Thermodynamic bound on the asymmetry of cross-correlations Naruo Ohga, Sosuke Ito, Artemy Kolchinsky, arXiv:2303.13116,

Ito group, The University of Tokyo Naruo Ohga

Aug. 4, 2023 YITP-YSF symposium

temy Kolchinsky, arXiv:2303.13116, (Phys. Rev. Lett., in press.)

Two-time correlations

Fluctuating system in steady state



Steady-state trajectory keeps fluctuating

a, *b* : Arbitrary two observables (Even under time-reversal)

Two noneq. features of two-time correlations 1 Asymmetry of cross-correlations 2 Oscillations

Two-time correlations

$$C_{ba}^{\tau} = \langle b(t+\tau) \, a(t) \rangle$$

$$\begin{bmatrix} C_{aa}^{\tau}, C_{bb}^{\tau} : \text{Auto-correlation} \\ C_{ba}^{\tau}, C_{ab}^{\tau} : \text{Cross-correlation} \end{bmatrix}$$

- Captures the temporal structure at the trajectory level
- Experimentally accessible in various systems

I Asymmetry of cross-correlations **2** Oscillations (Eigenvalues)

Discussion & Summary

Asymmetry in cross-correlations

Equilibrium

Noneq. Steady state

$C_{ba}^{\tau} = C_{ab}^{\tau}$ for any a, b(Microscopic reversibility)

 $C_{ba}^{\tau} \neq C_{ab}^{\tau}$ for some a, b(Breaking of micro. rev.)



Asymmetry in cross-correlations

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(Breaking of micro. rev.)

 $C_{ba}^{\tau} \neq C_{ab}^{\tau}$: fundamental signature of nonequilibrium

Biochemical reaction



D. Sisan+ Biophys J. 2010

Biological motion Turbulence



A. Jachens+, J. Fluid. Mech. 2006



2016



Goal

Asymmetry $C_{ba}^{\tau} \neq C_{ab}^{\tau}$: Biochemical reaction



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2016

Any universal & **quantitative** relation? Thermodynamic driving Temperature Mechanical gradient force_z





J. R. Gomez-Solano+, EPL 2010

Chemical potential gradient







Dimensionless measure of the asymmetry of cross-correlation (New in this study)

Asymmetry of cross-corr.

$$\chi_{ba} = \lim_{\tau \to 0} \frac{C_{ba}^{\tau} - C_{ab}^{\tau}}{2\sqrt{(\Delta_{\tau}C_{aa})(\Delta_{\tau}C_{bb})}}$$
Decay of auto-corr.

$$\begin{bmatrix} C_{ba}^{\tau} = \langle b(t+\tau) a(t) \rangle \\ \Delta_{\tau}C_{aa} = C_{aa}^{0} - C_{aa}^{\tau} \\ = \frac{1}{2} \langle [a(t+\tau) - a(t)]^{2} \rangle \\ \text{(The change speed of } a) \end{bmatrix}$$
Invariant under rescaling of a, b
and time. Experimentally accessible

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Strength of thermodynamic driving (Standard in discrete systems)

Cycle *c*

Cyclic sequence of distinct states connected with allowed transitions

Cycle affinity \mathcal{F}_c

- = The sum of thermodynamic forces over one turn of the cycle
- = Dissipation in the environment per one turn of the cycle

Determined by environmental parameters









- Cycle affinity \mathcal{F}_c
- = The sum of thermodynamic forces over the cycle
- = Dissipation per one cycle



$$|\chi_{ba}| \leq \max \frac{1}{12}$$







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$$|\chi_{ba}| \leq \max \frac{\tan}{ta}$$

The magnitude of asymmetry is universally related to affinity!







