

YHAL F08 035

Vector Meson Theory = π -n Proton, Anomalous
Magnetic Moment, 計算.

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Vector Meson Theory = \exists Proton, Anomalous
 Magnetic Moment, 計算。

3次, 攝動計算ヲ行フ。次, 3次の過程が可能ナル。

(1) $P + (H) \xrightarrow{NM} N + M_k^+ + (H) \xrightarrow{ME} N + M_k^+ + (H) \xrightarrow{NM} P + (H)$

(2) $P + (H) \xrightarrow{NM} N + M_k^+ + (H) \xrightarrow{NM} P + M_k^0 + M_k^+ + (H) \xrightarrow{ME} P + (H)$

(3) $P + (H) \xrightarrow{ME} P + M_k^+ + M_k^0 + (H) \xrightarrow{NM} N + M_k^+ + (H) \xrightarrow{NM} P + (H)$

P: Proton; N: Neutron; M_k^+ : Neutron; M_k^0 : 運動量 \vec{k} ヲ

持ッテ正荷電, Meson; M_k^- : 運動量 \vec{k} ヲ持ッテ負荷電

, Meson; (H): 空方向-強Hヲ持ッテ-挿+磁場。

\xrightarrow{ME} , \xrightarrow{NM} " 挿+ Meson + Electro-magnetic field,

Heavy Particle + Meson, interaction = \exists Processヲ

示ス。

1. Interaction Energies + Matrix Elements.

$$H_{NM}^I = \tilde{u} \cdot \frac{f}{g} \frac{1}{x} [\Pi(\vec{\sigma}, \text{curl } \vec{\phi}) + \Pi^*(\vec{\sigma}, \text{curl } \vec{\phi}^*)] \cdot u,$$

$$H_{NM}^{II} = \tilde{u} \cdot \frac{g}{x} [\Pi \text{div } \vec{\psi} + \Pi^* \text{div } \vec{\psi}^*] \cdot u$$

$$H_{NM} = H_{NM}^I + H_{NM}^{II}$$

Π^* : Proton, Neutron = $\frac{1}{2} \sigma \cdot u$ operator,

$\vec{\phi}$: Meson, wave function, Transverse part.

$\vec{\psi}$: Meson, wave function, longitudinal part.

$\vec{\sigma}$: Pauli, spin vector (of the heavy particle).

$\vec{\phi}, \vec{\phi}^*$ は、振動場である。

$$\left. \begin{aligned} \vec{\phi} &= \frac{1}{\sqrt{V}} \sum_{\vec{k}} i \sqrt{\frac{2\pi\hbar^2 c^2}{\epsilon_k}} (a_{\vec{k}} - b_{\vec{k}}^*) \cdot \vec{j} \cdot e^{i(\vec{k}, \vec{r})} \\ \vec{\phi}^* &= \frac{1}{\sqrt{V}} \sum_{\vec{k}} -i \sqrt{\frac{2\pi\hbar^2 c^2}{\epsilon_k}} (a_{\vec{k}}^* - b_{\vec{k}}) \cdot \vec{j} \cdot e^{-i(\vec{k}, \vec{r})} \end{aligned} \right\} \quad (3)$$

$\epsilon_k = \hbar c \sqrt{k^2 + \mu^2}$; $\mu = \frac{mc}{\hbar}$; μ : Meson, rest mass
 \vec{j} : \vec{k} = 単位 + unit vector.

- $[a_{\vec{k}}, a_{\vec{k}'}^*] = [\epsilon_{\vec{k}}, \epsilon_{\vec{k}'}^*] = \delta_{\vec{k}\vec{k}'}$
- 他, a 's, b 's, a 's と b 's 互に commute する。
- $a_{\vec{k}}^*$: $M_{\vec{k}}^+$ の 1 個 emit する operator.
- $b_{\vec{k}}^*$: $M_{\vec{k}}^-$ の " " " " " " " "
- $a_{\vec{k}}$: $M_{\vec{k}}^+$ の 1 個 absorb する operator.
- $b_{\vec{k}}$: $M_{\vec{k}}^-$ の " " " " " " " "

