_0 \right) e^{i(kr - k'r')} \quad (1)$$

$$\beta = \frac{2\epsilon Z Z' e^2}{k v} = \frac{m Z Z' e^2}{k^2 R}$$

$$e^{2i\eta_0} = \frac{\Gamma(1+i\alpha)}{\Gamma(1-i\alpha)}$$

上式、 $k \rightarrow k'$  平行  $\pi + i\beta$  限、 $\Gamma$  充分大  $\neq k \rightarrow k'$  成る。

$\beta - i\alpha =$

$$A(\alpha) = \frac{k\alpha}{2i} \frac{e^{i\beta \cos \theta - i\beta \log k\alpha(1+\cos \theta)}}{1 - \frac{\beta^2}{i\beta\alpha(1+\cos \theta)}} = \frac{J_1(k\alpha \sin \theta) \cos \theta}{\dots} \quad (2)$$

\*  $A(\alpha)$  の  
 $\theta = \pi$   
 $\beta$   
 $i\beta\alpha(1+\cos \theta)$   
 1項、非常  
 小  $\neq \beta$   
 neglect して  
 主要項

$\pi$  の  $\beta$ ,  $\alpha, \beta = \dots$  積分  $\neq \dots$ ,  $z = r \cos \theta = -d = \pi$

$$\tan \theta > a + i\alpha$$

$$\tan \theta < a + i\alpha$$

$\pi$  の  $\beta$  + wave function  $\neq$  得る。

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之ヲ算算ノ形式ニ立証明シテ見ル！  
 マッ ①ノ項一カニ ②ノカケテ見分ズル。

$$4 \approx \int_0^{\pi/2} \int_0^{\pi/2} \left( 1 - \frac{j^2}{i R r (1 - \sin \alpha \cos(\theta - \beta) - \cos \theta \cos \alpha)} \right) e^{i R r (\sin \theta \sin \alpha \cos(\theta - \beta) + \cos \theta \cos \alpha)} \times e^{i \log R r (1 - \sin \theta \sin \alpha \cos(\theta - \beta) - \cos \theta \cos \alpha)} A(\alpha) d\alpha d\theta \quad (3)$$

$A(\alpha) = \int_0^{\pi/2} J_1(R r \sin \alpha) + \text{factor}$  等々  $\alpha > \frac{10^4}{R r} \approx 10^{-4}$

マッ 3位ノ zero 十カケテ見分ズル。

$(1 - \sin \theta \sin \alpha \cos(\theta - \beta) - \cos \theta \cos \alpha)$  1.  $\theta = \alpha, \theta = \beta$  1.  $\theta = \pi - \alpha, \theta = \pi - \beta$  zero

マッ 他ノ場合ニ positive 等々 十カケテ見分ズル。  
 被積分項ニ singularity 等々 十カケテ見分ズル。

マッ 値ニ  $\frac{j^2}{R r (1 - \sin \theta \sin \alpha \cos(\theta - \beta) - \cos \theta \cos \alpha)}$  1. max 11.  $\theta = \beta, \alpha = \frac{\pi}{2}, \theta = \pi$  11  
 被積分項  $\frac{j^2}{R r (1 - \sin \theta)}$  等々 十カケテ見分ズル。 1. 連続等々 十カケテ見分ズル。  $\frac{j^2}{R r (1 - \sin \theta)}$

マッ 11. 等々 十カケテ見分ズル。

$$1 - \sin \theta > \frac{j^2}{R r} \approx 10^{-7} \quad (\because d_{max} \approx 10^3)$$

マッ 11. 等々 十カケテ見分ズル。

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①  $A(\alpha)$  は  $\delta$ -function, 極 =  $\alpha$  へ ~~積分~~  $\int_{-\infty}^{\infty} A(\alpha) d\alpha$

$$\chi \approx \left(1 - \frac{d^2}{iRr(1-\cos\theta)}\right) e^{i\theta \log Rr(1-\cos\theta)}$$

$$\times \int_0^{2\pi} e^{iRr(\sin\theta \sin\alpha \cos(\theta-\beta) + \cos\theta \cos\alpha)} A(\alpha) d\alpha d\beta$$

$J_0(z) = \frac{1}{\pi} \int_0^{\pi} e^{iz \cos\beta} d\beta$  (公式)  $\Rightarrow A(\alpha) = \beta = \pi/2$  / 積分  $\Rightarrow$  通関数  $J_0$

$$\chi \approx 2\pi \left(1 - \frac{d^2}{iR(r-2)}\right) e^{i\theta \log R(r-2)}$$

$$\times \int_0^{\pi} e^{iRr \cos\theta \cos\alpha} J_0(Rr \sin\theta \sin\alpha) A(\alpha) d\alpha$$

② 式,  $A(\alpha) = \alpha$ , 1. + 1 部分  $\Rightarrow$  簡単 = +177

$$A(\alpha) = \frac{R\alpha}{2\pi} \frac{e^{-i\theta \log 2R\alpha}}{1 - \frac{d^2}{i \cdot 2R\alpha}}$$

$\Rightarrow$  7 行 2 行 上 1 式 = 177  $A(r, \alpha)$

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$$\psi \approx Ra \frac{1 - \frac{t^2}{i(Rr-Rz)}}{1 - \frac{t^2}{i.2Rd}} e^{i\theta \log \frac{t-z}{2d}}$$

$$\times \int_0^{\frac{\pi}{2}} e^{iR(r \cos \theta + d) \cos \alpha} J_1(Ra \sin \alpha) J_0(Rr \sin \theta \sin \alpha) \cos \alpha \, d\alpha$$

この式は  $\frac{t^2}{Rr(1-\sin \theta)} \ll 1$  の場合  $r \sin \theta \ll t - \frac{t^2}{R}$  成立する式である

(4)  $r \cos \theta = z = -d$  の場合

$$\psi(r, \theta) \approx Ra \int_0^{\frac{\pi}{2}} J_1(Ra \sin \alpha) J_0(Rr \sin \theta \sin \alpha) \cos \alpha \, d\alpha$$

$$\approx Ra \int_0^{\infty} J_1(Rax) J_0(Rr \sin \theta x) \, dx$$

この式は前記で証明した通り、  
 $Ra \int_0^{\infty} J_1(Rax) J_0(Rr \sin \theta x) \, dx = \begin{cases} \sin \theta < a \\ r \sin \theta > a \end{cases}$

したがって、 $r \cos \theta = z = -d$  の場合、 $r \sin \theta$  が  $a$  より小さいと  $\psi = \sin \theta$ 、 $r \sin \theta$  が  $a$  より大きいと  $\psi = r \sin \theta$  となる。

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一般 (4) 式、 $\rho$  に対する波動関数 (slit に対する plane wave) の wave function, incident part を  $\rho$  に対する波動関数として、 $\rho$  に対する  
出来かたを示す。