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Implications





Numerical example





Discrete-state Markov process States i = 1, ..., N

Transition matrix R

 $R_{ij} = \text{Transition rate from } j \text{ to } i \quad (i \neq j)$ $R_{ii} = -\sum_{j:j \neq i} R_{ji}$ (Escape rate)

Time evolution

 $d\mathbf{p}(t)/dt = \mathsf{R}\mathbf{p}(t) \implies \mathbf{p}(t) = e^{\mathsf{R}t}\mathbf{p}(0)$ Steady-state distribution q: Rq = 0

Two-time correlation

$$C_{ba}^{\tau} = \langle b(\tau)a(0) \rangle = \sum_{ij} [e^{\mathsf{R}\tau}]_{ij} q_j b_i a_j$$

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Two-time correlation $C_{ba}^{\tau} = \langle b(\tau)a(0) \rangle = \sum_{ii} [e^{\mathsf{R}\tau}]_{ij} q_j b_i a_j$ q: Steady-state distribution Proof Use $[e^{R\tau}]_{ii} \approx 1 + \tau R_{ii}$ $\chi_{ba} = \lim_{\tau \to 0} \frac{C_{ba}^{\tau} - C_{ab}^{\tau}}{2\sqrt{\Delta_{\tau}C_{aa}} \Delta_{\tau}C_{bb}}}$ $\frac{2\sum_{ij} R_{ij}q_j (b_i a_j - a_i b_j)}{R_{ij}q_j [(a_i - \alpha)^2]}$ Area $\sum_{ij} R_{ij} q_j \left[(a_i - a_j)^2 + (b_i - b_j)^2 \right]$ Length

- 1 1

- :
- :

 (a_{i}, b_{i})

 (a_j, b_j) \mathcal{A}

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1 Asymmetry of cross-correlations 2 Oscillations (Eigenvalues)

Discussion & Summary

Fluctuating oscillations



Oscillations should be coherent in time for reliable biochemical functionality.





Fluctuating oscillations



Oscillations should be coherent in time for reliable biochemical functionality.



Finite-time corr. $C_{ba}^{\tau} = \sum_{ij} [e^{R\tau}]_{ij} q_j b_i a_j$ $\lambda_{\alpha} = -\mu_{\alpha} + i\omega_{\alpha}$: α -th eigenvalue of R $\rightarrow C_{ba}^{\tau}$ is a superposition of $\exp(\lambda_{\alpha}\tau) = \exp(-\mu_{\alpha}\tau) \exp(i\omega_{\alpha}\tau)$ Relaxation Oscillation

Assume that one of the modes α is dominant.

Quantitative measure of coherence # of oscillations before the auto- \langle correlation decays

 $(\mu_{\alpha})^{-1}$ Decay time ω_{α} Period $2\pi\mu_{\alpha}$ $2\pi |\omega_{\alpha}|$



Fluctuating oscillations



Any universal & quantitative relation?

Thermodynamic driving

Temperature Mechanical gradient force_z



J. R. Gomez-Solano+, EPL 2010

Chemical potential ATP gradient $\Delta \mu$

ADP + Pi





Conjectured by Barato and Seifert (2017) based on numerical evidence. Not rigorously proven before.

The coherence of oscillation is universally bounded by affinity



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Proof $\begin{aligned} -\mu_{\alpha} + i\omega_{\alpha} &= \frac{\partial}{\partial \tau} \exp(\lambda_{\alpha} \tau) \Big|_{\tau=0} = \frac{\partial}{\partial \tau} C_{z^{*}z}^{\tau} \Big|_{\tau=0} \\ Complex-valued \\ "Observable" \\ z_{i} &= u_{i}^{(\alpha)} / q_{i} \\ & \bullet \end{aligned}$



Main Result for any a, b

$$|\chi_{ba}| \le \max_{c} \frac{\tanh(\mathcal{F}_{c}/2n_{c})}{\tan(\pi/n_{c})} \le \max_{c} \frac{\mathcal{F}_{c}}{\mathcal{F}_{c}}$$

$$\begin{vmatrix} = \frac{\partial}{\partial \tau} \left[(C_{aa}^{\tau} + C_{bb}^{\tau}) + i (C_{ba}^{\tau} - C_{ab}^{\tau}) \right] \end{vmatrix}_{\tau=0}$$

ued Real-valued
" "Observable"
 $b_i = \operatorname{Im} z_i$
nvector $a_i = \operatorname{Re} z_i$