2-成分, spin 及び 上下反対符号, spin 7 成分 \rightarrow heavy particle, wave function \neq \neq

$$\left. \begin{aligned} u^{(A)} &= \alpha e^{i(\vec{k}, \vec{r})} = \begin{pmatrix} \alpha_+ \\ \alpha_- \end{pmatrix} \cdot e^{i(\vec{k}, \vec{r})} \\ u^{(B)} &= \beta \cdot e^{i(\vec{k}, \vec{r})} = \begin{pmatrix} \beta_+ \\ \beta_- \end{pmatrix} \cdot e^{i(\vec{k}, \vec{r})} \end{aligned} \right\} \quad (4)$$

↑ z 成分

$$\left. \begin{aligned} \alpha_+ &= 1 \\ \alpha_- &= 0 \end{aligned} \right\} \quad \left. \begin{aligned} \beta_+ &= 0 \\ \beta_- &= 1 \end{aligned} \right\} \quad (4')$$

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$$\tilde{\alpha}^* + \beta^* = 1 \quad (\text{unit matrix}) \quad (5)$$

initial state is intermediate state =  $\pi^+ e^- \mu^-$  heavy particle, wave function  $\neq u^{(A)}, u^{(B)}, u^{(C)} \approx \pi^+ \approx \bar{\psi} = \bar{u}$

$$H_{MM}^{NM} = \frac{1}{\kappa} \frac{1}{\sqrt{L^3}} \sum_{\vec{k}} \sqrt{\frac{2\pi\hbar^2 c^2}{\epsilon_k}} \tilde{u}^{(A)} \cdot (a_{\vec{k}}^* - b_{\vec{k}}) (\vec{\sigma} \cdot [\vec{j} \times \vec{k}]) \cdot u^{(A)} \prod^* e^{-i(\vec{k}-\vec{k}') \cdot \vec{r}} + \text{compl. conj.} \quad (6)$$

$$(3) \rightarrow (1) = \lambda \vec{r} \quad H_{ME}^{ME} = \frac{\hbar c}{2L^3} \sum_{\vec{k}, \vec{k}'} \frac{1}{\sqrt{\epsilon_k \epsilon_{k'}}} (a_{\vec{k}} - b_{\vec{k}}^*) (\vec{\sigma} \cdot [\vec{j} \times \vec{k}]) \cdot u^{(A)} \prod^* e^{-i(\vec{k}-\vec{k}') \cdot \vec{r}} + \text{compl. conj.} \quad (7)$$

(6), (7) の 2 次, Matrix elements を得る.

$$\begin{aligned} H_{M_k^+ \text{ emission}}^{MM} &= \frac{1}{\sqrt{L^3}} \frac{1}{\kappa} \sqrt{\frac{2\pi\hbar^2 c^2}{\epsilon_k}} \tilde{u}^{(A)} \cdot (\vec{\sigma} \cdot [\vec{j} \times \vec{k}]) \cdot A \equiv H_{k0} \\ H_{M_k^- \text{ absorption}}^{MM} &= \frac{1}{\sqrt{L^3}} \frac{1}{\kappa} \sqrt{\frac{2\pi\hbar^2 c^2}{\epsilon_k}} \tilde{u}^{(A)} \cdot (\vec{\sigma} \cdot [\vec{j} \times \vec{k}]) \cdot I \equiv H_{0k} \\ H_{M_k^+ \rightarrow M_k^+}^{ME} &= \frac{\hbar c}{2L^3} \frac{1}{\sqrt{\epsilon_k \epsilon_{k'}}} \int \{ (\vec{j} \times \vec{k}) \cdot [\vec{A} \times \vec{j}'] + (\vec{j} \times \vec{k}') \cdot [\vec{A} \times \vec{j}] \} \\ &\quad \cdot e^{i(\vec{k}-\vec{k}') \cdot \vec{r}} d\vec{r} \equiv H_{kk} \end{aligned}$$

$$H_{M_k^- \text{ absorption}}^{MM} = -\frac{1}{\sqrt{L^3}} \frac{1}{\kappa} \sqrt{\frac{2\pi\hbar^2 c^2}{\epsilon_k}} \tilde{u}^{(A)} \cdot (\vec{\sigma} \cdot [\vec{j} \times \vec{k}]) \cdot I \equiv H'_{k0}$$

$$H_{M_k^- \text{ absorption}}^{MM} = -\frac{1}{\sqrt{L^3}} \frac{1}{\kappa} \sqrt{\frac{2\pi\hbar^2 c^2}{\epsilon_k}} \tilde{u}^{(A)} \cdot (\vec{\sigma} \cdot [\vec{j} \times \vec{k}]) \cdot A \equiv H'_{k0k}$$