Proof $\begin{aligned}
-\mu_{\alpha} + i\omega_{\alpha} &= \frac{\partial}{\partial \tau} \exp(\lambda_{\alpha} \tau) \Big|_{\tau=0} = \frac{\partial}{\partial \tau} C_{z^{*}z}^{\tau} \Big|_{\tau} \\
\frac{|\omega_{\alpha}|}{\mu_{\alpha}} &= \lim_{\tau \to 0} \frac{|C_{ba}^{\tau} - C_{ab}^{\tau}|}{\Delta_{\tau} C_{aa} + \Delta_{\tau} C_{bb}} \leq \lim_{\uparrow \tau \to 0} \frac{1}{2\sqrt{\tau}} \\
x + y &\geq 2\sqrt{xy}
\end{aligned}$

We prove the conjecture as a corollary!

Main Result for any *a*, *b*

$$|\chi_{ba}| \le \max_{c} \frac{\tanh(\mathcal{F}_{c}/2n_{c})}{\tan(\pi/n_{c})} \le \max_{c} \frac{2}{2}$$

$$= \frac{\partial}{\partial \tau} \left[\left(C_{aa}^{\tau} + C_{bb}^{\tau} \right) + i \left(C_{ba}^{\tau} - C_{ab}^{\tau} \right) \right]_{\tau=0}$$

$$\frac{|C_{ba}^{\tau} - C_{ab}^{\tau}|}{\sqrt{(\Delta_{\tau} C_{aa})(\Delta_{\tau} C_{bb})}} = |\chi_{ba}| \le \text{(affinity bounce}$$



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Asymmetry of cross-correlations 2 Oscillations (Eigenvalues)



Affinity bounds vs. entropy bounds



Complementary ways to relate statistical and thermodynamic signatures

Affinity bounds vs. entropy bounds

Cycle affinity

Strength of

Environment (e.g., chemic

Dependence on the steady state

Determined by

Independen⁻

Response to microscopic kinetic changes

 \bigvee

Macroscop

 \rightarrow Complementary characterizations of the thermodynamic cost

sy \mathcal{F}_c	Entropy production Σ^{i}
driving Ital parameters cal potentials)	Resulting dissipation rate
t	Dependent
	Sensitive
bic setup	ب Microscopic behavio





Summary

Result

The asymmetry of cross-correlation is universally related to affinity

 $|\chi_{ba}| \le \max_{c} \frac{\tanh(\mathcal{F}_c/2n_c)}{\tan(\pi/n_c)} \le \max_{c} \frac{\mathcal{F}_c}{2\pi}$

$$\chi_{ba} = \lim_{\tau \to 0} \frac{C_{ba}^{\tau} - C_{ab}^{\tau}}{2\sqrt{\Delta_{\tau}C_{aa}} \Delta_{\tau}C_{bb}} \begin{bmatrix} C_{ba}^{\tau} = \langle b(t+\tau) a(t) \rangle \\ \Delta_{\tau}C_{aa} = C_{aa}^{0} - C_{aa}^{\tau} \end{bmatrix}$$

Proof: Use the isoperimetric inequality between area and length

Details \rightarrow N. Ohga, S. Ito, A. Kolchinsky, arXiv:2303.13116 See also N. Shiraishi, arXiv:2304.12775 T. Van Vu, V. T. Vo, K. Saito, arXiv:2305.18000

 \mathcal{F}_c : Cycle affinity (Strength of thermodynamic driving) n_c : # of states on cycle c

Corollary

The coherence of oscillation is universally bounded by affinity

$$\frac{|\omega_{\alpha}|}{\mu_{\alpha}} \le \max_{c} \frac{\tanh(\mathcal{F}_{c}/2n_{c})}{\tan(\pi/n_{c})} \le \max_{c} \frac{\mathcal{F}_{c}}{2\pi}$$

Conjectured by Barato & Seifert (2017) $\lambda_{\alpha} = -\mu_{\alpha} + i\omega_{\alpha}$: Eigenvalue of R

Proof: Extract λ_{α} from a complex-valued correlations