$$H_{M_k^+ M_k^- \text{ emission}}^{ME} = -\frac{\hbar c}{2L^3} \frac{1}{\sqrt{\epsilon_k \epsilon_{k'}}} \int \{ (\vec{j} \times \vec{k}) \cdot [\vec{A} \times \vec{j}'] + (\vec{j} \times \vec{k}') \cdot [\vec{A} \times \vec{j}] \} \cdot e^{i(\vec{k}-\vec{k}') \cdot \vec{r}} d\vec{r} \equiv H_{kk}^e \quad A$$



$$W_I = -\frac{\pi f^2 e}{\hbar c \kappa^2} \sum_{\vec{k}} \frac{1}{kR} (k^2 + \kappa^2)(k^2 + \kappa^2) \left\{ (\vec{k}\vec{k}') (\vec{A}, \vec{k} + \vec{k}') \right. \\ \left. + k^2 (\vec{A}, \vec{k}) + k'^2 (\vec{A}, \vec{k}') \right\} e^{i(\vec{k}-\vec{k}', \vec{r})} d\vec{r}$$

$$- \frac{i\pi f^2 e}{\hbar c \kappa^2} \sum_{\vec{k}} \frac{1}{kR} (k^2 + \kappa^2)(k^2 + \kappa^2) \left\{ (\vec{A}, \vec{k} + \vec{k}') [\vec{k} \times \vec{k}]_z \right. \\ \left. - (\vec{A}, [\vec{k} \times \vec{k}']) (k_z + k'_z) \right\} e^{i(\vec{k}-\vec{k}', \vec{r})} d\vec{r}$$

$$L_1 \rightarrow \infty \quad \text{t sum} \quad \frac{1}{L^3} \sum_{\vec{k}} \rightarrow \frac{1}{(2\pi)^3} \int d\vec{k} \\ \int_V d\vec{r} \rightarrow \int_{\text{全空間}} d\vec{r} \quad \text{t} \rightarrow \tau$$

$$W_I = -\frac{\pi f^2 e}{64\pi^6 \hbar c \kappa^2} \int \frac{1}{(k^2 + \kappa^2)(k'^2 + \kappa^2)} \left\{ (\vec{k}\vec{k}') (\vec{A}, \vec{k} + \vec{k}') + k^2 (\vec{A}, \vec{k}) + k'^2 (\vec{A}, \vec{k}') \right\} \\ \cdot e^{i(\vec{k}-\vec{k}', \vec{r})} d\vec{r} d\vec{k} d\vec{k}' \\ - \frac{i\pi f^2 e}{64\pi^6 \hbar c \kappa^2} \int \frac{1}{(k^2 + \kappa^2)(k'^2 + \kappa^2)} \left\{ (\vec{A}, \vec{k} + \vec{k}') [\vec{k} \times \vec{k}]_z \right. \\ \left. - (\vec{A}, [\vec{k} \times \vec{k}']) (k_z + k'_z) \right\} e^{i(\vec{k}-\vec{k}', \vec{r})} d\vec{r} d\vec{k} d\vec{k}'$$

1st. integral,  $\vec{k}$  &  $\vec{k}'$  space = 赤々々 polar angle  $\Rightarrow \int d\Omega$   
 積分  $\Rightarrow \Rightarrow \Rightarrow$  vanishes. 又、部分積分等、用いて計算スル

Wanted

$$\int \frac{1}{(k^2 + \kappa^2)(k'^2 + \kappa^2)} \left\{ (\vec{A}, \vec{k} + \vec{k}') [\vec{k} \times \vec{k}]_z - (\vec{A}, [\vec{k} \times \vec{k}']) (k_z + k'_z) \right\} e^{i(\vec{k}-\vec{k}', \vec{r})} d\vec{r} d\vec{k} d\vec{k}'$$

$$= + 4\pi^3 i H \int \frac{(\vec{k} + \vec{k}', \vec{k}) + (k_z + k'_z) \cdot k_z}{(k^2 + \kappa^2)(k'^2 + \kappa^2)} \delta(\vec{k} - \vec{k}') d\vec{k} d\vec{k}'$$

$$= 8\pi^3 i H \int \frac{k^2 + k_z^2}{(k^2 + \kappa^2)^2} d\vec{k}$$

$$= 8\pi^3 i H \frac{4\pi}{3} \int \frac{k^4 dk}{(k^2 + \kappa^2)^2}$$

$$\begin{aligned} \therefore W_I(H) &= - \frac{i \pi f e^2}{64 \pi^5 \cdot hc \kappa^2} \cdot \frac{128 \pi^4 \cdot i H}{3} \int \frac{k^4 dk}{(\kappa^2 + \nu^2)^2} \\ &= - \frac{2 f^2 e H}{3 \pi \cdot hc \kappa^2} \int \frac{k^4 dk}{(\kappa^2 + \nu^2)^2} \end{aligned}$$

3. Process (2)  $B \bar{e} \nu (3) = \exists \sim$  Perturbation Energy.

$W_I, W_{II} \sim \int \vec{p} \cdot \vec{p} = \nu^2$

$$W_{II} = \sum_{k, \vec{k}} \sum_{j, j'} \frac{H_{k0}^\alpha \cdot H_{k0}^\alpha}{\epsilon_k (\epsilon_k + \epsilon_k)}$$

$\hat{z} = (7) \rightarrow$  用  $e^{-t}$

$$W_{II} = \frac{\pi f^2 e \hbar^3}{\kappa^2 L^6} \sum_{k, \vec{k}} \sum_{j, j'} \frac{1}{\epsilon_k^2 \epsilon_k (\epsilon_k + \epsilon_k)} \int \tilde{A}(\vec{\sigma}, [\vec{j} \times \vec{k}]) I \cdot \tilde{I}(\vec{\sigma}, [\vec{j} \times \vec{k}]) A \cdot [([\vec{j} \times \vec{k}], [\vec{A} \times \vec{j}]) + ([\vec{j} \times \vec{k}], [\vec{A} \times \vec{j}])] \cdot e^{i(\vec{k} - \vec{k}, \vec{r})} d\vec{r}$$

$|\vec{r}| \vec{p} = \nu^2$

$$W_{III} = \frac{\pi f^2 e \hbar^3}{\kappa^2 L^6} \sum_{k, \vec{k}} \sum_{j, j'} \frac{1}{(\epsilon_k + \epsilon_k) \cdot \epsilon_k \epsilon_k^2} \int \tilde{A}(\vec{\sigma}, [\vec{j} \times \vec{k}]) I \cdot \tilde{I}(\vec{\sigma}, [\vec{j} \times \vec{k}]) A \cdot [([\vec{j} \times \vec{k}], [\vec{A} \times \vec{j}]) + ([\vec{j} \times \vec{k}], [\vec{A} \times \vec{j}])] \cdot e^{i(\vec{k} - \vec{k}, \vec{r})} d\vec{r}$$

$$\frac{1}{\epsilon_k^2 \epsilon_k (\epsilon_k + \epsilon_k)} + \frac{1}{(\epsilon_k + \epsilon_k) \epsilon_k \epsilon_k^2} = \frac{1}{\epsilon_k^2 \epsilon_k^2} = \frac{1}{\hbar^2 c^4 (\kappa^2 + \nu^2)}$$

$\vec{r} \cdot \vec{p} = \nu^2$

$$W_{II} + W_{III}$$

$$= \frac{\pi f^2 e}{\hbar c \kappa^2 L^6} \sum_{k, \vec{k}} \sum_{j, j'} \frac{1}{(\kappa^2 + \nu^2)(\kappa^2 + \nu^2)} \int \tilde{A}(\vec{\sigma}, [\vec{j} \times \vec{k}]) I \cdot \tilde{I}(\vec{\sigma}, [\vec{j} \times \vec{k}]) A \cdot [([\vec{j} \times \vec{k}], [\vec{A} \times \vec{j}]) + ([\vec{j} \times \vec{k}], [\vec{A} \times \vec{j}])] \cdot e^{i(\vec{k} - \vec{k}, \vec{r})} d\vec{r}$$

$$= W_I$$

P.

従って全振動エネルギー  $W(H)$  は

$$W(H) = W_I + W_{II} + W_{III}$$

$$= 2W_I.$$

$$= \frac{4f^2 e^2 H}{3\pi \hbar c^3} \int \frac{k^4 dk}{(k^2 + \kappa^2)^2}$$

